

Dark Photon Study at BESIII

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on behalf of BESIII Collaboration

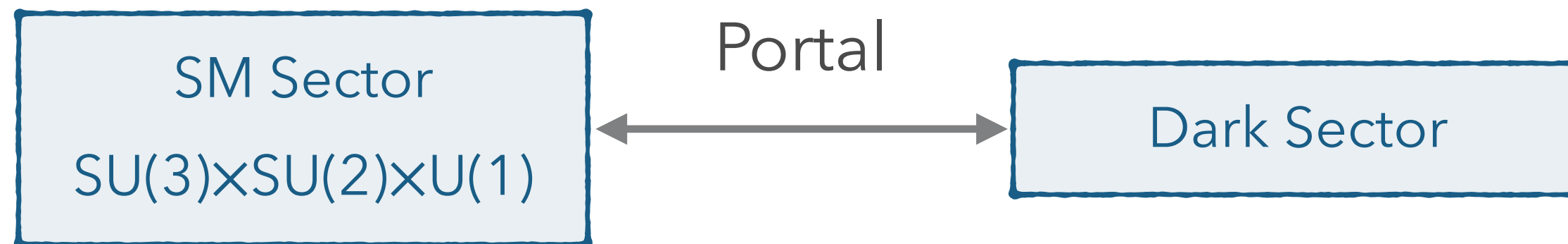
第三届地下和空间粒子物理与宇宙物理前沿问题研讨会

Conference on frontiers of underground and space particle physics and cosmophysics

2024年5月7日-11日，四川西昌

Dark Sector

- Consisting of (light) particles do not interact with the known strong, weak, or electromagnetic forces
- Weakly coupled to visible sector through a mediator or “portal”



| Portal | Operators |
|------------------------------------|--|
| Dark photon (Spin 1) | $\frac{\epsilon}{2} F_{\mu\nu} F'^{\mu\nu}$ |
| Axion (Spin 0) | $\frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{1}{f_a} \partial_\mu \bar{\psi} \gamma^\mu \gamma_5 \psi$ |
| Scalar (Spin 0) | $(\mu S + \lambda S^2) H^\dagger H$ |
| Sterile neutrino (Spin 1/2) | $y_N \bar{L} H N$ |

[arXiv:2005.01515v3 \[hep-ph\]](https://arxiv.org/abs/2005.01515v3)

Dark sectors beyond Dark Matter

In fact, beyond the DM motivation, many other open problems in particle physics let us think about dark particles.

- Models to address the **strong CP problem**. Axions and axion-like particles;
- Models to address the **gauge hierarchy problem** (relaxion);
- **SUSY** extended models (Next-to-Minimal-Supersymmetric-Standard-Model);
- Models for **baryogenesis**;
- Models for **neutrino** mass generation;
- Models addressing **anomalies in data**;
(g-2)_μ, galactic center excess for Dark Matter, B-physics anomalies, ...).

Some of these particles are naturally light thanks to approximate global symmetries.

From a phenomenological point of view, the signatures to search for are often similar

S.Gori

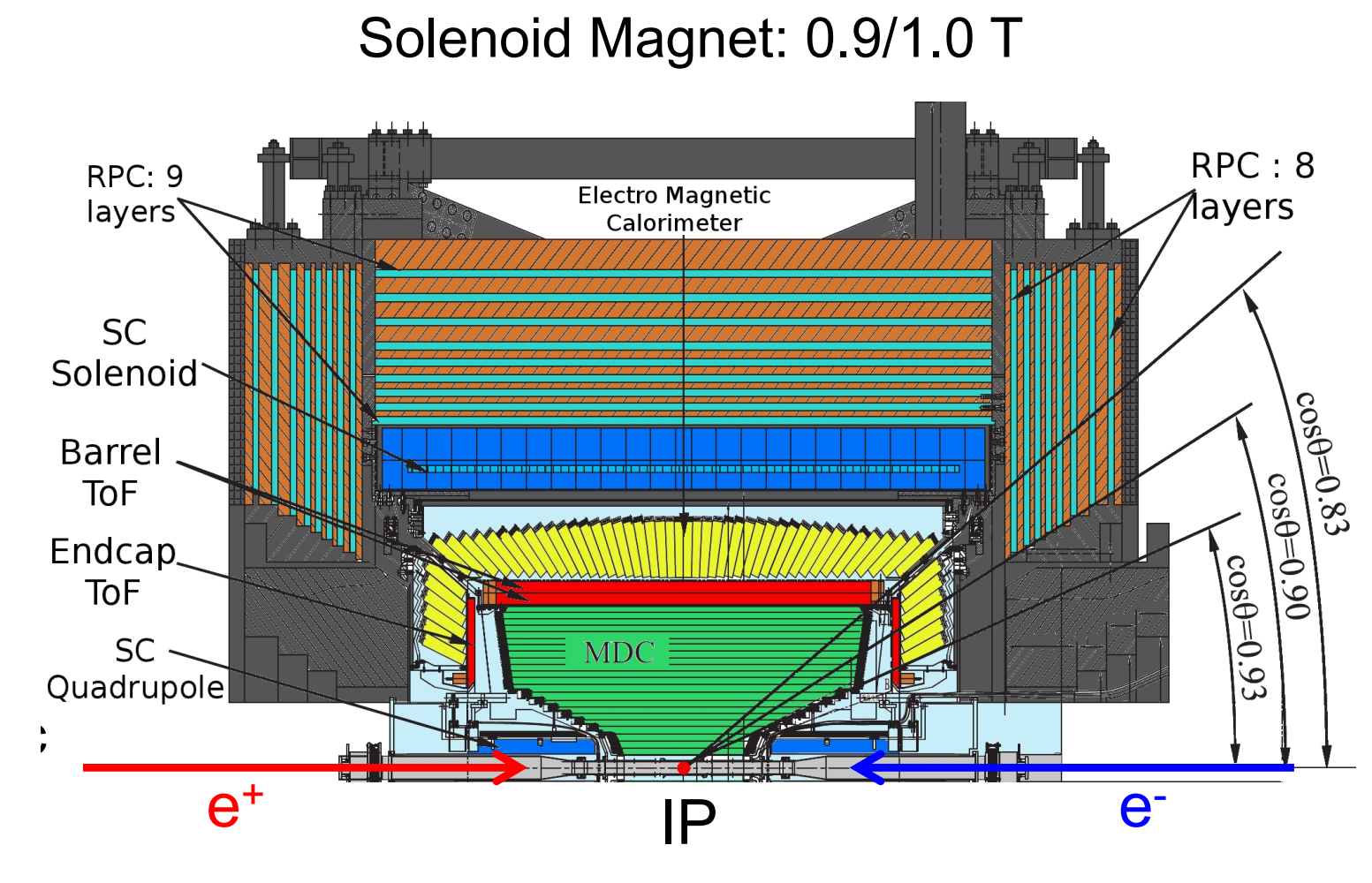
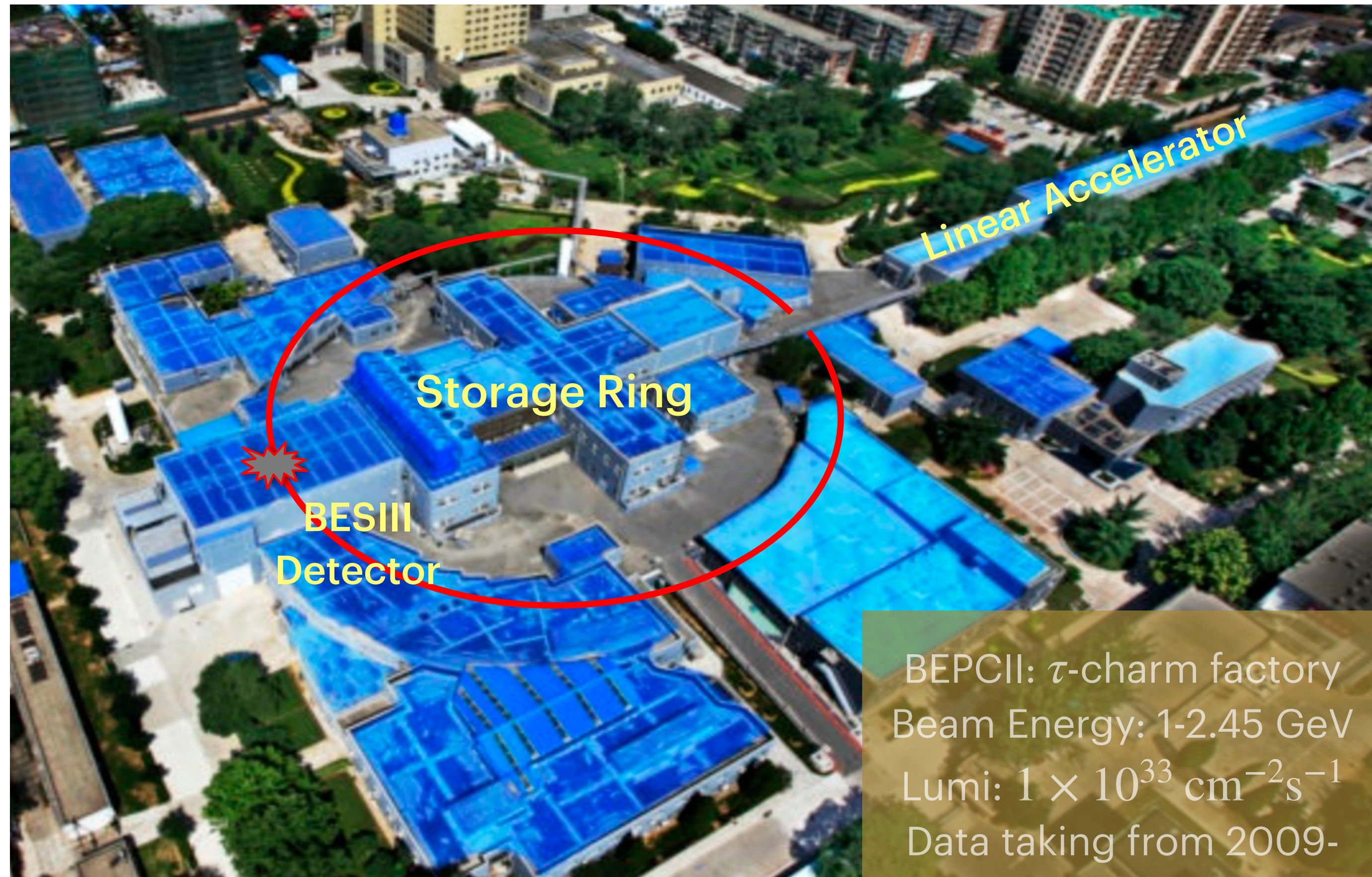
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Borrowed from S. Gori, Dark Sector at Snowmass HPS Collaboration Meeting, Nov.6, 2022

Dark Photon

- New Abelian gauge group U(1) force carrier
- Kinetic mixing with SM U(1) with mixing coefficient ϵ *B. Holdom, PLB 166,196 (1986)*
- Typical mix strength: $10^{-2} - 10^{-5}$, could be smaller
- Expected mass scale: MeV/ c^2 -GeV/ c^2 , can be accessible by high intensity e^+e^- collider experiments
- Massless dark photon and massive dark photon
 - Massless kind: does not couple directly to any of the SM currents and interacts with ordinary matter only through operators of dimension higher than four
 - Massive kind: couples to ordinary matter through a current

BEP CII and BES III



MUC $\sigma_{R\phi}$: 2 cm

TOF

σ_T : 80 ps
 110 ps (60 ps)

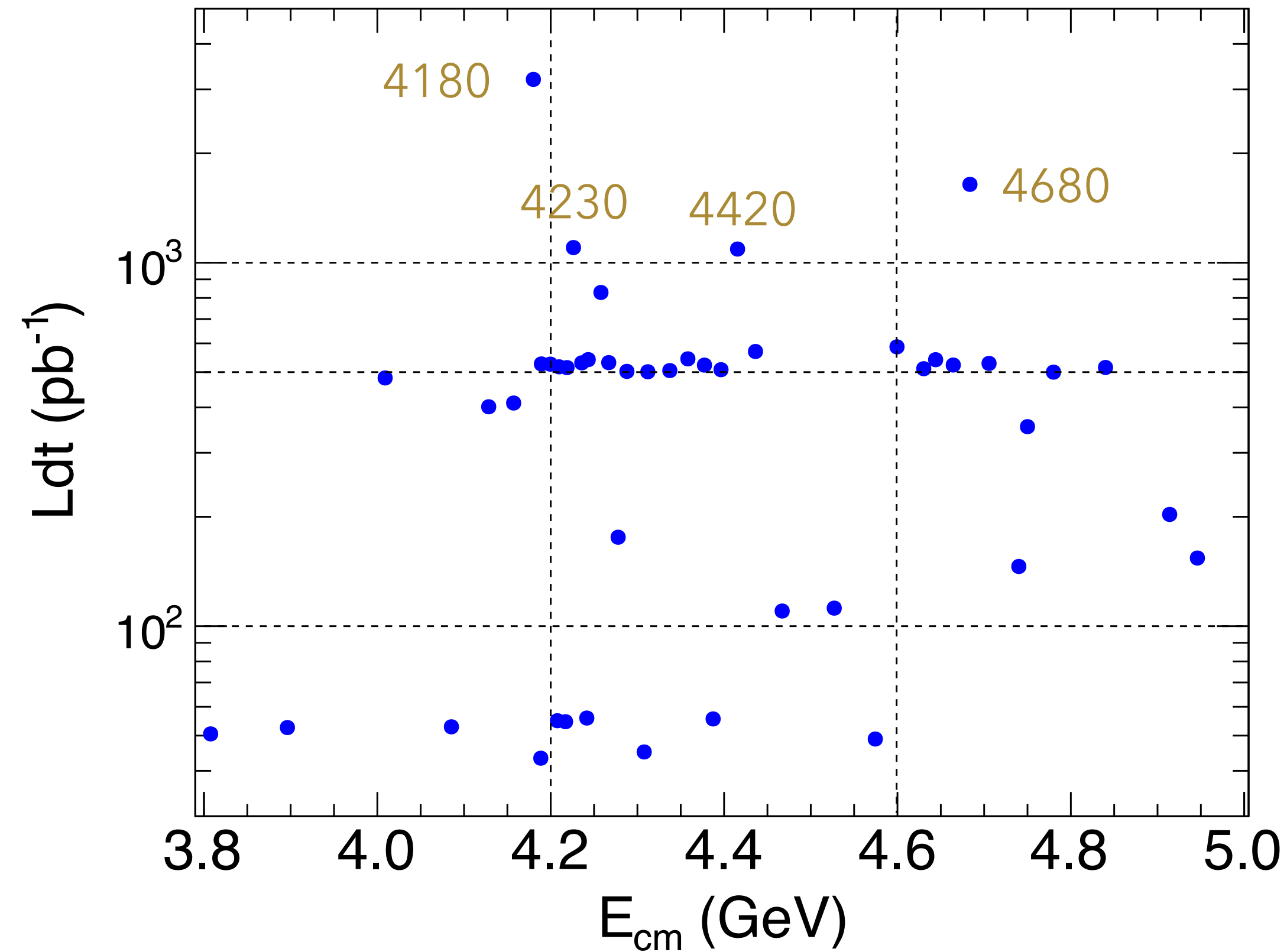
MDC

dE/dx: 6%
 σ_p/p : 0.5% at 1GeV/c

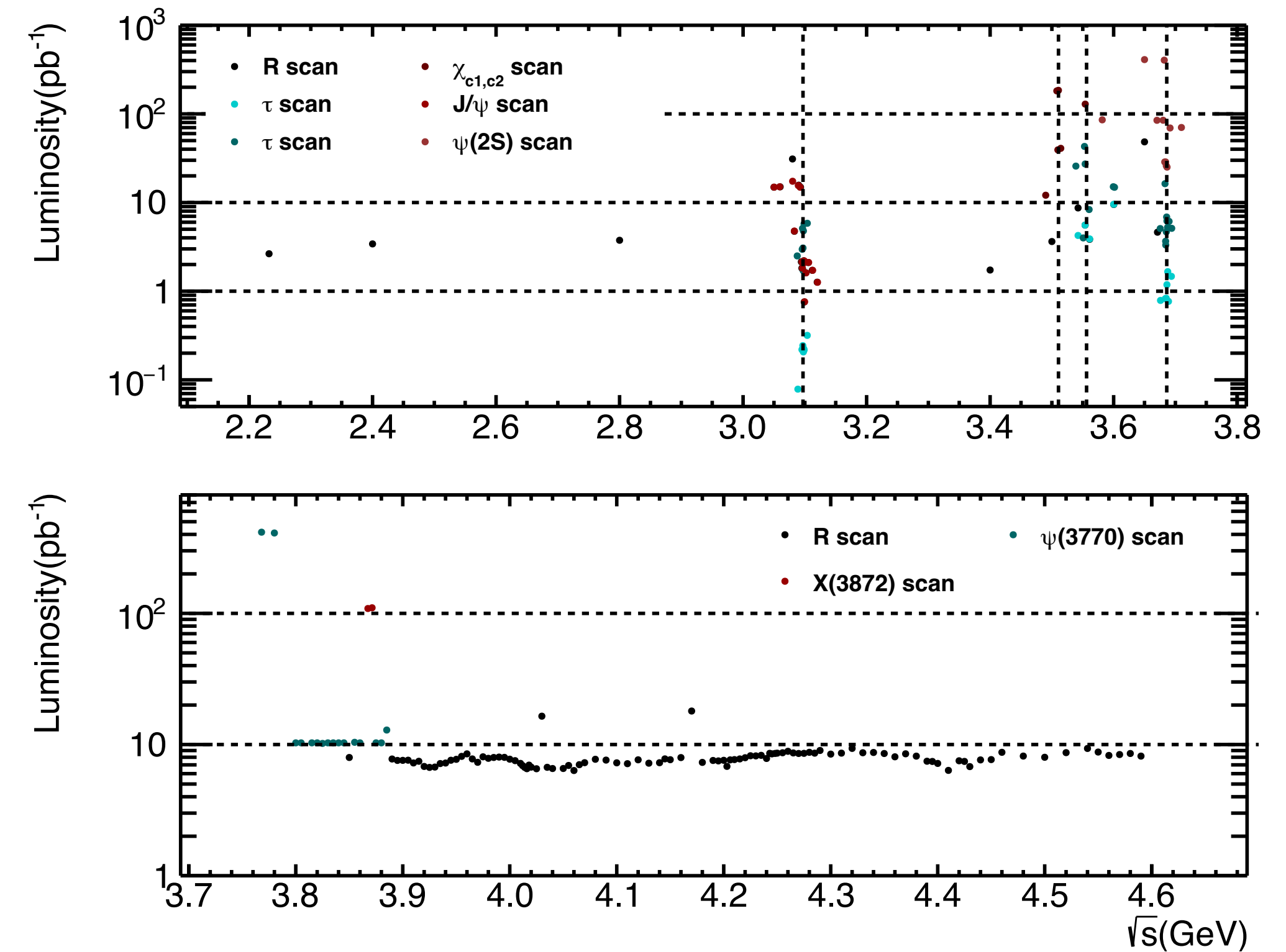
EMC

$\Delta E/E$: at 1GeV
 2.5%
 5.0%
 σ_z : 0.6 cm/ \sqrt{E}

BESIII Data Samples



46 sample, $\sim 22 \text{ fb}^{-1}$



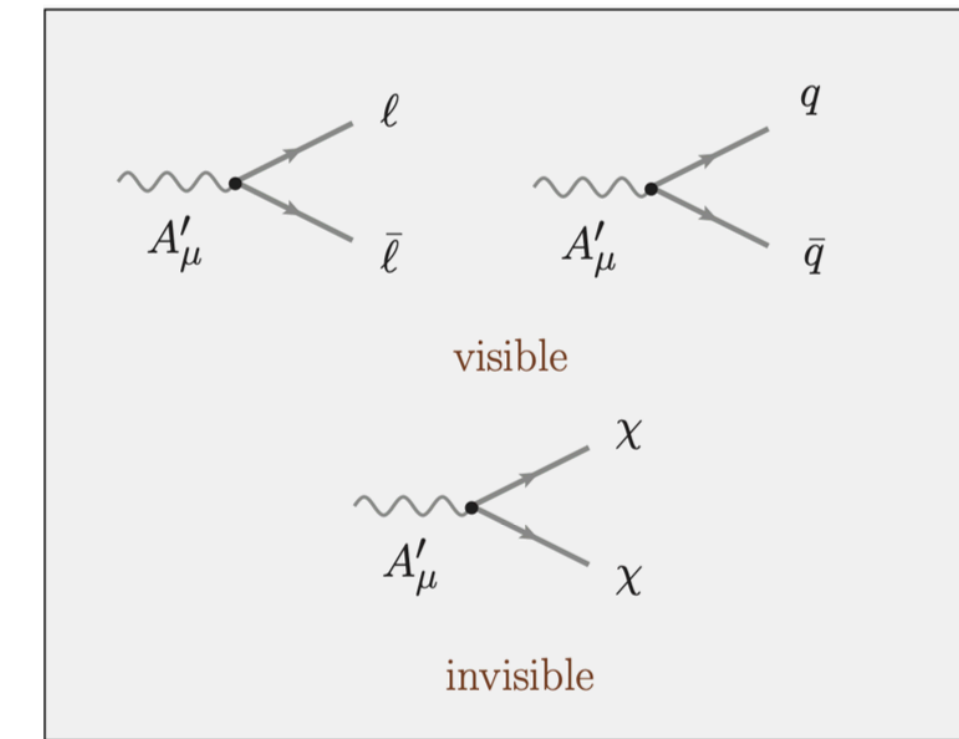
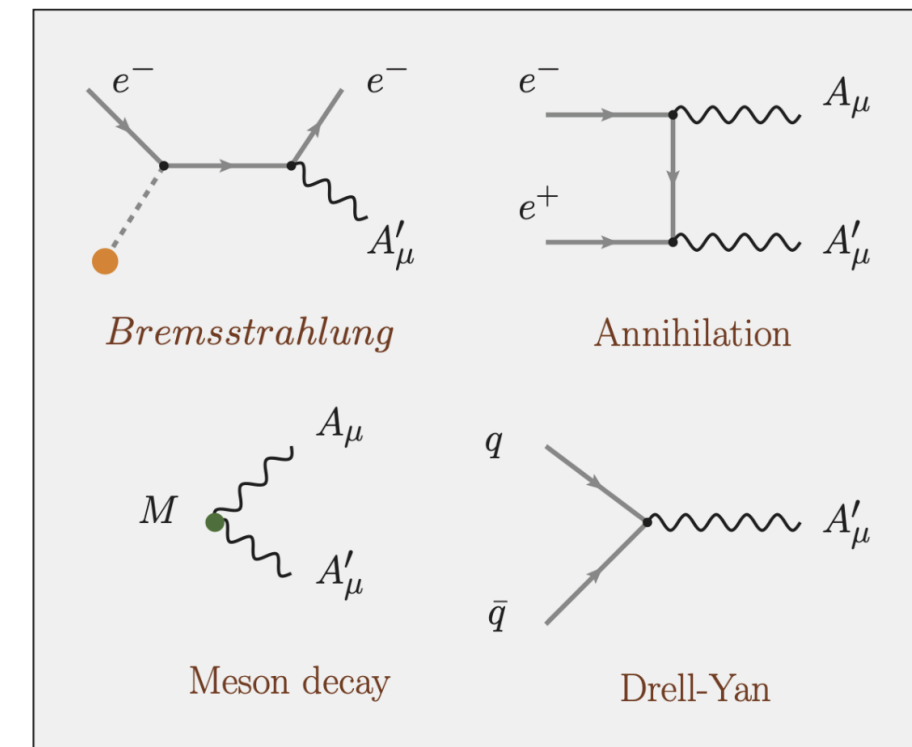
+ Small scan sample, $\sim 3.5 \text{ fb}^{-1}$

+ 10 Billion J/ψ , 2.7 Billion $\psi(3686)$, $20 \text{ fb}^{-1} \psi(3770)$

Dark Photon Studies at BESIII

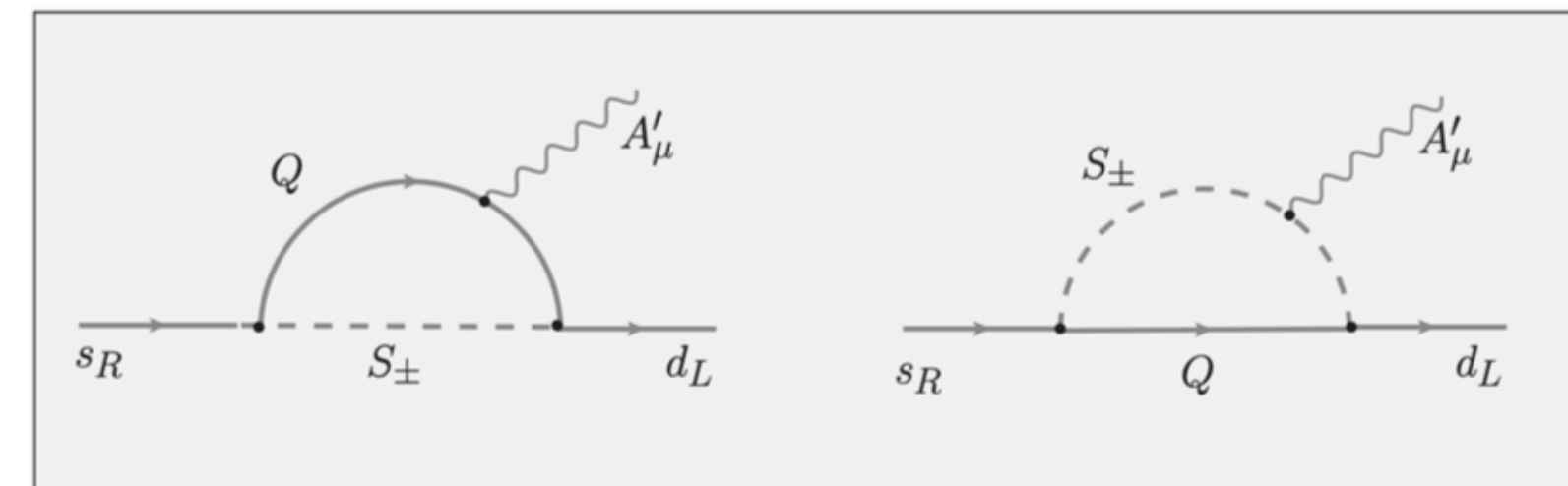
- Massive dark photon

- $e^+e^- \rightarrow \gamma\gamma', \gamma' \rightarrow l^+l^-$ *PLB774, 252 (2017)*
 - $J/\psi \rightarrow \gamma'\eta, \gamma' \rightarrow e^+e^-$ *PRD99, 012006 (2019)*
 - $J/\psi \rightarrow \gamma'\eta', \gamma' \rightarrow e^+e^-$ *PRD99, 012013 (2019)*
 - $e^+e^- \rightarrow \gamma\gamma', \gamma \rightarrow \text{invisible}$ *PLB839, 137785 (2023)*
 - $J/\psi \rightarrow \gamma'\eta', \gamma' \rightarrow \gamma\pi^0$ *PRD102, 052005 (2022), not included in this talk*



- Massless dark photon in hyperon decays

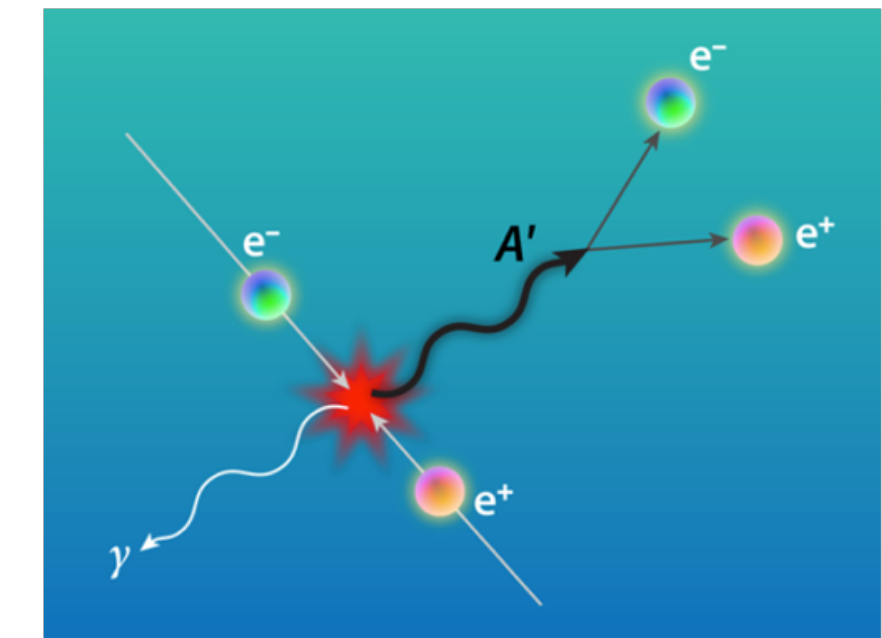
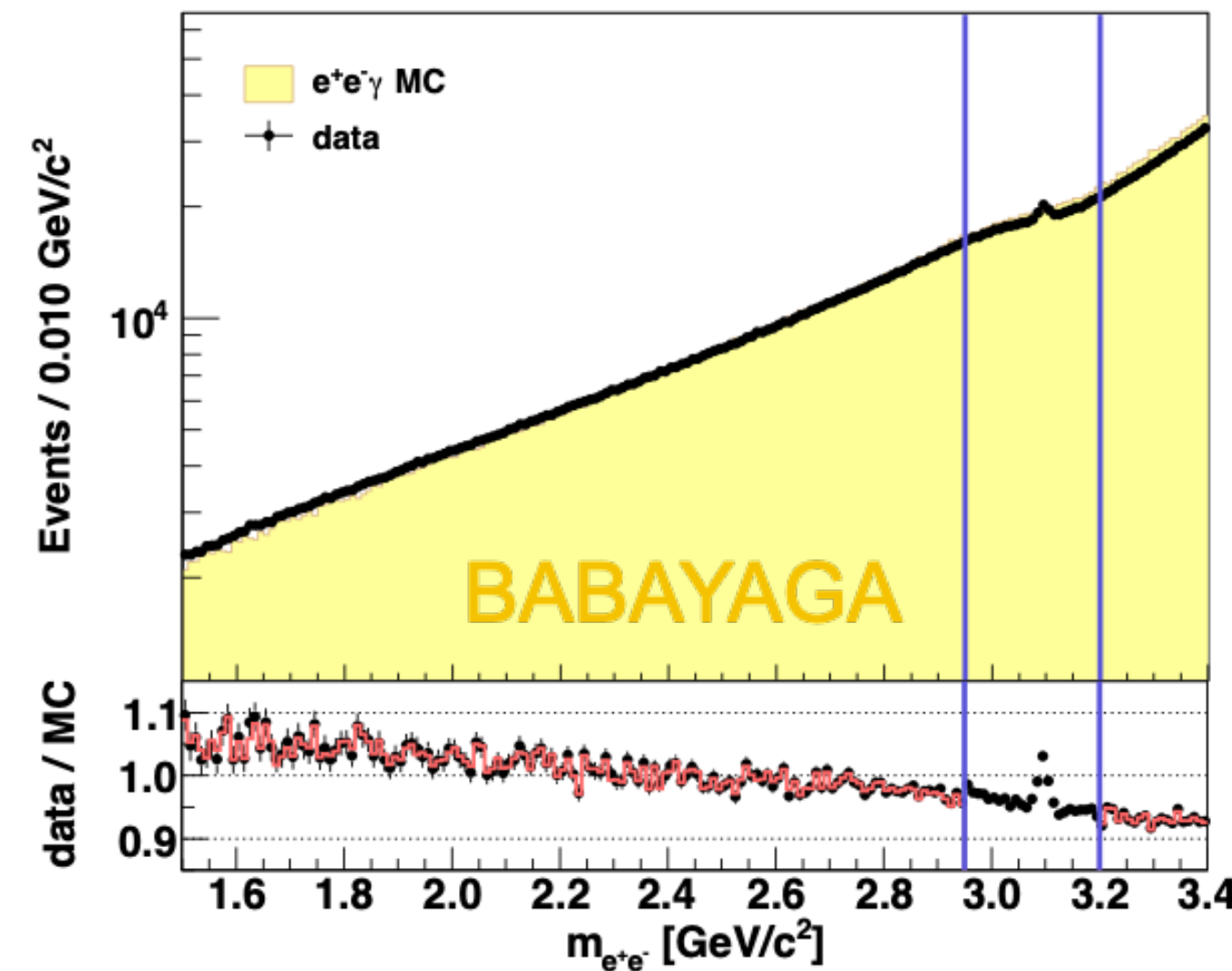
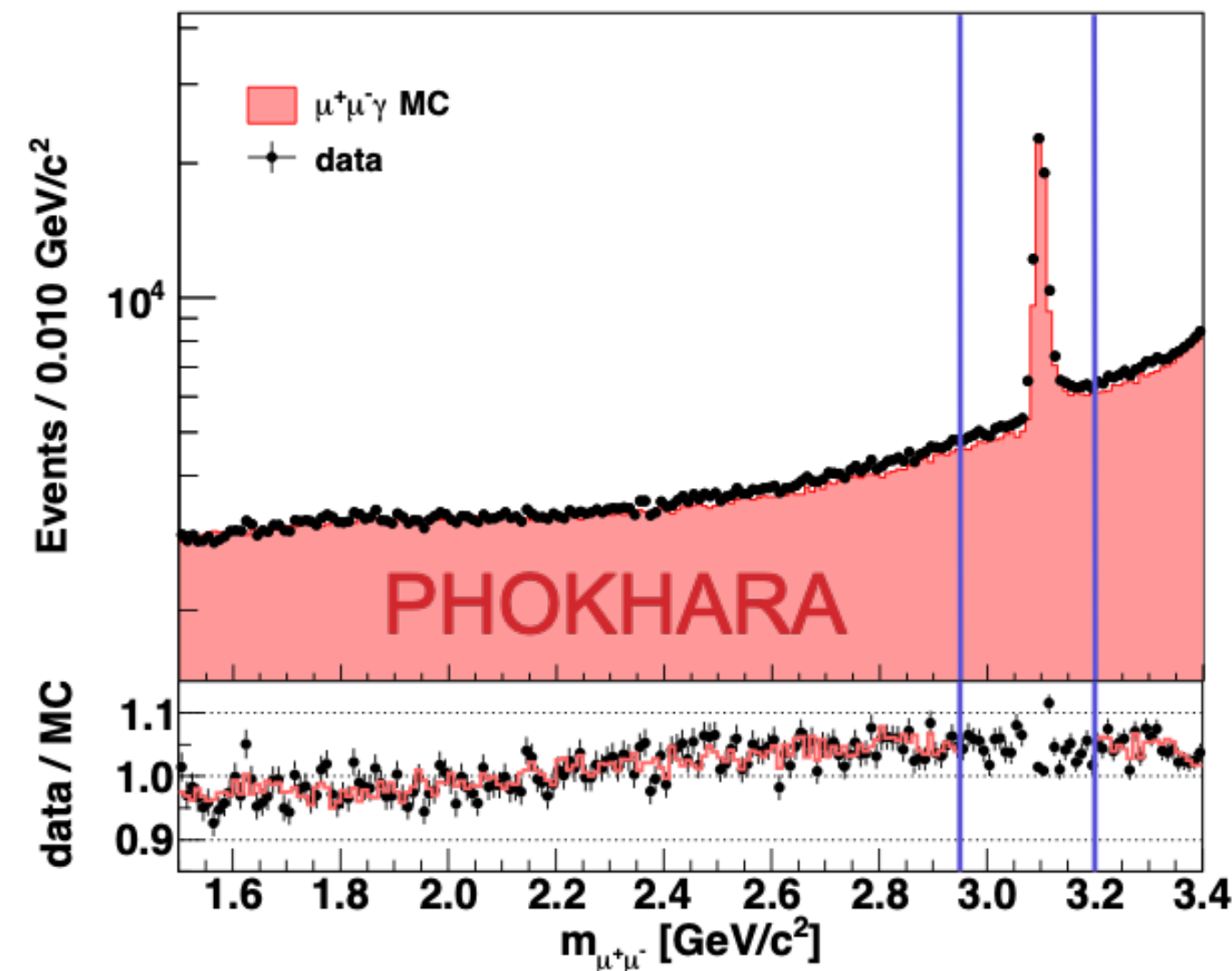
- $\Lambda_c \rightarrow p + \text{invisible}$ *PRD106, 072008 (2002)*
 - $\Sigma \rightarrow p + \text{invisible}$ *PLB852, 138614 (2024)*



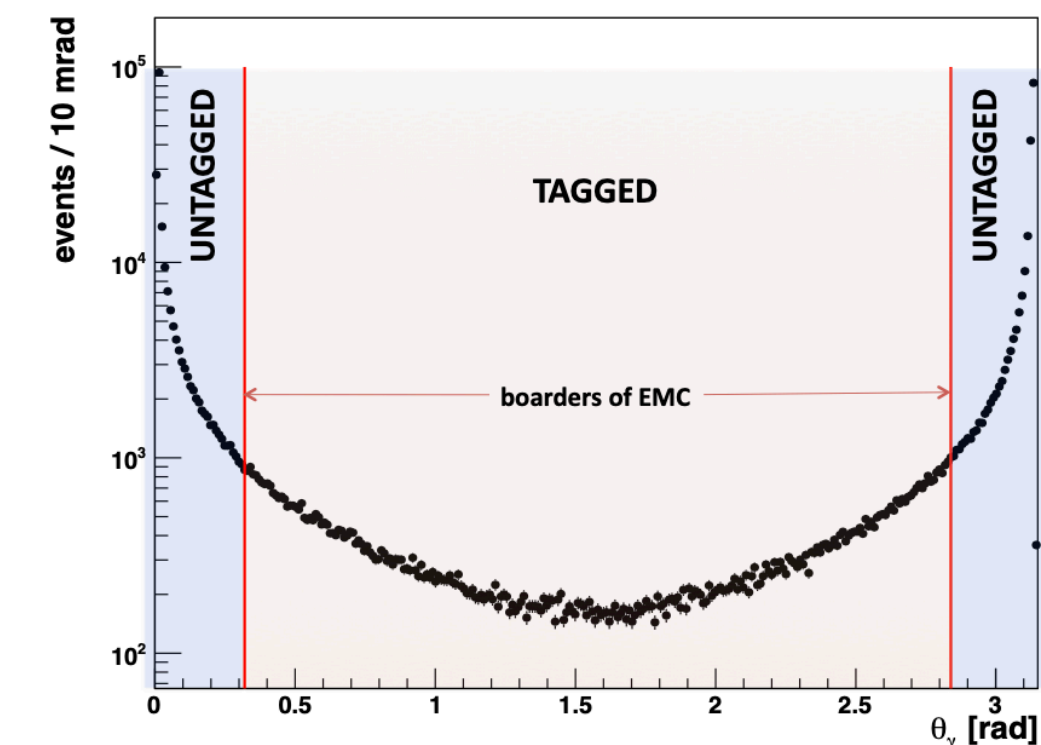
Massive Dark Photon Decays to Lepton Pair

- $e^+e^- \rightarrow \gamma_{\text{ISR}}\gamma', \gamma' \rightarrow l^+l^-$, based on 2.93 fb⁻¹ data sample taken at $\psi(3770)$, untagged method
- Search for narrow structure on top of the continuum QED background

PLB774, 252 (2017)



APS/Alan Stonebreaker



- Cover mass region: 1.5 GeV/c² to 3.4 GeV/c²
 - <1.5 GeV/c²: $\pi^+\pi^-$ background dominant
 - >3.4 GeV/c²: hadronic $q\bar{q}$ background dominant

Massive Dark Photon Decays to Lepton Pair

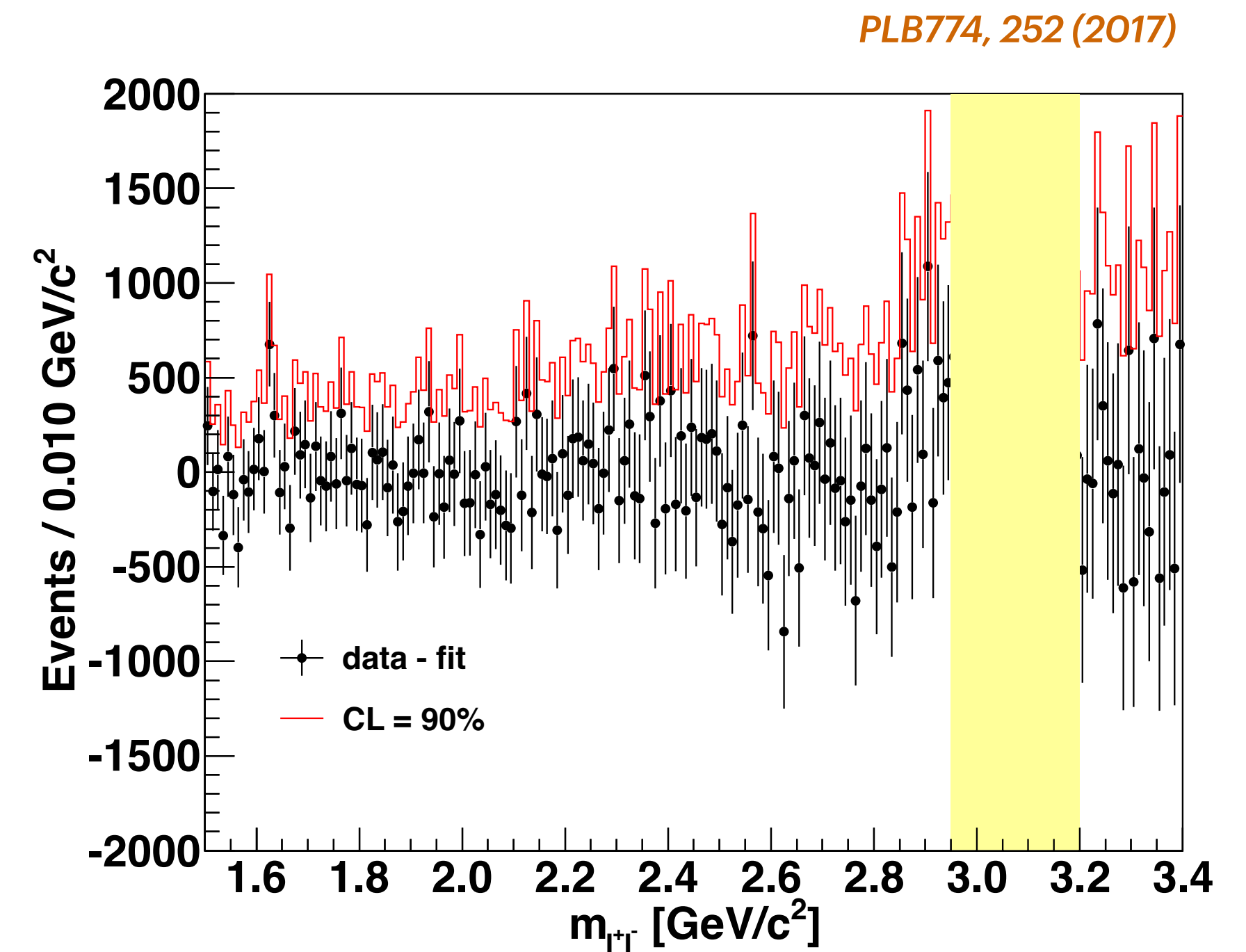
- Number of signal events determined by fitting the $M(l^+l^-)$ distribution, combined statistical significance less than 3σ in the explored mass region
- 90% confidence level limit obtained with profile likelihood approach with systematic uncertainty included
- Kinetic mixing parameter ϵ determined with

$$\frac{\sigma_i(e^+e^- \rightarrow \gamma' \gamma_{\text{ISR}} \rightarrow l^+l^- \gamma_{\text{ISR}})}{\sigma_i(e^+e^- \rightarrow \gamma^* \gamma_{\text{ISR}} \rightarrow l^+l^- \gamma_{\text{ISR}})} = \frac{3\pi}{2N_f^{l^+l^-}} \cdot \frac{\epsilon^2}{\alpha} \cdot \frac{m_{\gamma'}}{\delta_m^{l^+l^-}}$$

Number of dark photon events (points to numerator)
 Number of QED events from annihilation process (points to denominator)
 Ratio of possible decay channels of the dark photon and the phase space (points to $N_f^{l^+l^-}$)
 Mixing coefficient (points to ϵ^2)
 Dark photon mass (points to $m_{\gamma'}$)
 Bin width (points to $\delta_m^{l^+l^-}$)
 Fine structure constant (points to α)

$$N_f^{l^+l^-} = \frac{\Gamma_{\text{tot}}}{\Gamma(\gamma' \rightarrow l^+l^-)} = 1 + \frac{\Gamma_{\mu\mu}}{\Gamma_{ee} + \Gamma_{\mu\mu}} \cdot (1 + R(\sqrt{s}))$$

J. D. Bjorken, R. Essig, P. Schuster, N. Toro, PRD 80, 075018 (2009)



Massive Dark Photon Decays to Lepton Pair

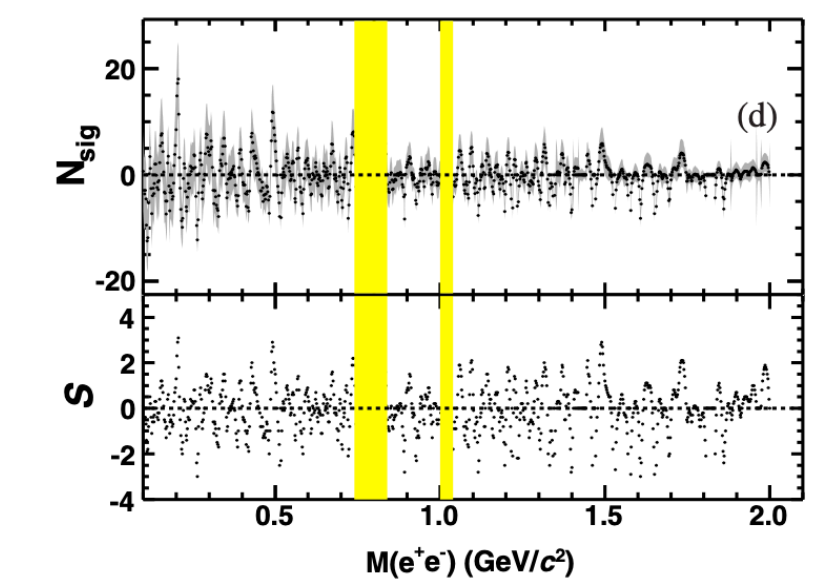
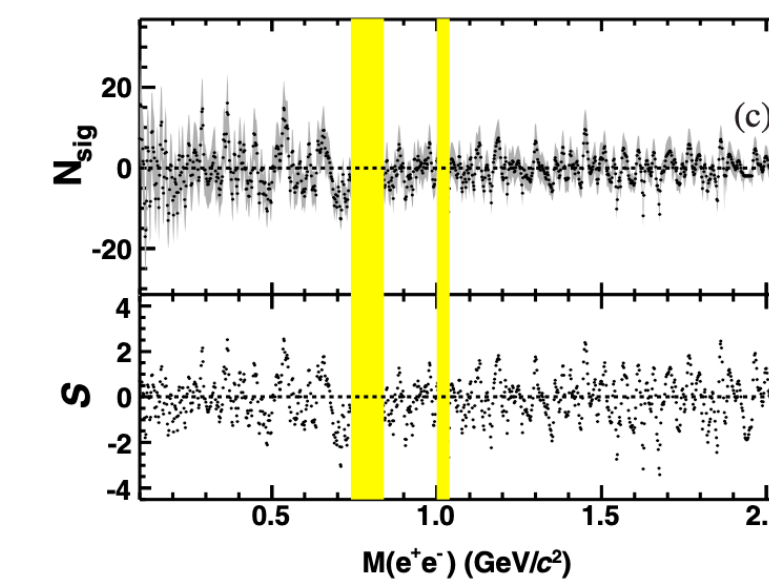
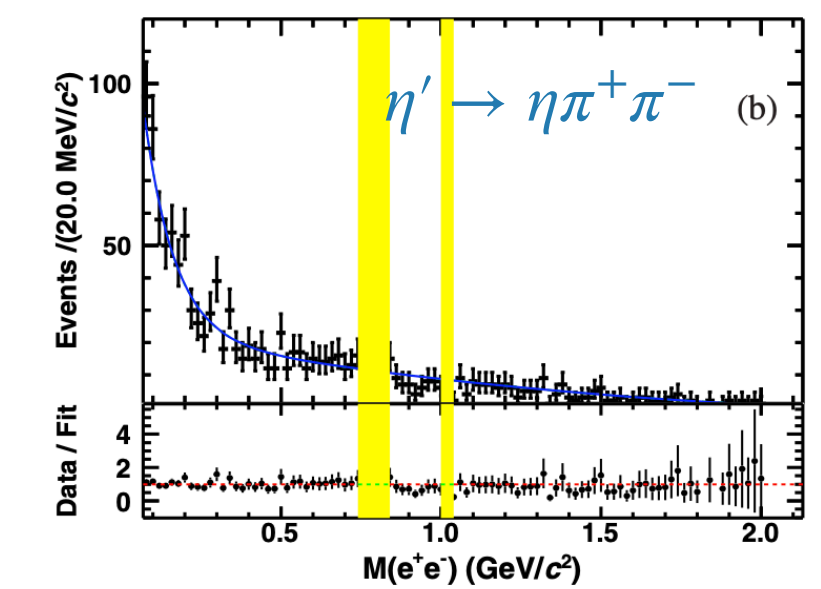
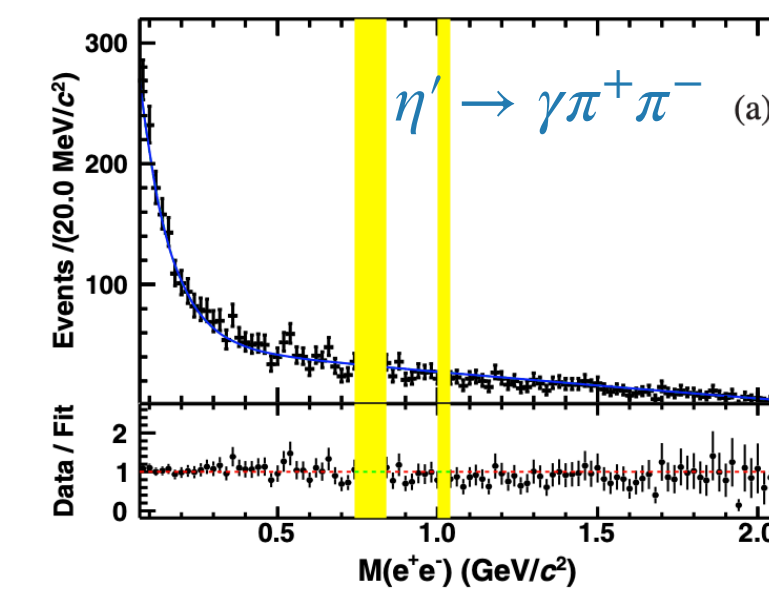
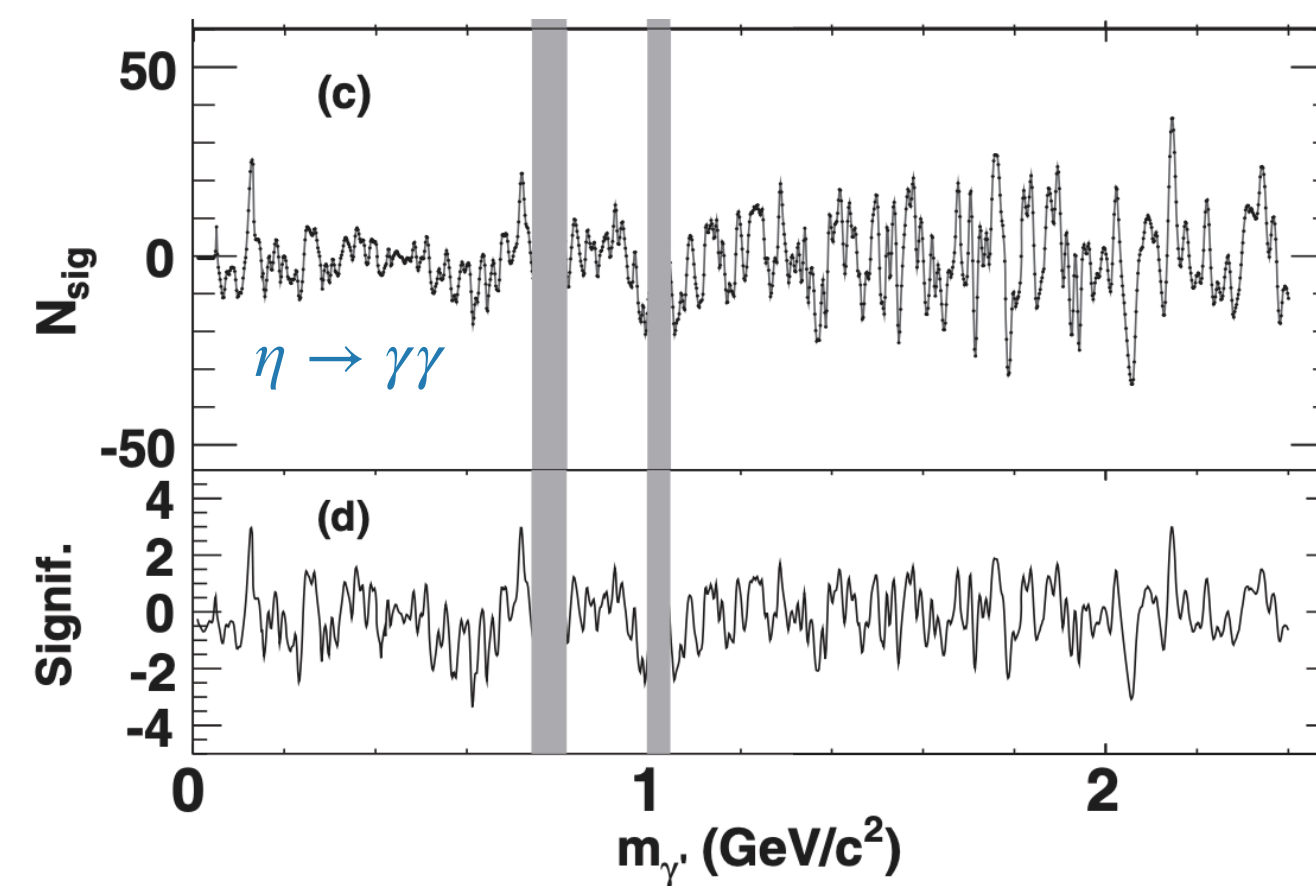
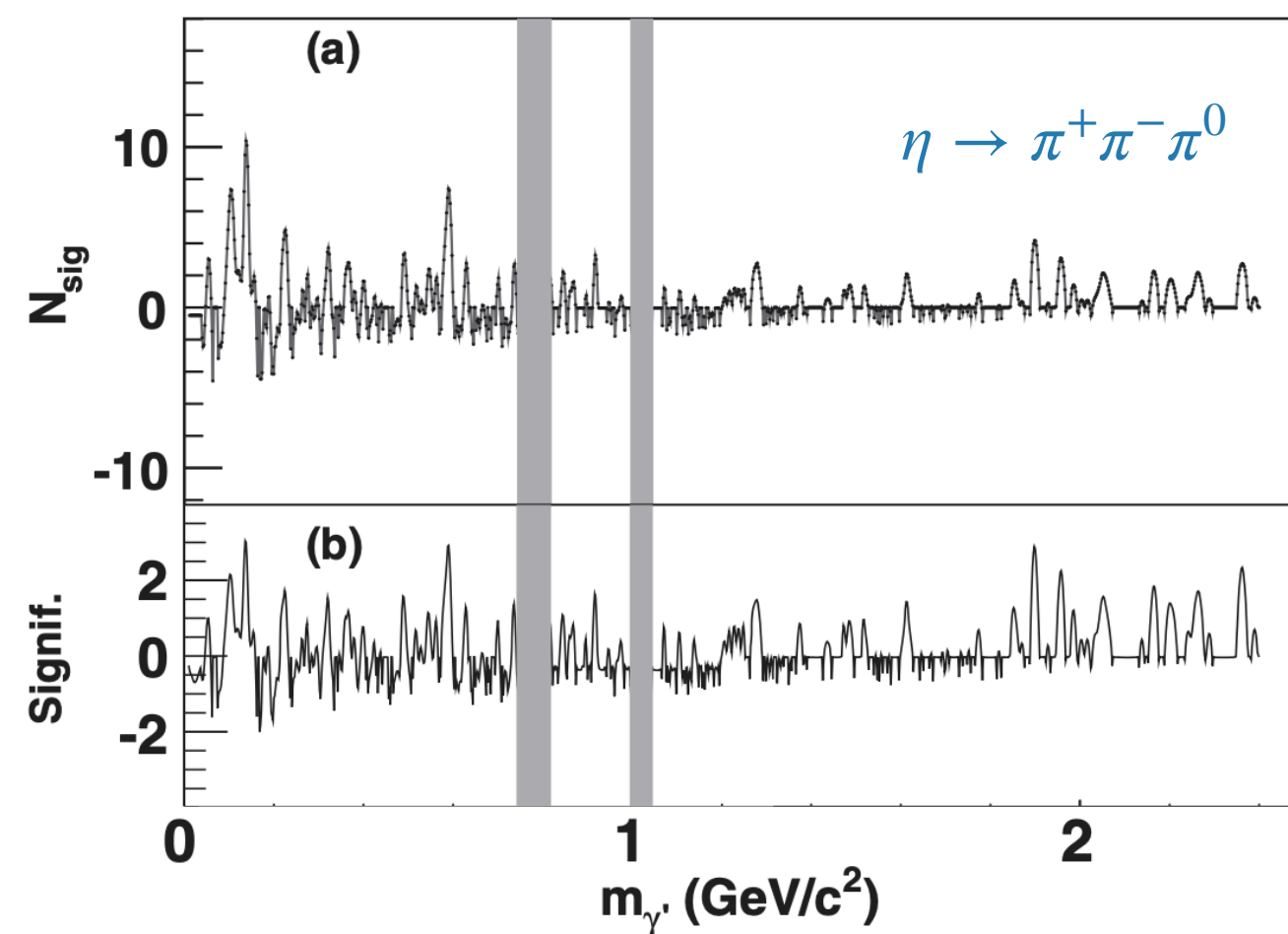
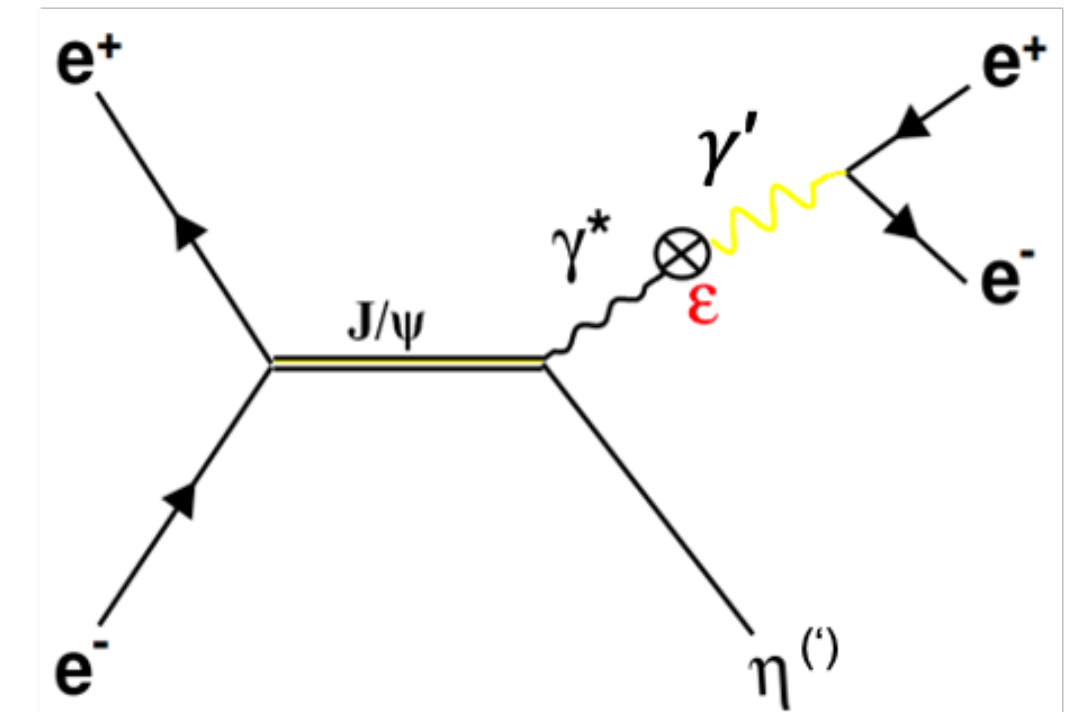
- Electromagnetic (EM) Dalitz decay, based on $(1310.67 \pm 7.0) \times 10^6$ J/ψ events

- $J/\psi \rightarrow \gamma' \eta \rightarrow e^+ e^- \eta, \eta \rightarrow \gamma \gamma / \pi^+ \pi^- \pi^0$ PRD99, 012006 (2019)

- $J/\psi \rightarrow \gamma' \eta' \rightarrow e^+ e^- \eta', \eta' \rightarrow \gamma \pi^+ \pi^- / \eta \pi^+ \pi^-$ PRD99, 012013 (2019)

- Fit to $M(e^+ e^-)$ distribution in $2 \text{ MeV}/c^2$ step, exclude ρ, ω, ϕ mass region

- $M(\gamma') \in [0.01, 2.40] \text{ GeV}/c^2$ in η mode, $[0.07, 2.13] \text{ GeV}/c^2$ in η' mode



Massive Dark Photon Decays to Lepton Pair

Branching fraction of $J/\psi \rightarrow \gamma' \eta^{(\prime)}$

Dark photon mass

Mixing coefficient

$$\frac{B(J/\psi \rightarrow \gamma' \eta^{(\prime)})}{B(J/\psi \rightarrow \gamma \eta^{(\prime)})} = \varepsilon^2 |F_{J/\psi \eta^{(\prime)}}(m_{\gamma'}^2)|^2 \cdot \frac{\Lambda^{3/2}(m_{J/\psi}^2, m_{\eta^{(\prime)}}^2, m_{\gamma'}^2)}{\Lambda^{3/2}(m_{J/\psi}^2, m_{\eta^{(\prime)}}^2, 0)}$$

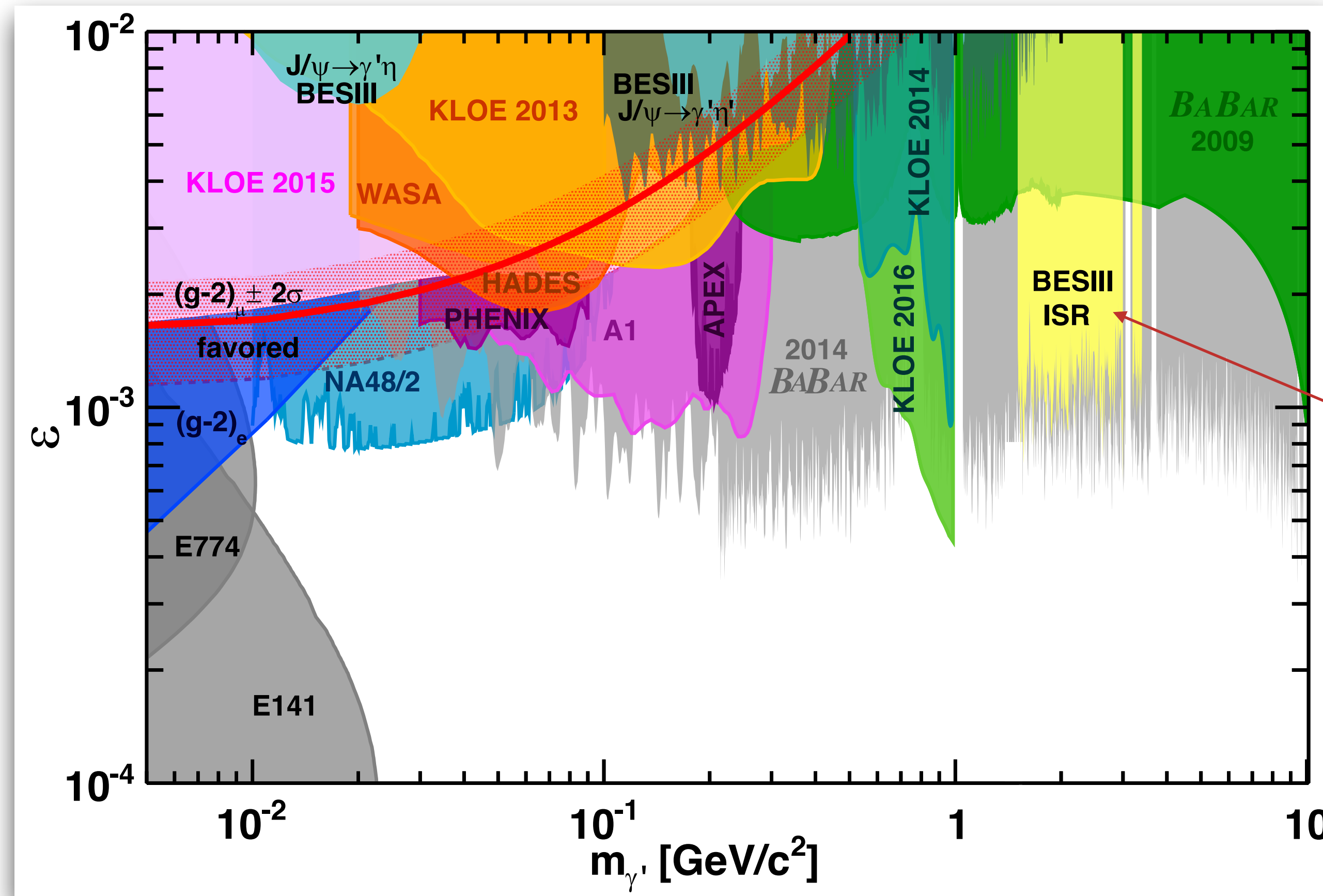
Branching fraction of $J/\psi \rightarrow \gamma \eta^{(\prime)}$

$$\Lambda^{3/2}(m_1^2, m_2^2, m_3^2) = \left(1 + \frac{m_3^2}{m_1^2 - m_2^2}\right)^2 - \frac{4m_1^2 m_3^2}{(m_1^2 - m_2^2)^2}$$

Form factor for $J/\psi \rightarrow \gamma^* \eta^{(\prime)}$ transition
evaluated at γ' mass

M. Reece and L. T. Wang JHEP07, 051 (2009)

Kinetic Mixing Coefficient ϵ

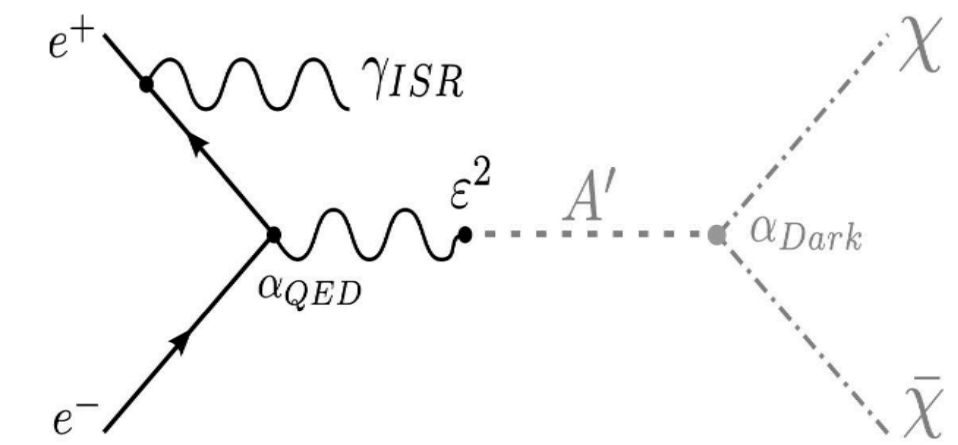


$\epsilon < (0.6 - 2.0) \times 10^{-3}$
 for $m_{\gamma'} \in [1.5, 3.4] \text{ GeV}/c^2$
 Results compatible with
 BaBar measurement based
 on $514 \text{ fb}^{-1} \Upsilon$ data sample

Larger $\psi(3770)$ ($2.93 \Rightarrow 20 \text{ fb}^{-1}$) and J/ψ (1 billion \Rightarrow 10 billion) data sample already collected

Massive Dark Photon Decays to Invisible

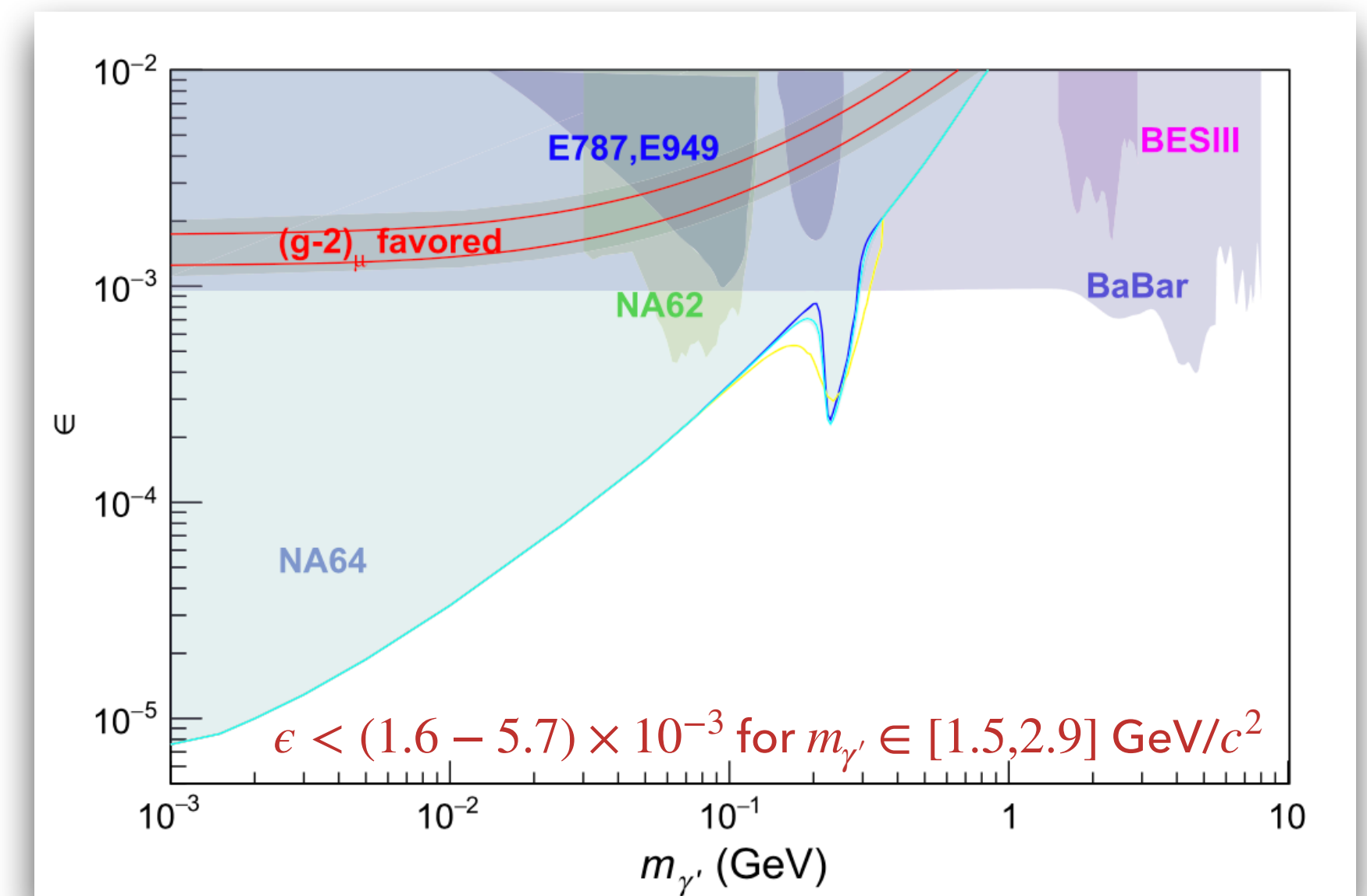
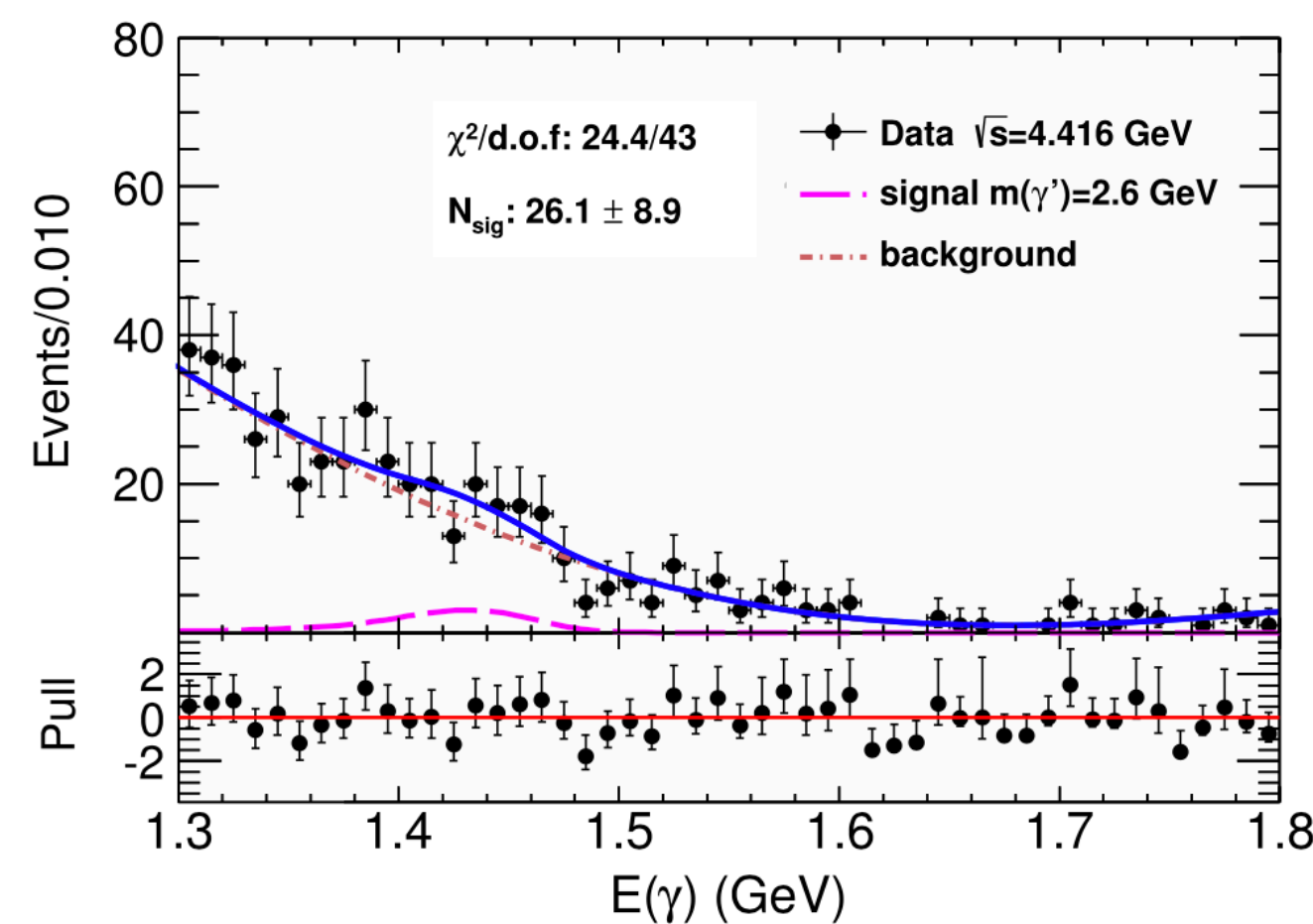
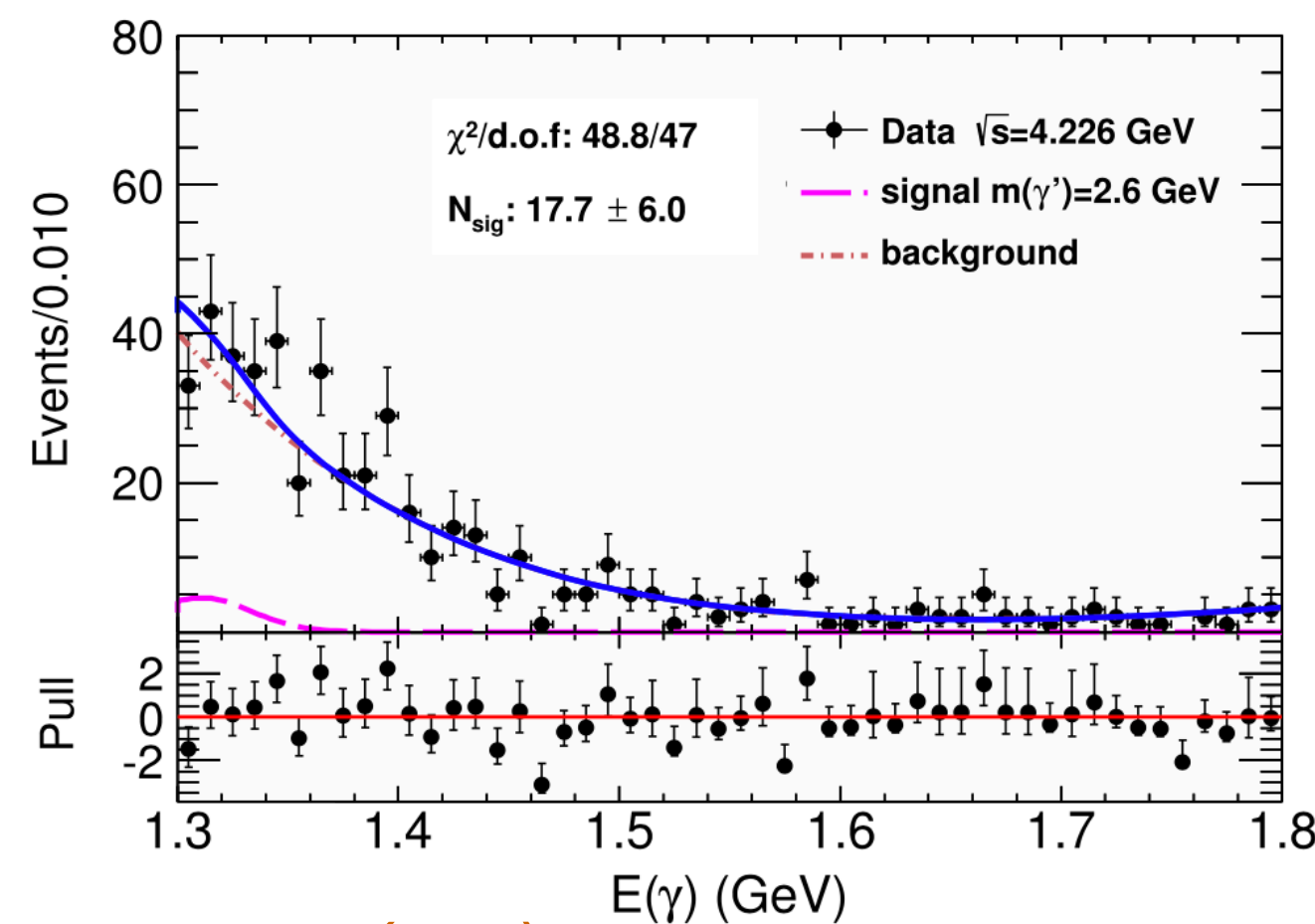
- Based on 14.9 fb⁻¹ data sample taken at $\sqrt{s} = 4.13 - 4.60$ GeV
- Dark photon candidate would be signified by the presence of a the monochromatic photon $E_\gamma = \frac{s - m_{\gamma'}^2}{2\sqrt{s}}$



- $E(\gamma) < 1.3$ GeV discarded due to low trigger efficiency
- Maximum local significance: 3.1σ at $m_{\gamma'} = 2.6$ GeV/c²

$$\sigma(e^+e^- \rightarrow \gamma\gamma') = \frac{2\pi\alpha^2}{s} \epsilon^2 \left(1 + \frac{2x}{(1-x)^2}\right) \log \frac{(1 + \cos\theta_c)^2}{(1 - \cos\theta_c)^2} - 2 \cos\theta_c$$

R. Essig, P. Schuster, N. Toro, PRD 80, 015003 (2009)



PLB839, 137785 (2023)

Massless Dark Photon in Λ_c Decays

- The massless dark photon is a gauge boson associated with a new unbroken $U(1)_d$ symmetry

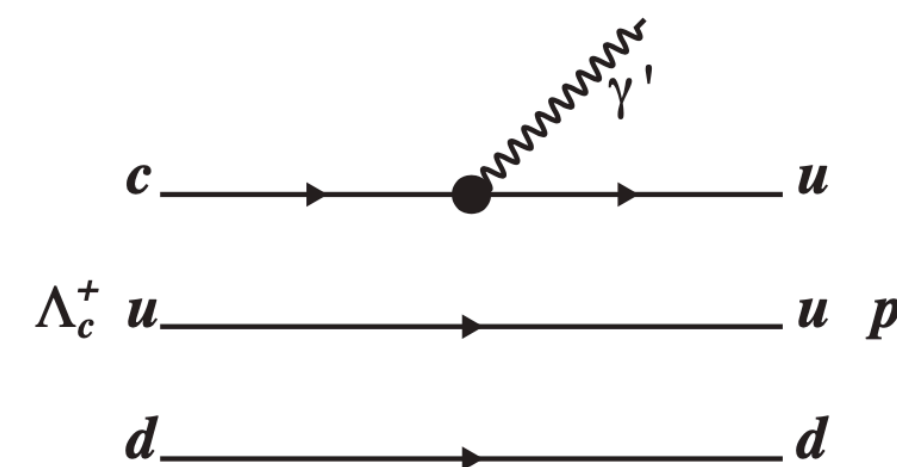
B. Holdom, PLB 166,196 (1986)

B. A. Dobrescu, PRL94, 151802 (2005)

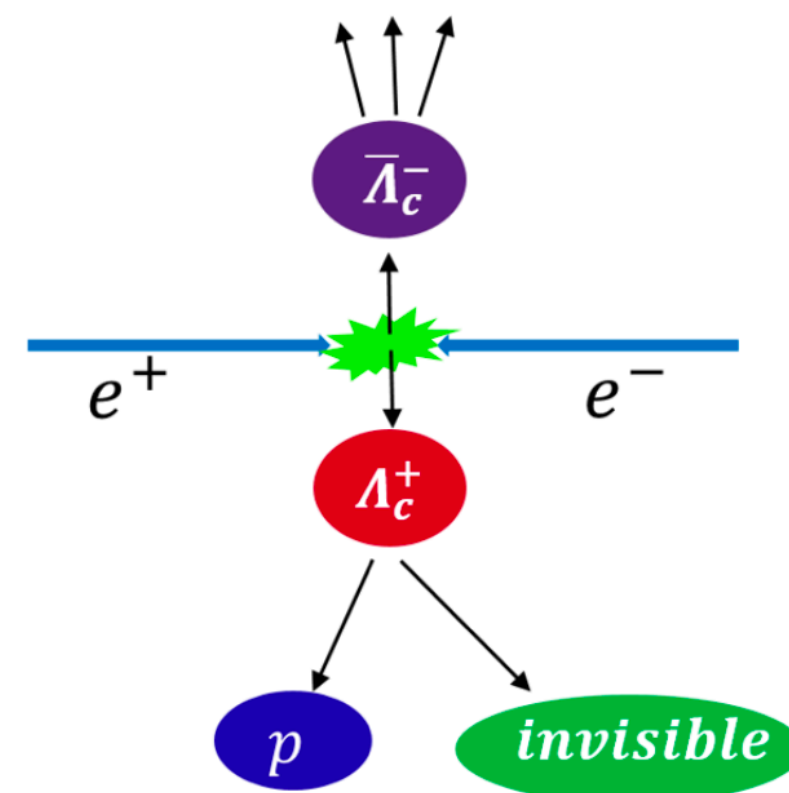
- Do not interact with SM particle directly, but can exert influence on the SM via higher-dimensional operators generated by loop diagrams

- May enhance the branching fractions of flavor changing neutral current (FCNC) decays

- Based on 4.5 fb^{-1} data sample at $\sqrt{s} = 4.6 - 4.7 \text{ GeV}$

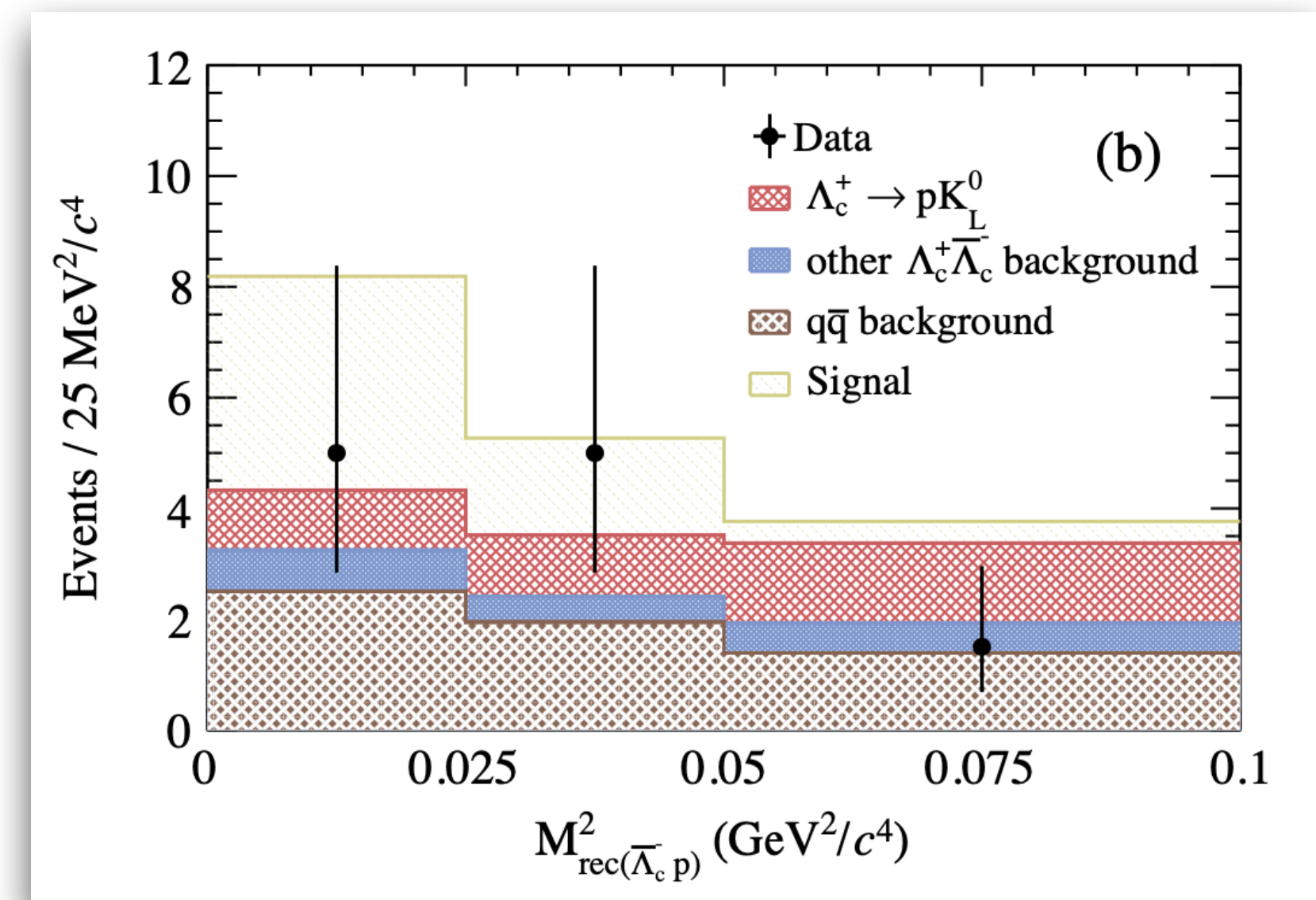


10 hadronic decay modes



105244 ± 384 tagged $\Lambda_c \bar{\Lambda}_c$ pair

$$B(\Lambda_c^+ \rightarrow p\gamma') = \frac{N_{\text{obs}} - N_{\text{bkg}}}{\sum_i N_i^{\text{ST}} \cdot (\epsilon_i^{\text{DT}}/\epsilon_i^{\text{ST}})}$$



$$N_{\text{obs}} = 13, N_{\text{bkg}} = 14.6 \pm 1.5$$

$$B(\Lambda_c^+ \rightarrow p\gamma') < 8.0 \times 10^{-5} \text{ at } 90\% \text{ C.L.}$$

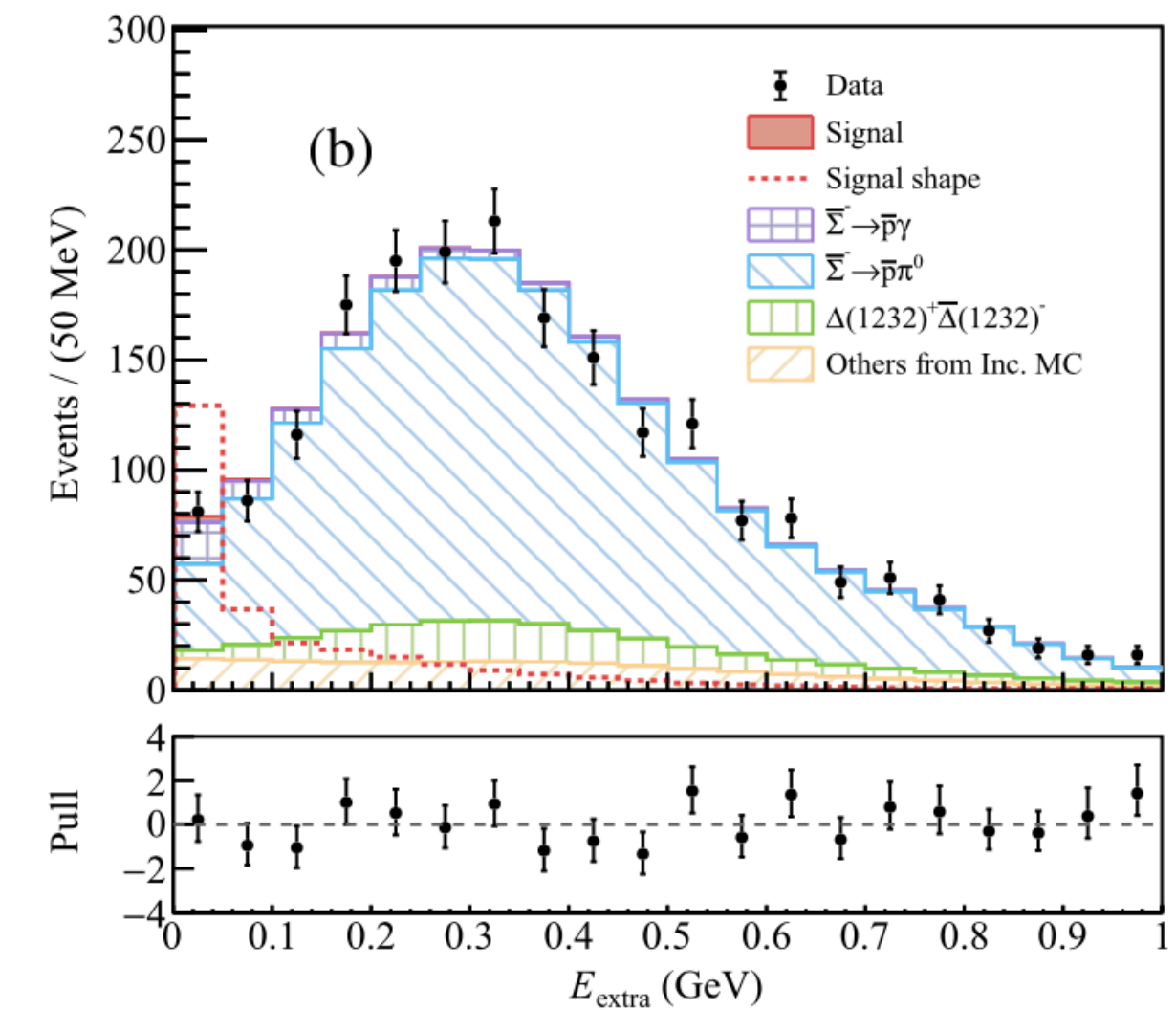
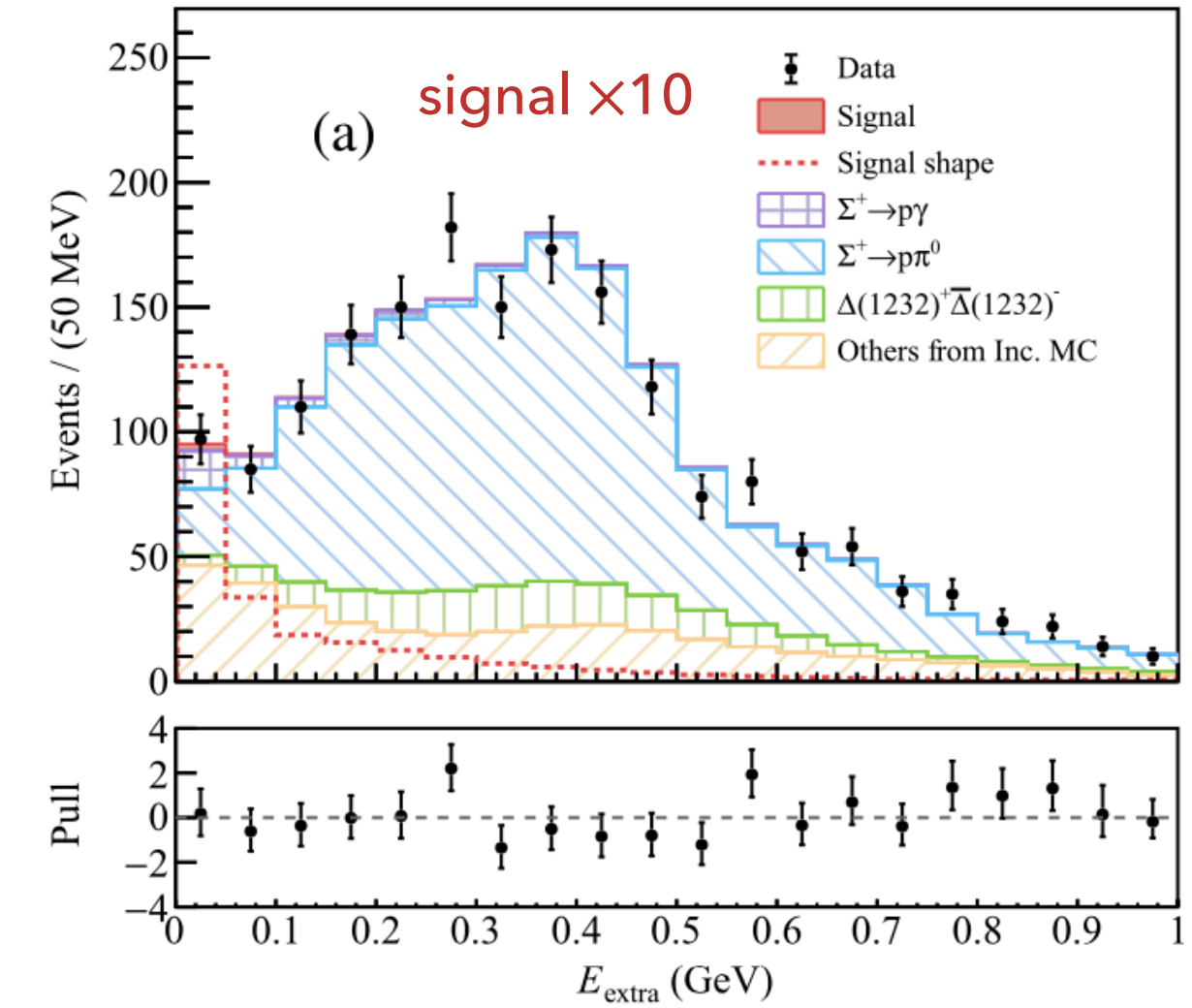
Massless Dark Photon in Σ Decays

- Based on 10 billion J/ψ events, about 4 million tagged Σ^\pm using $p\pi^0$ mode
- In search for massless dark photon, QCD axion, ...

- $$B_{\text{sig}} = \frac{N_{\text{DT}}^{\text{obs}}/\epsilon_{\text{DT}}}{N_{\text{ST}}^{\text{obs}}/\epsilon_{\text{ST}}}$$

- DT signal extracted from the missing energy, defined as $E_{\text{extra}} = E_{\text{extra}}^{\text{DT}\pi^0/\gamma} + E_{\text{extra}}^{\text{other}}$
 - $E_{\text{extra}}^{\text{DT}\pi^0/\gamma}$ expected to be 0 for signal events
 - $E_{\text{extra}}^{\text{other}}$: 93% contribution from interaction between the \bar{p} and the detector, determined using data-driven approach
- $B(\Sigma^+ \rightarrow p + \text{invisible}) < 3.2 \times 10^{-5}$

PLB852, 138614 (2024)



Summary

- Electron-Positron colliders provide **clean environment** for the study of particles beyond SM, in both **visible and invisible decays**
- BESIII works on tau-charm energy region, can probe **lower mass region**, and can provide **better mass resolution**
- Massive and massless dark photon has been studied in visible or invisible final states, based on part of the data samples
- More results to be expected in the near future, with larger data samples

Thank You!