



Dark Photon Dark Matter Detection with Radio Telescopes

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Based on 2010.15836 [PRL 2021], 2207.05767 [PRL 2023],
2301.03622 [NC 2024]

第三届地下和空间粒子物理与宇宙物理前沿问题研讨会

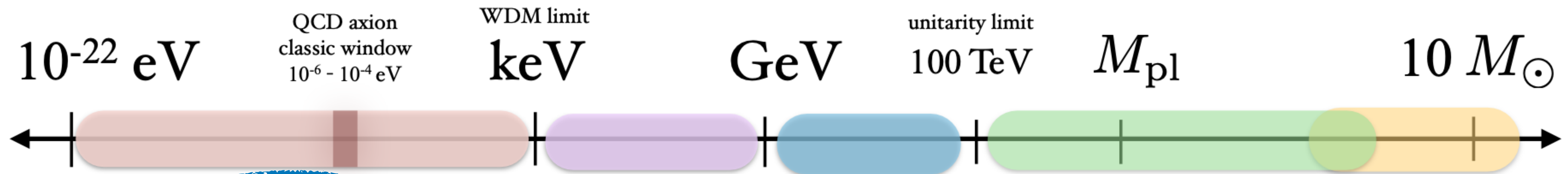
2024-05-09

What is the nature of dark matter?

Unknown matter and energy $\sim 95\%$

The dark matter candidate models

1904.07915, TASI lecture



“Ultralight” DM
non-thermal
bosonic fields

“Light” DM
dark sectors
sterile ν
can be thermal

WIMP

Composite DM
(Q-balls, nuggets, etc)

**Primordial
black holes**

↕
Today's focus


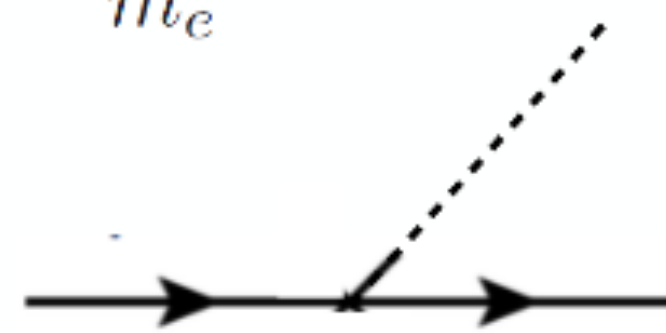
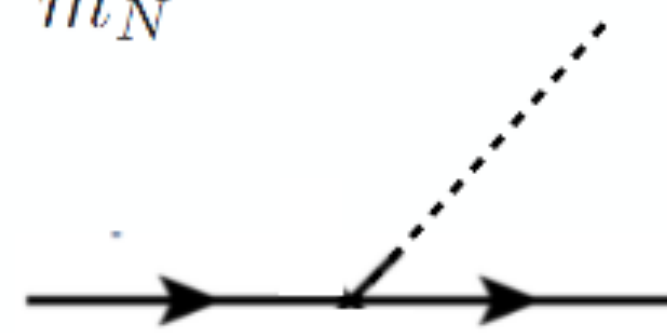
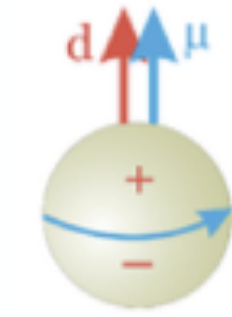


HEP at a cross-road: explore all directions!

Ultralight Bosonic Dark Matter

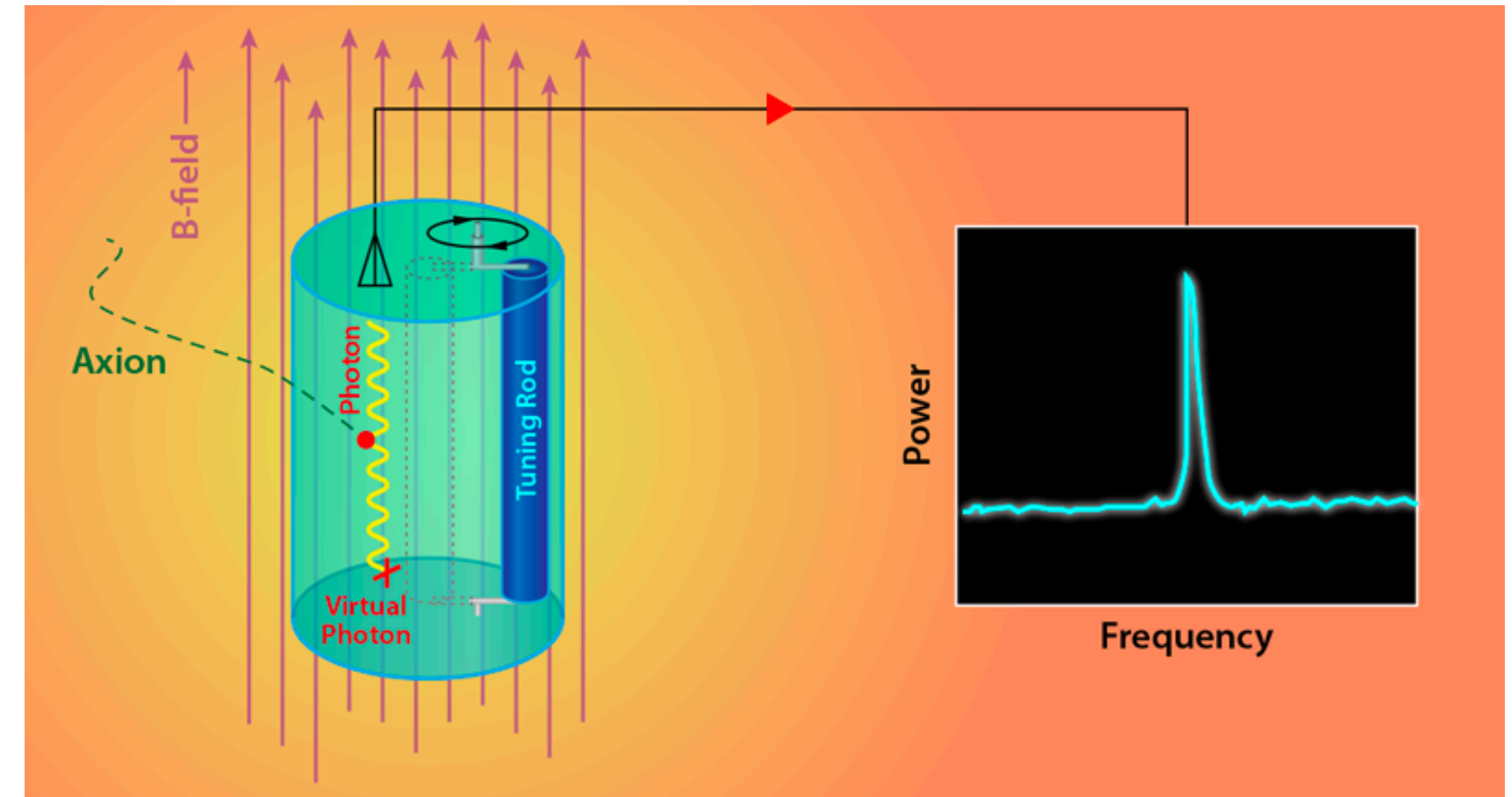
- Ultralight: $m \lesssim \text{keV}$, ultralight due to shift symmetry (pseudo-Nambu Goldstone, e.g. Axion)
- Bosonic: Pauli-exclusion for fermionic DM
- Exits as classical fields ($m \lesssim \mathcal{O}(1) \text{ eV}$)
- Typical models:
 - Pseudo-scalar: Axion, Axion-like Particle
 - Dark Scalar: dilaton-like coupling
 - Vector: Kinetic Mixing Dark Photon, $U(1)_{B-L}$ dark photon etc

超轻暗物质：赝标量轴子

photon coupling	electron coupling	nucleon coupling	Neutron electric dipole
$-\frac{g_{a\gamma}}{4} F_{\mu\nu} \tilde{F}^{\mu\nu} a$ 	$\frac{g_{ae}}{m_e} [\bar{e} \gamma^\mu \gamma^5 e] \partial_\mu a$ 	$\frac{g_{aN}}{m_N} [\bar{N} \gamma^\mu \gamma^5 N] \partial_\mu a$ 	$\propto \frac{1}{m_n} [F_{\mu\nu} \bar{n} \sigma^{\mu\nu} \gamma^5 n] \frac{A}{f_A}$ 

$$\mathcal{L}_{\text{ALP}} = g_{ag} \frac{a}{f_a} G\tilde{G} + g_{a\gamma} \frac{a}{f_a} F\tilde{F} + g_{af} \frac{\partial_\mu a}{2f_a} \bar{f} \gamma^\mu \gamma^5 f$$

$$g_{a\gamma\gamma} a F_{\mu\nu} \epsilon^{\mu\nu\alpha\beta} F_{\alpha\beta} \sim g_{a\gamma\gamma} a \vec{E} \cdot \vec{B}$$



矢量型超轻暗物质：暗光子

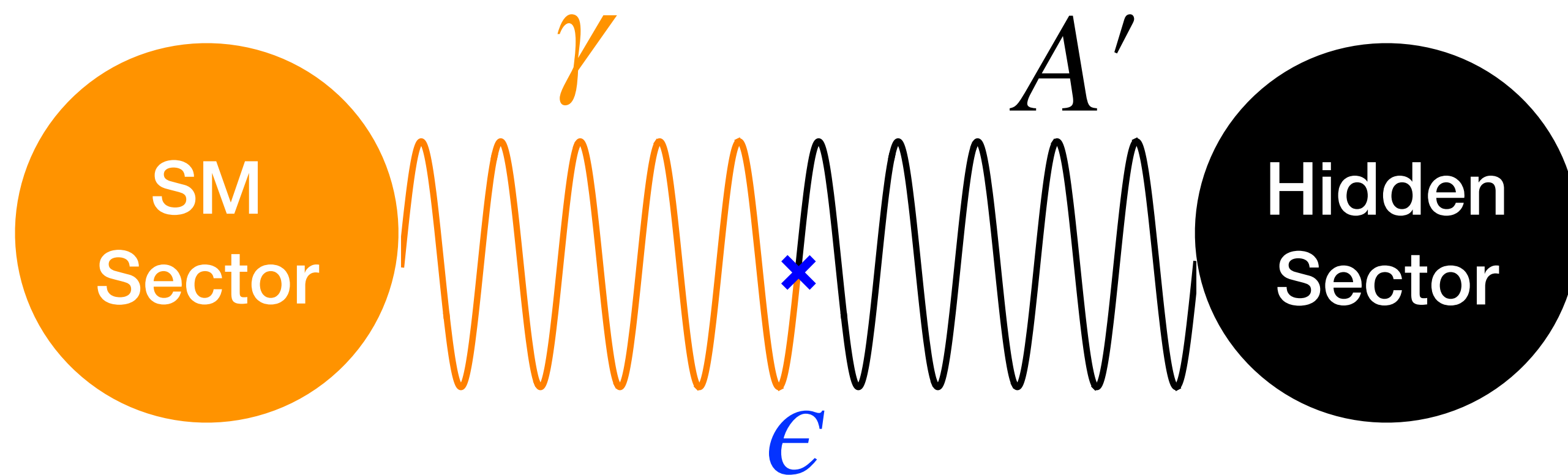
- Maxwell Equations: $\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + 0 \times A^\mu A_\mu - eA_\mu j_{\text{em}}^\mu$

- Extra U(1) extension of Maxwell Equations

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + 0 \times A^\mu A_\mu - eA_\mu j_{\text{em}}^\mu - \frac{1}{4}F'_{\mu\nu}F'^{\mu\nu} + \frac{1}{2}m_{A'}^2 A'^\mu A'_\mu - \frac{1}{2}\epsilon F'_{\mu\nu}F^{\mu\nu}$$

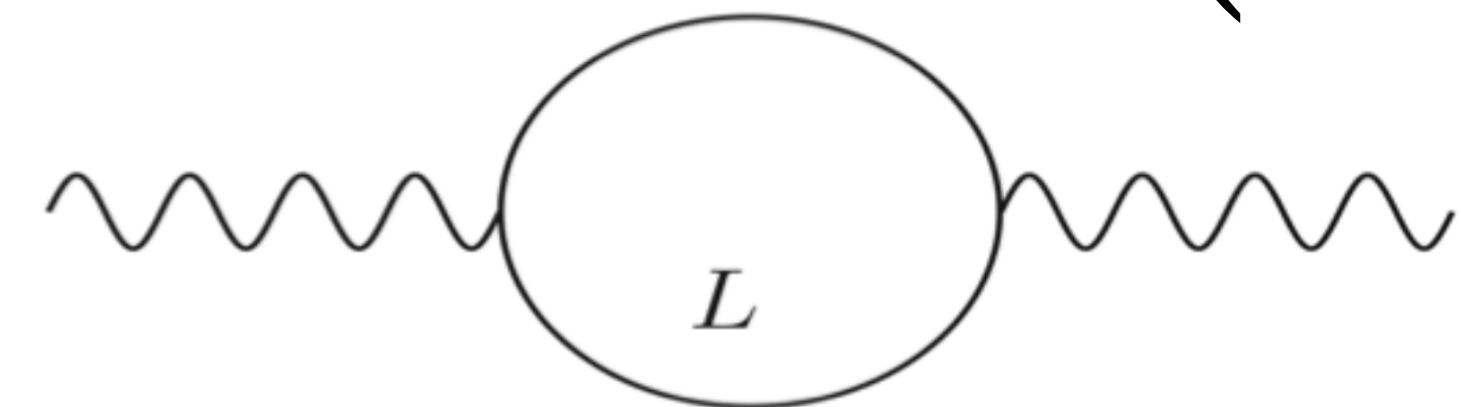
A proper low energy model from UV physics
Log dependence of UV scale

$$\epsilon \sim -\frac{gg'}{16\pi^2} \log\left(\frac{m_L^2}{\mu^2}\right)$$



Standard Model Sector U(1)

γ

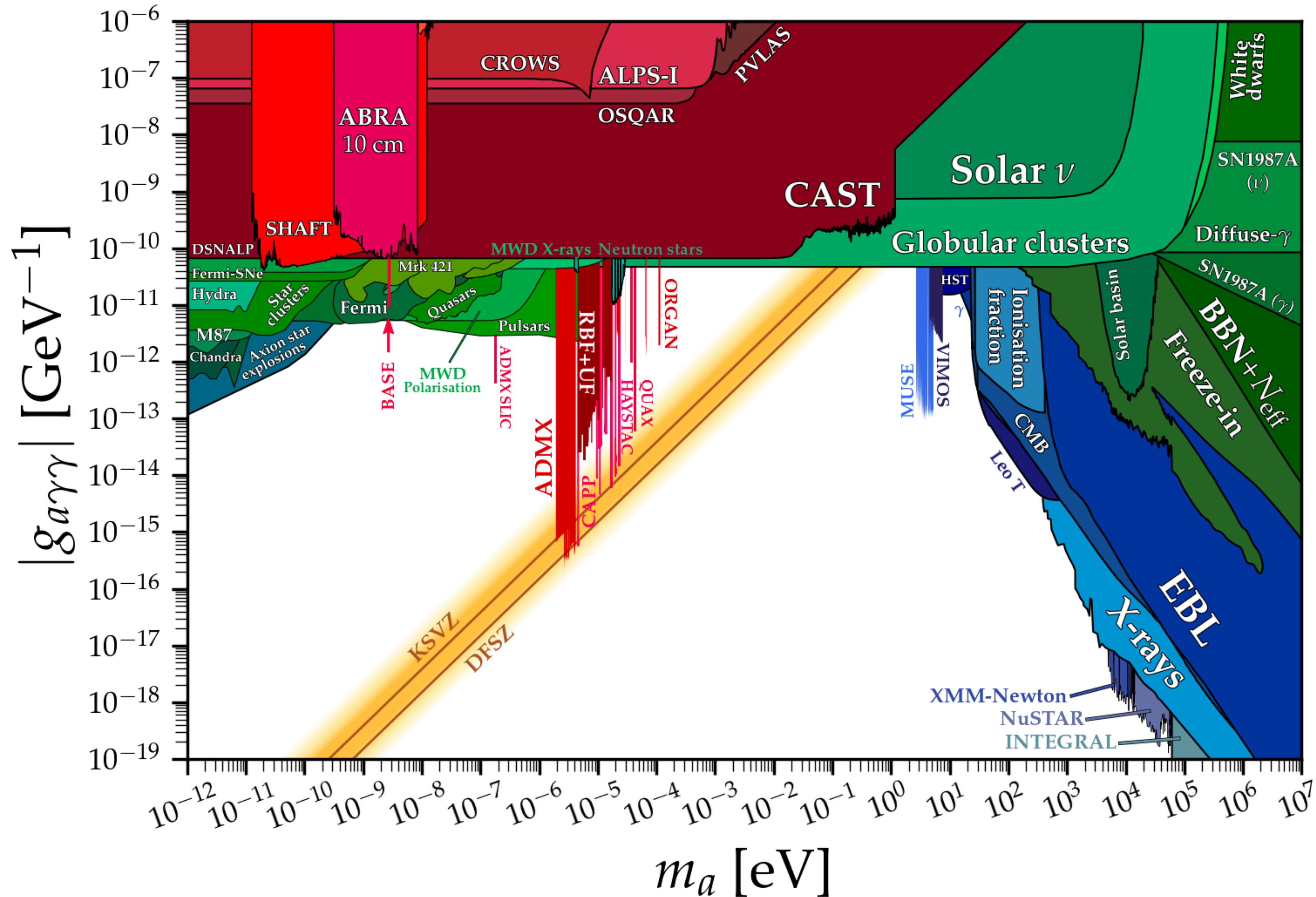


Heavy Charged Leptons L
(carry U(1)_d charge)

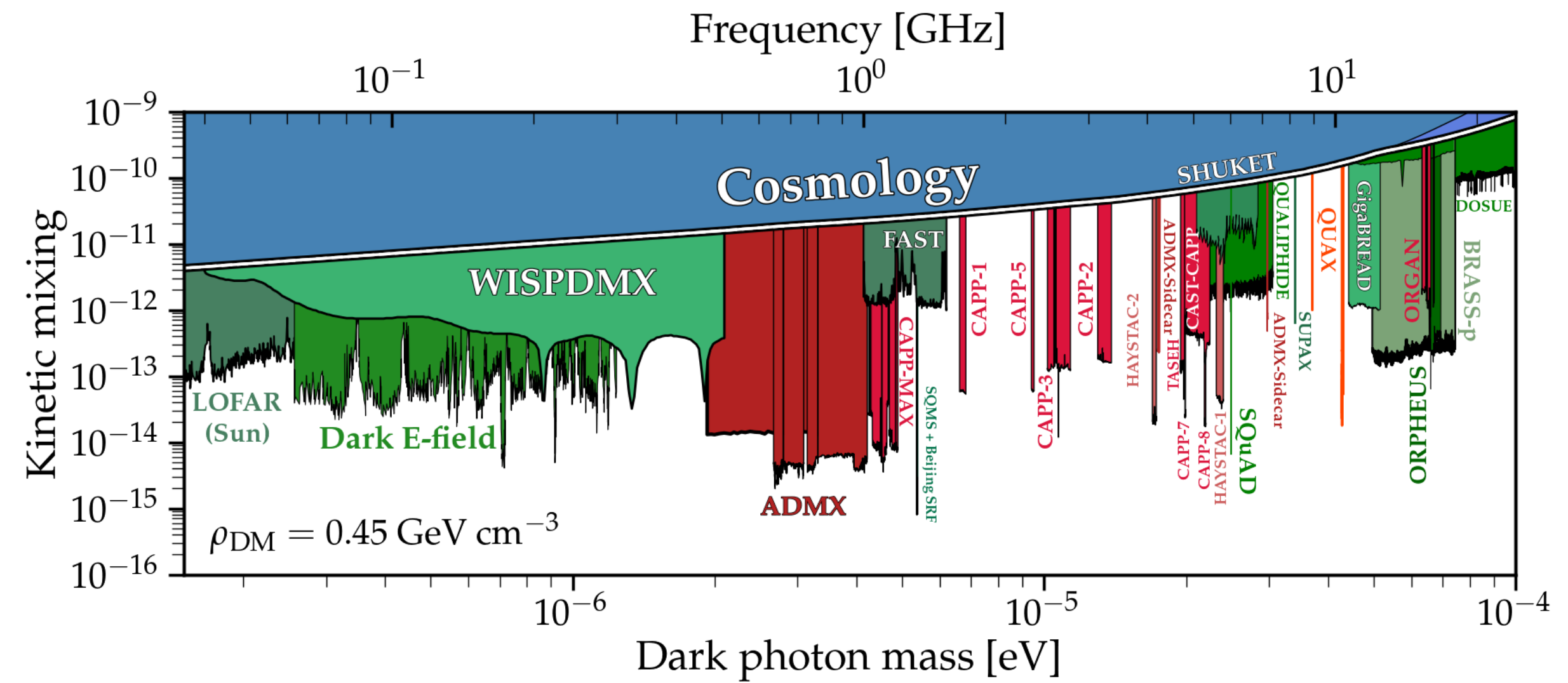
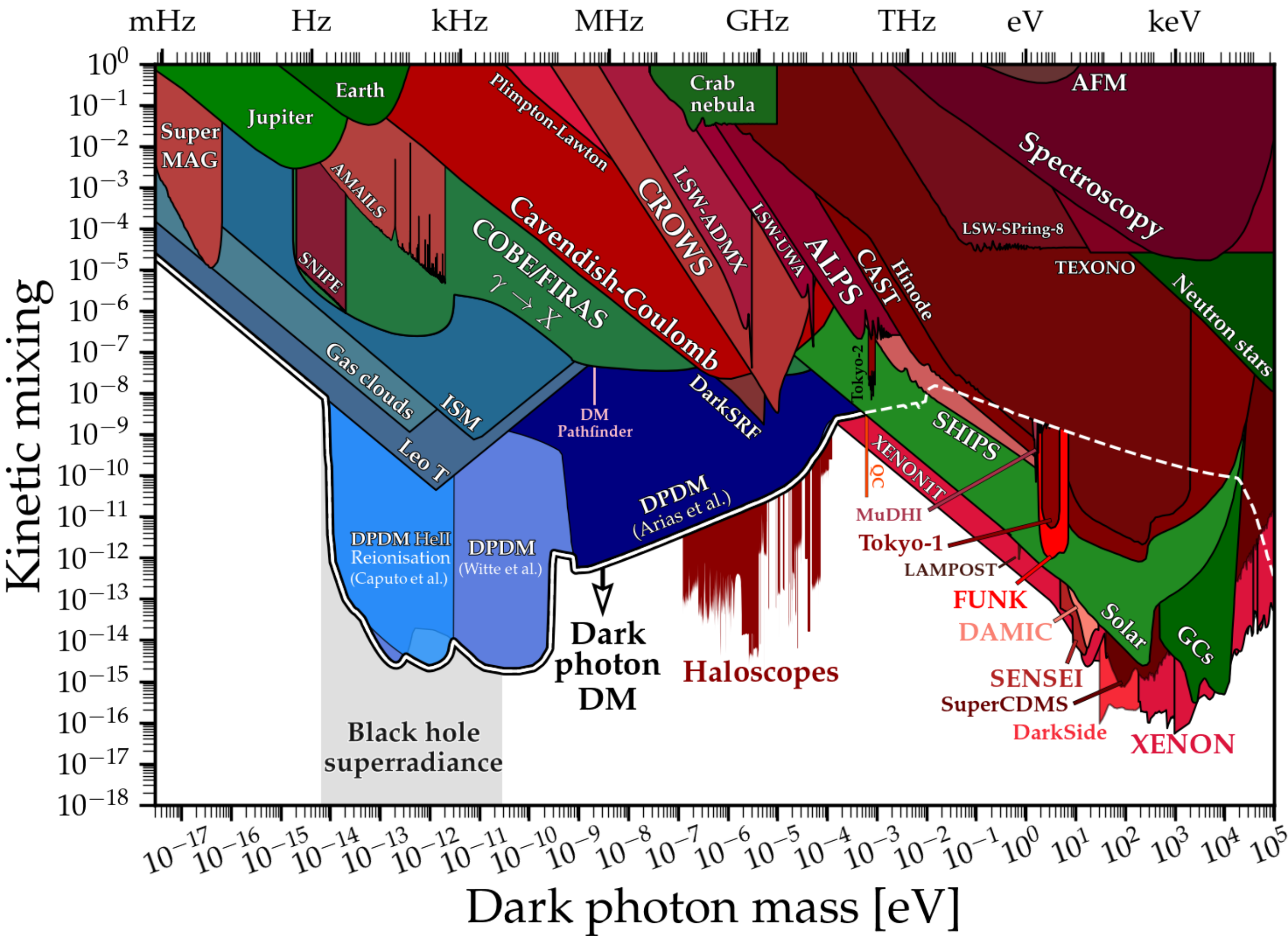
Dark Photon
(aka A' , U , Z_d , ...)

- Two free parameters: $m_{A'}$ and ϵ

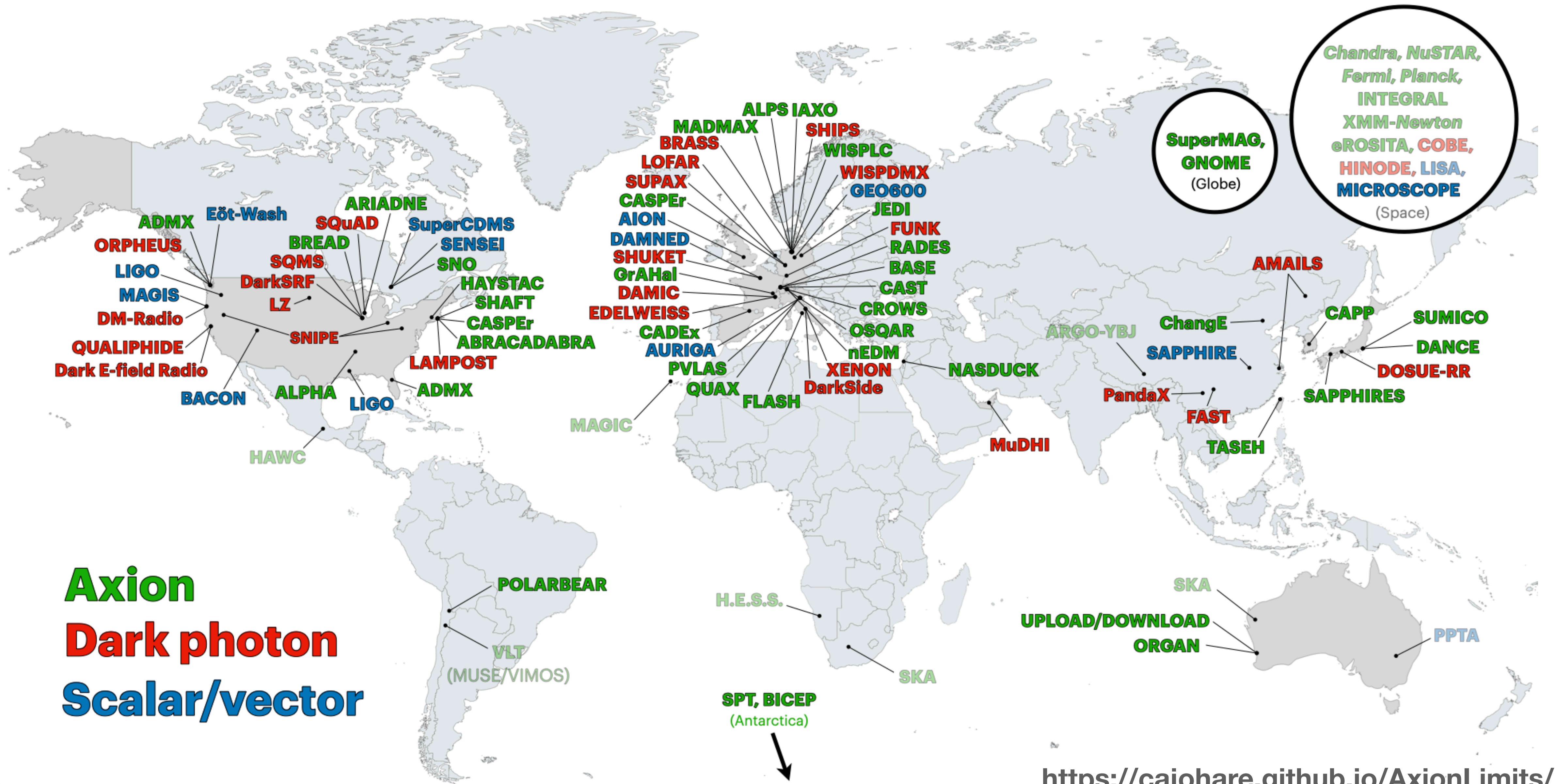
超轻玻色型暗物质探寻



超轻玻色型暗物质探寻

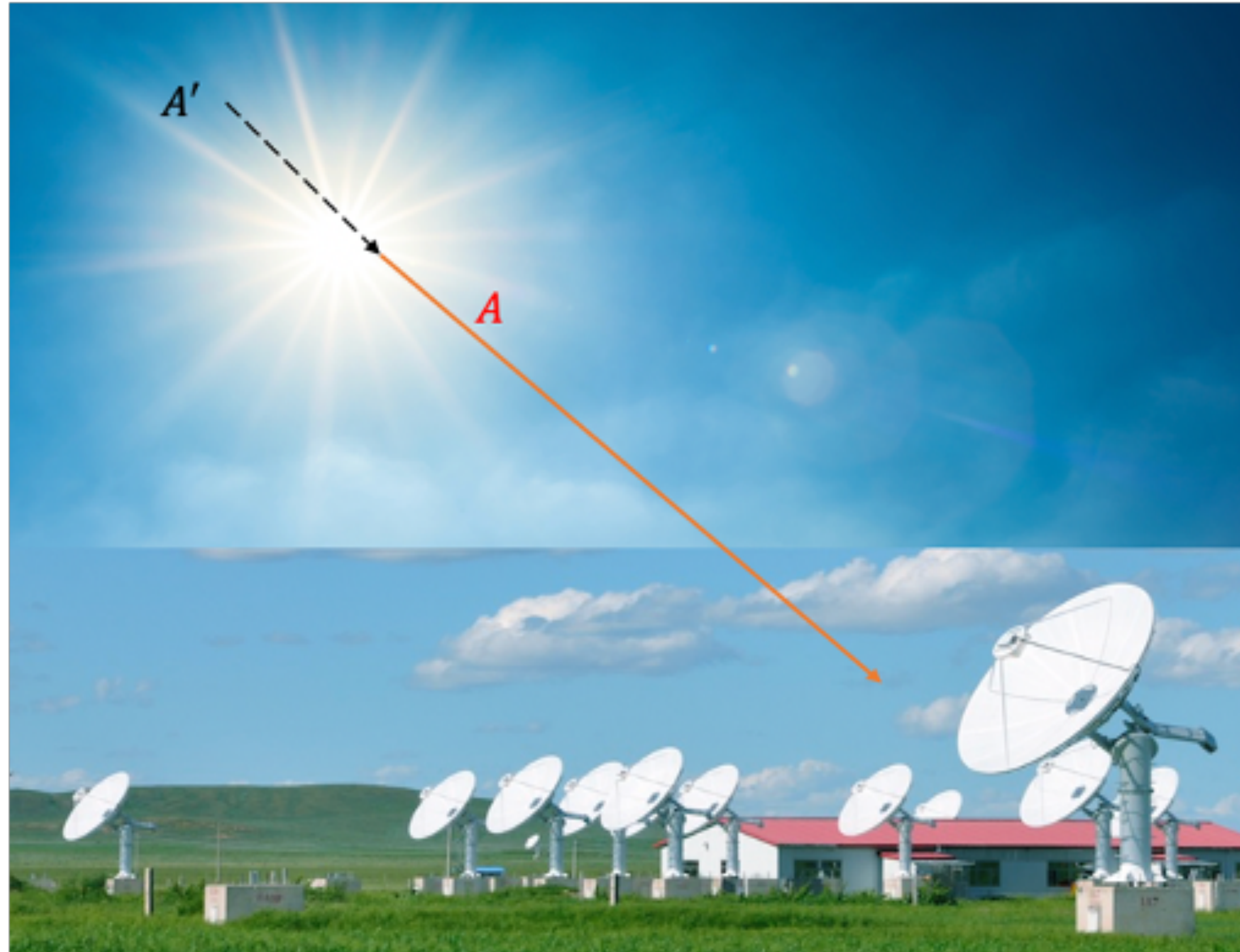


超轻玻色型暗物质探寻



Radio astronomy and ultralight bosonic dark matter

太阳物理



Radio telescope at Earth



**Dark photon dark matter
resonant conversion at solar corona**

An, Huang, JL, Xue, 2010.15836 (PRL 2021)
An, Chen, Ge, Liu, Luo, 2301.03622 (NC 2024)
Editor Highlights

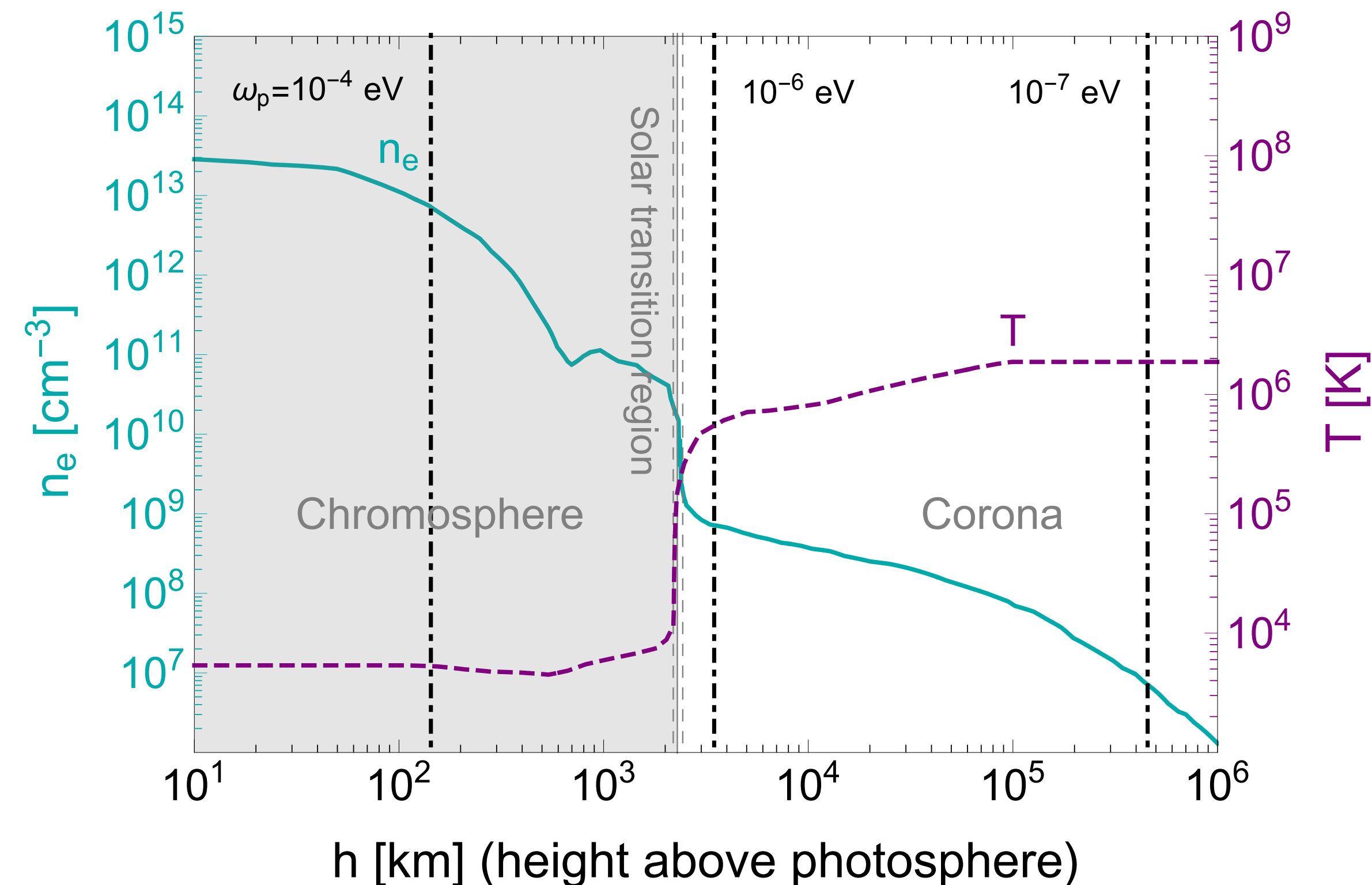
**Dark photon dark matter
conversion at radio telescope**

An, Ge, Guo, Huang, JL, Lu, 2207.05767 (PRL 2023, Featured in Physics)
reported by APS Physics, Phys.org, Physics Today

The dark photon dark matter conversion at solar corona

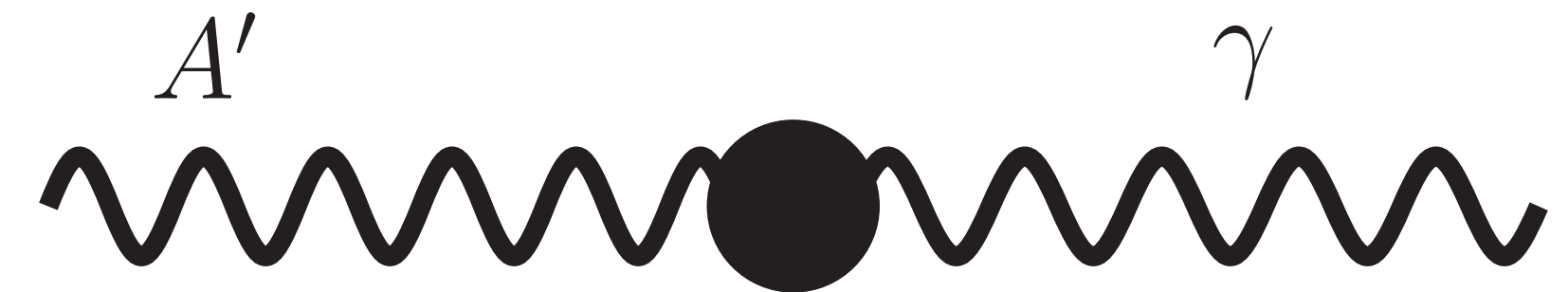
- The plasma frequency

$$\omega_p = \left(\frac{4\pi\alpha n_e}{m_e} \right)^{1/2} = \left(\frac{n_e}{7.3 \times 10^8 \text{ cm}^{-3}} \right)^{1/2} \mu\text{eV}$$



- Resonant conversion $A' \rightarrow \gamma$

- When $m_{A'} = \omega_p$



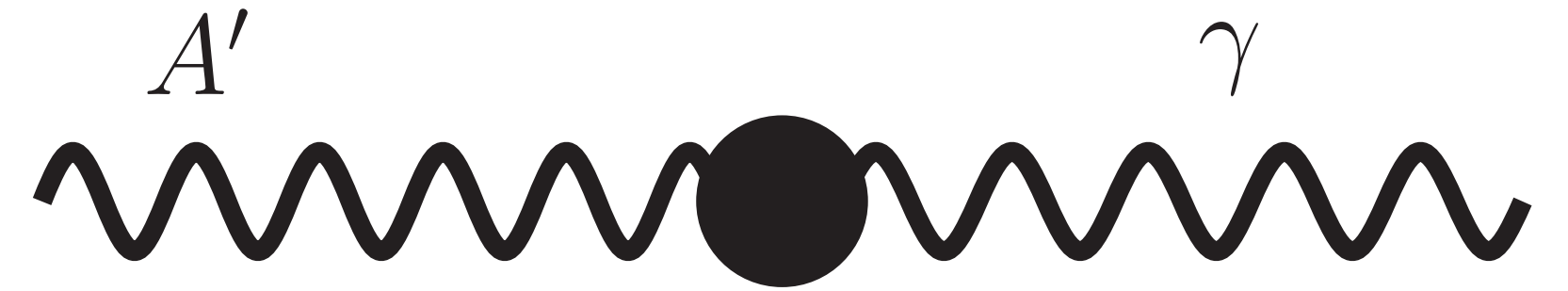
- For any A' mass, it can happen at a radius r_c

- $m_{A'} = \omega_p(r_c)$

- Can set limits for mass range $m_{A'} \in [10^{-8}, 10^{-5}] \text{ eV}$

The dark photon dark matter conversion at solar corona

- The resonant conversion probability (QFT method)



$$P_{A' \rightarrow \gamma}(v_r) = \frac{1}{3} \int \frac{dt}{2\omega} \frac{d^3p}{(2\pi)^3 2\omega} (2\pi)^4 \delta^4(p_{A'}^\mu - p_\gamma^\mu) \sum_{\text{pol}} |\mathcal{M}|^2$$

$$= \frac{2}{3} \times \pi \epsilon^2 m_{A'} v_r^{-1} \left| \frac{\partial \ln \omega_p^2(r)}{\partial r} \right|_{\omega_p(r)=m_{A'}}^{-1}$$

$$\mathcal{M} = -\epsilon m_{A'}^2 \left(\xi_\gamma^*(p) \cdot \xi_{A'}(p) \right)$$

$$\frac{1}{3} \sum_{\text{pol}} |\mathcal{M}|^2 = \frac{2}{3} \epsilon^2 m_{A'}^4$$

$$\int dt \delta(E_{A'} - E_\gamma) = 2\omega^{-1} \left(\frac{\partial \ln \omega_p^2}{\partial t} \right)^{-1}$$

$$\mathcal{L} \supset -\frac{1}{2} \epsilon F'_{\mu\nu} F^{\mu\nu}$$

- Due to the forced 4-momentum conservation, it applies to resonant conversion only.
- The [wave method](#) can work and is in agreement with QFT calculation

$$\left[\omega^2 - k^2 - \begin{pmatrix} \omega_p^2 & -\epsilon m_{A'}^2 \\ -\epsilon m_{A'}^2 & m_{A'}^2 \end{pmatrix} \right] \begin{pmatrix} A(r, t) \\ A'(r, t) \end{pmatrix} = 0$$

Sensitivity of the radio telescope

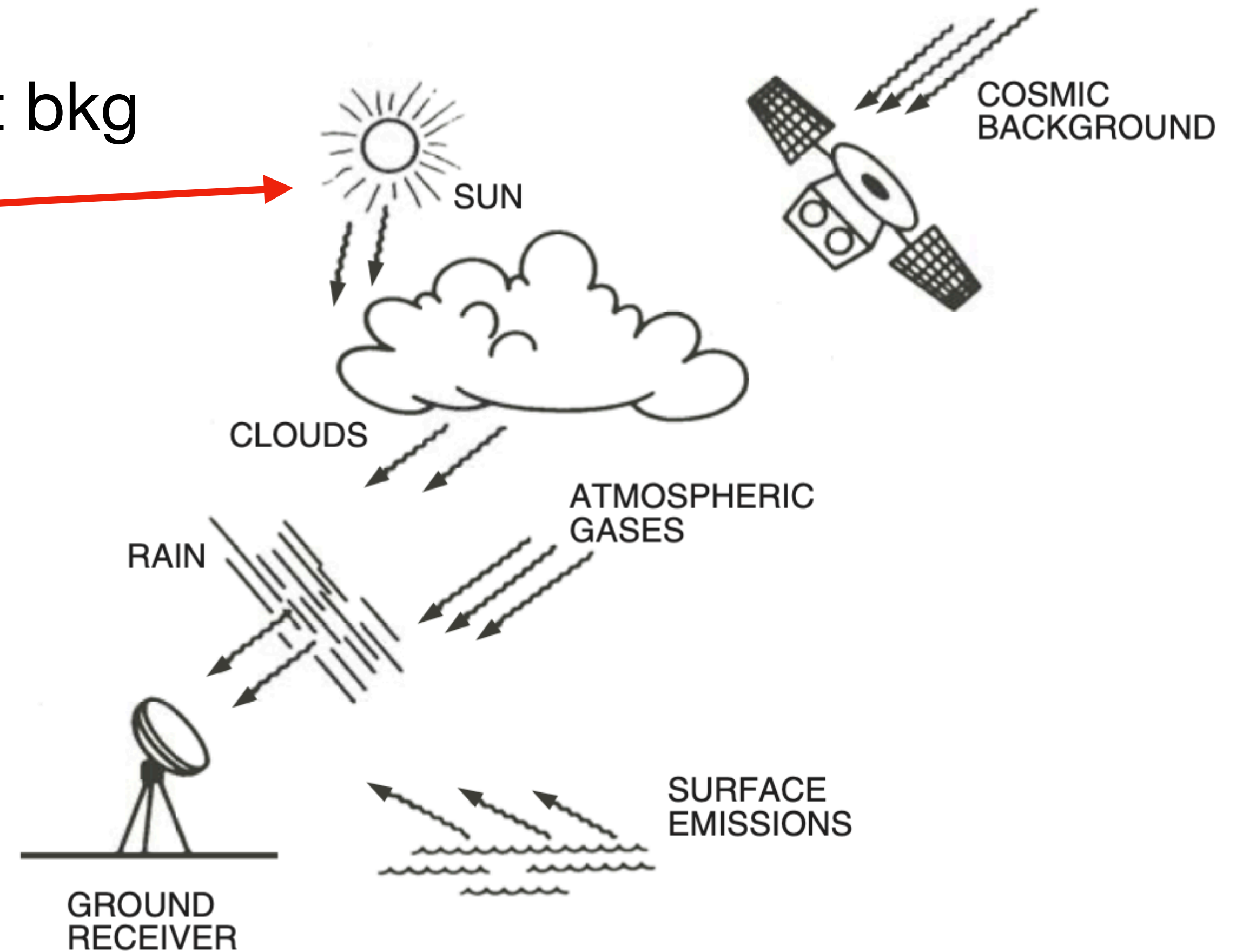
- The system equivalent flux density
- For solar observation, Sun is the largest bkg

$$SEFD = 2k_B \frac{T_{\text{sys}} + T_{\odot}^{\text{nos}}}{A_{\text{eff}}}$$

- The minimum detectable flux density

$$S_{\text{min}} = \frac{SEFD}{\eta_s \sqrt{n_{\text{pol}} \mathcal{B} t_{\text{obs}}}}$$

Name	f [MHz]	B_{res} [kHz]	$\langle T_{\text{sys}} \rangle$ [K]	$\langle A_{\text{eff}} \rangle$ [m ²]
SKA1-Low	(50, 350)	1	680	2.2×10^5
SKA1-Mid B1	(350, 1050)	3.9	28	2.7×10^4
SKA1-Mid B2	(950, 1760)	3.9	20	3.5×10^4
LOFAR	(10, 80)	195	28,110	1,830
LOFAR	(120, 240)	195	1,770	1,530

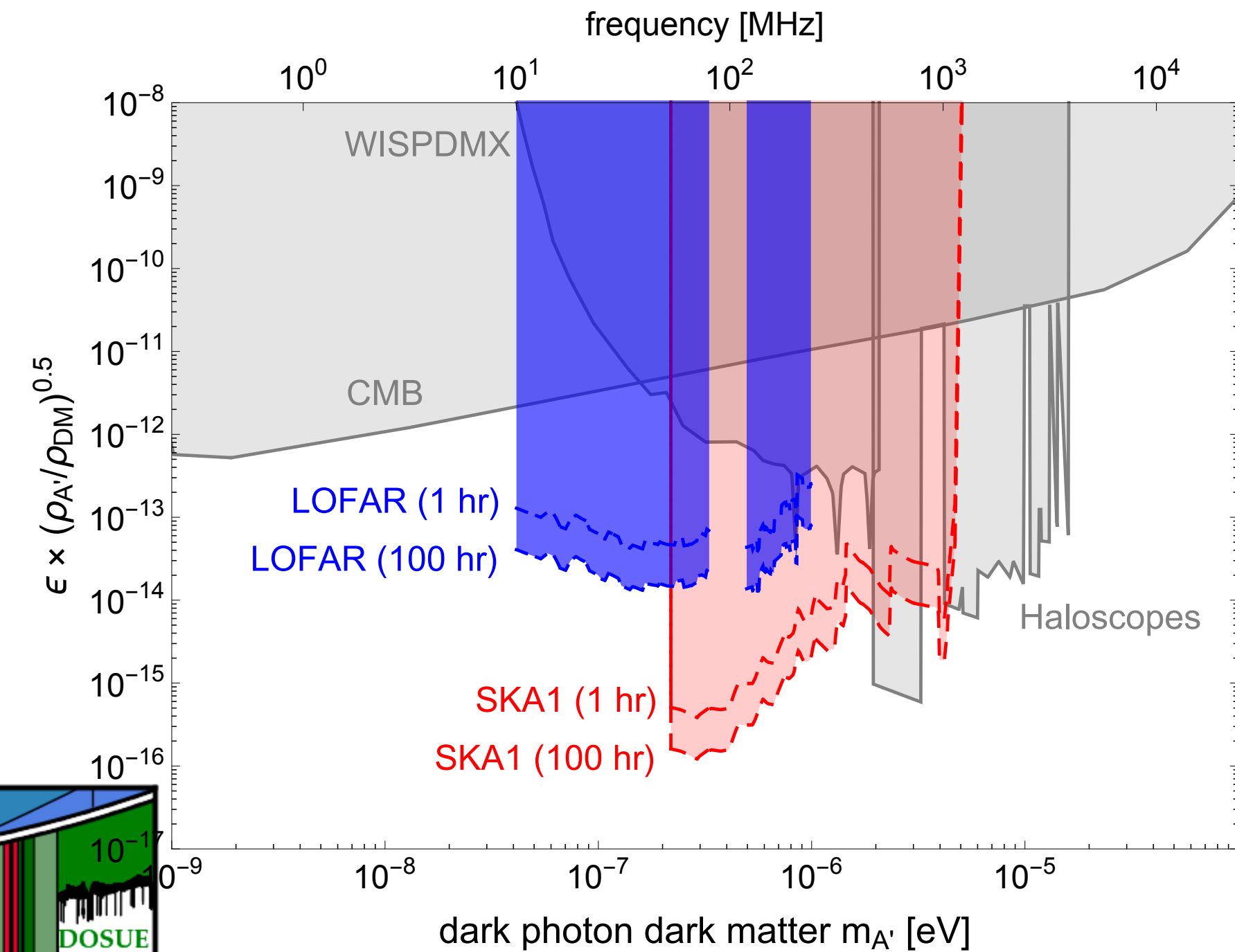
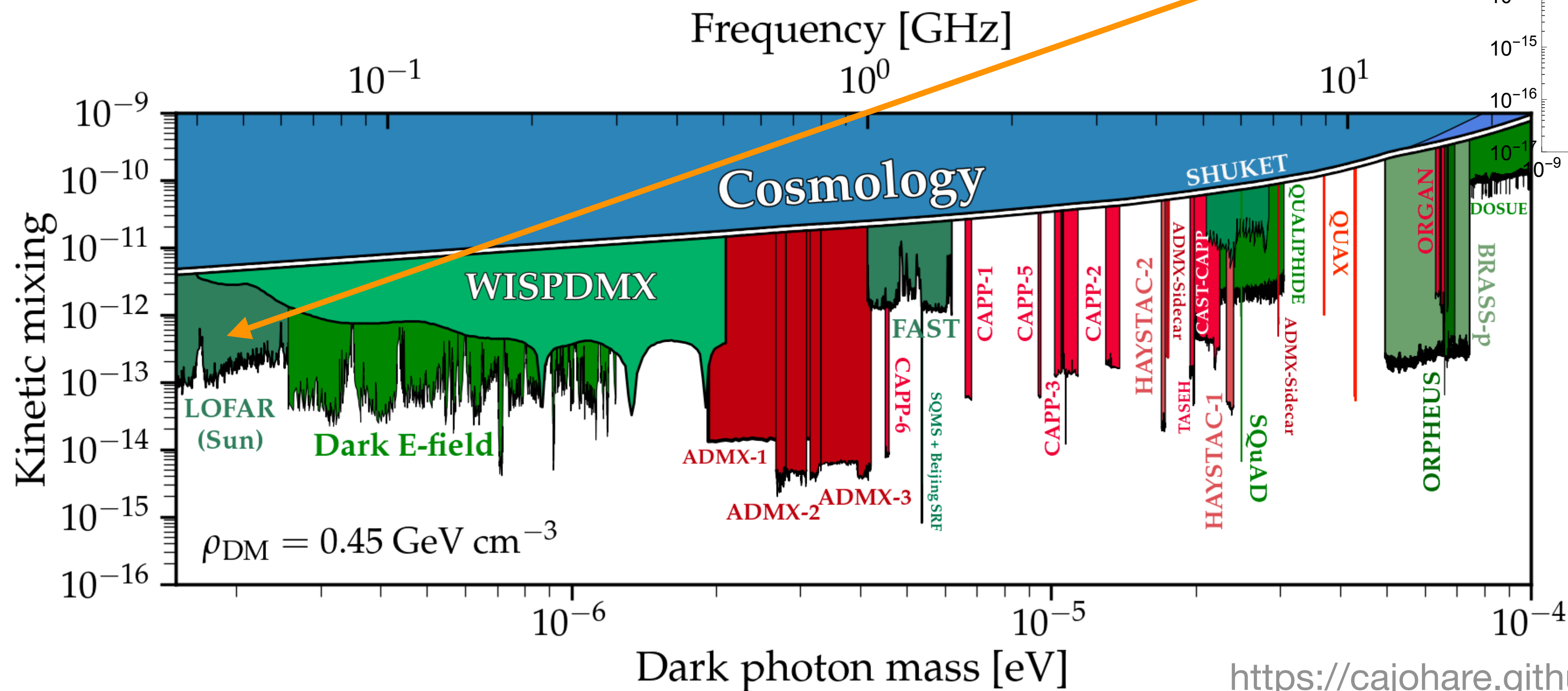


The sensitivity of DPDM from solar resonant conversion

- Attenuation of the mono-photon signal is considered,

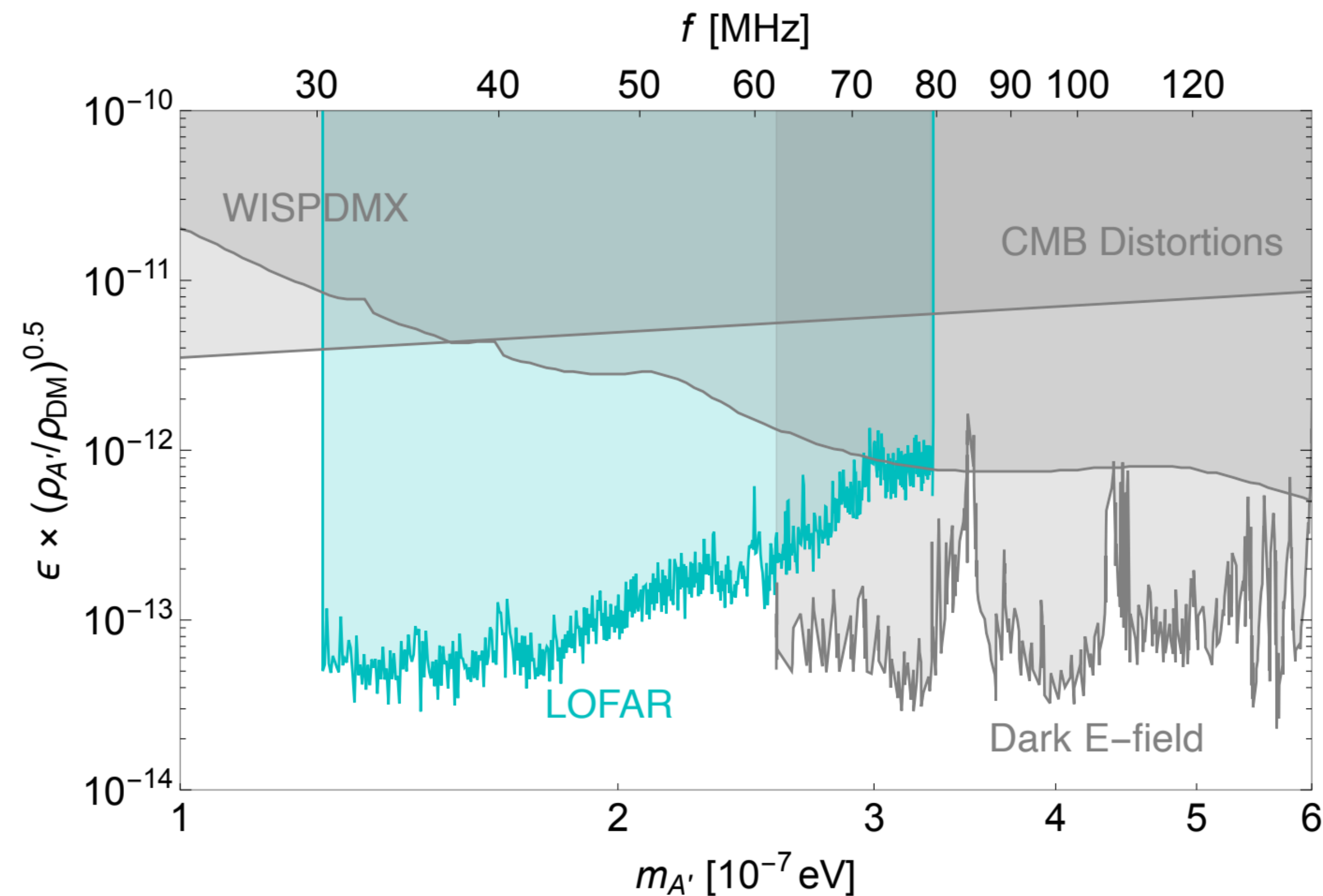
$$S_{\text{sig}} \times P_s = S_{\text{min}}$$

- 10 MHz lower end from LOFAR threshold
- 1 GHz higher end due to photon attenuation

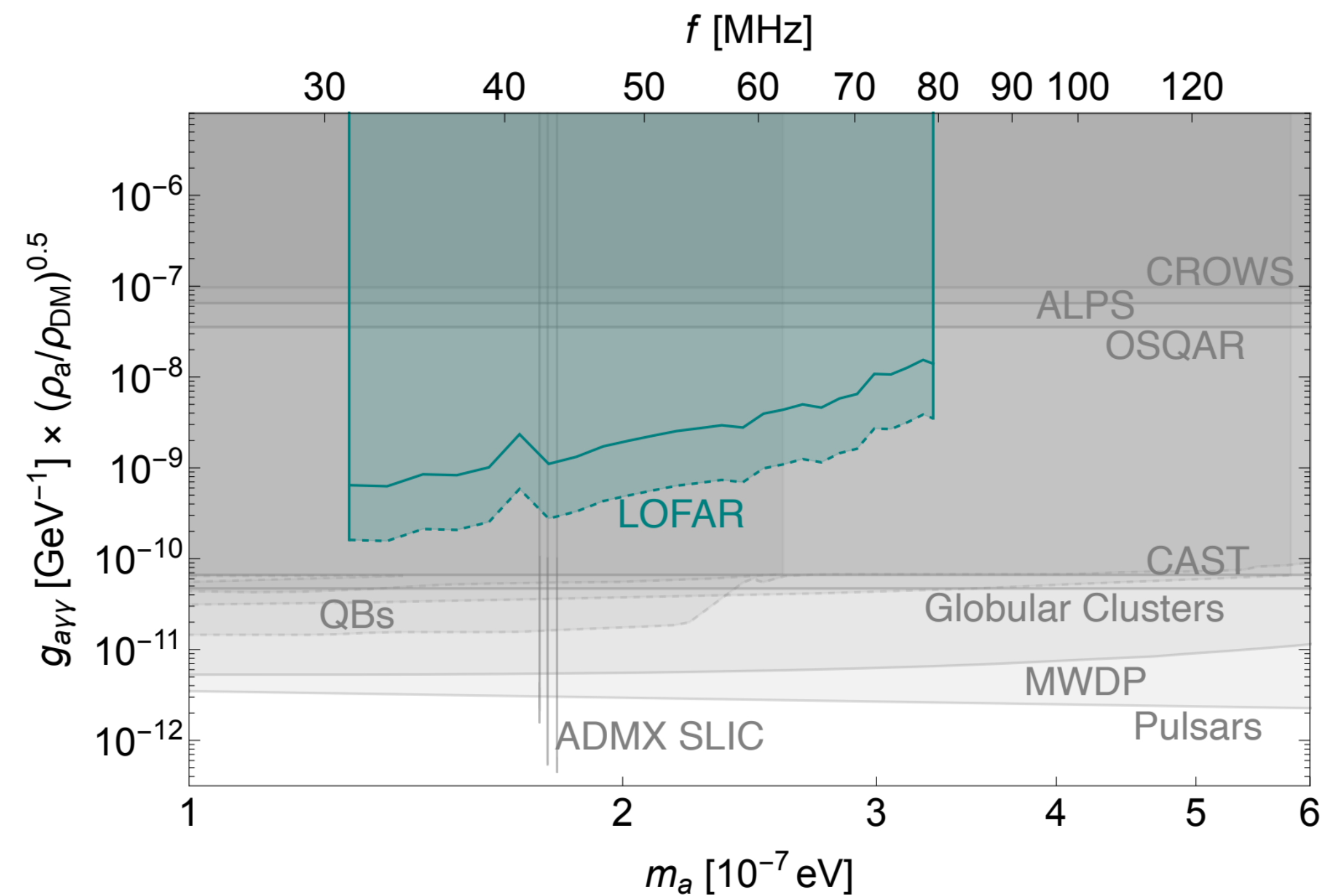


Phys.Rev.Lett. 126 (2021) 18, 181102

The sensitivity of DPDM from LOFAR data



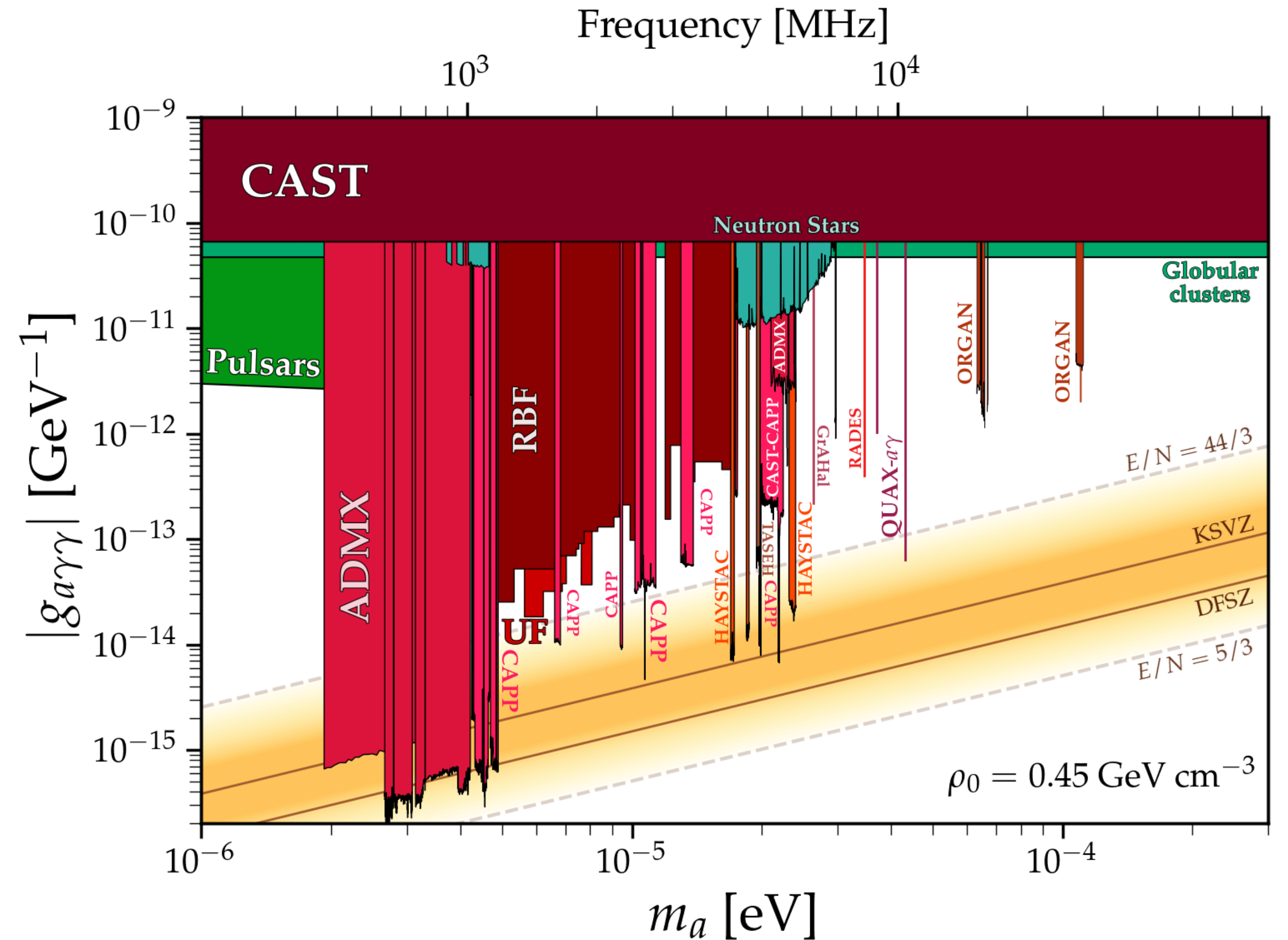
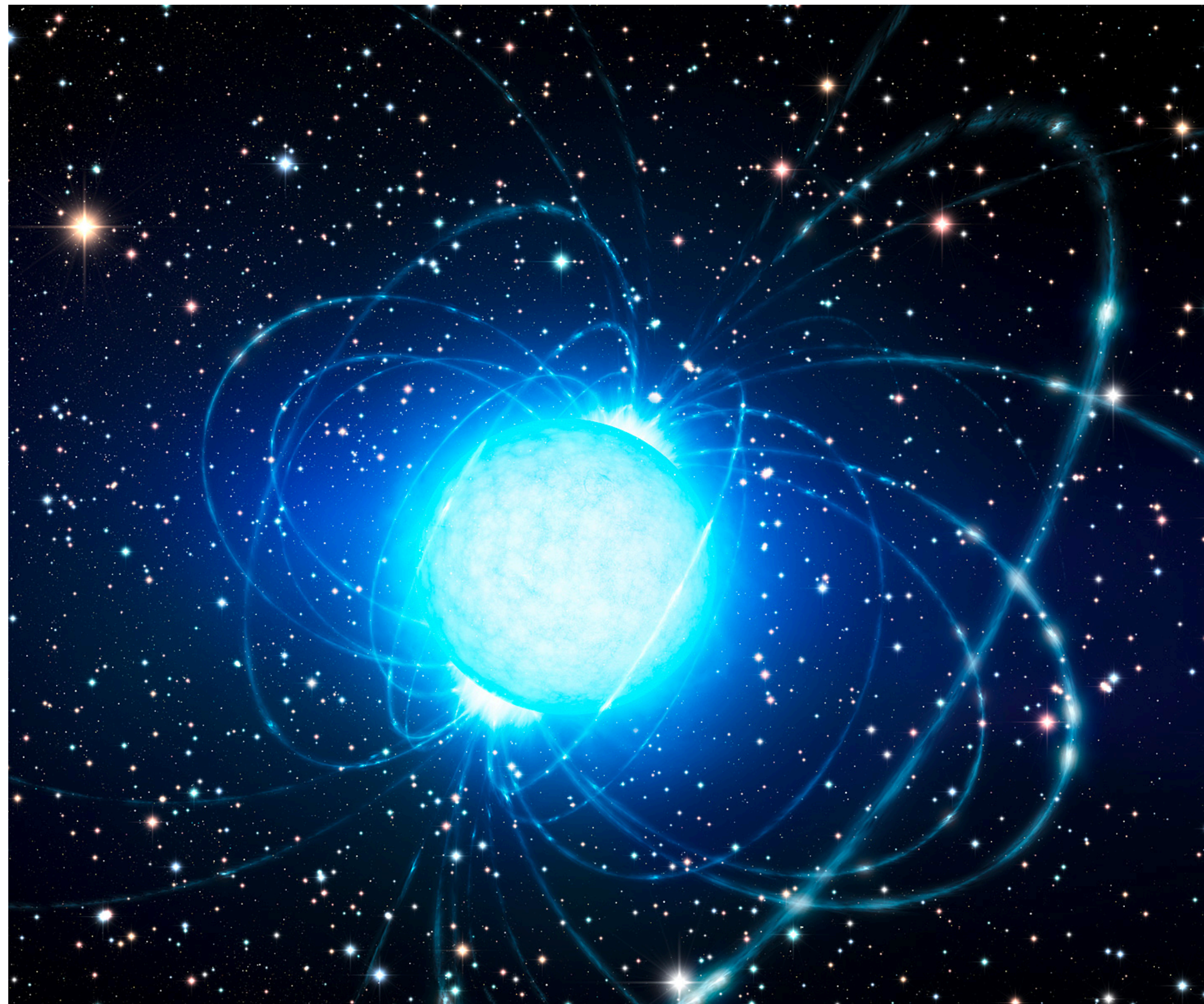
暗光子暗物质



轴子暗物质

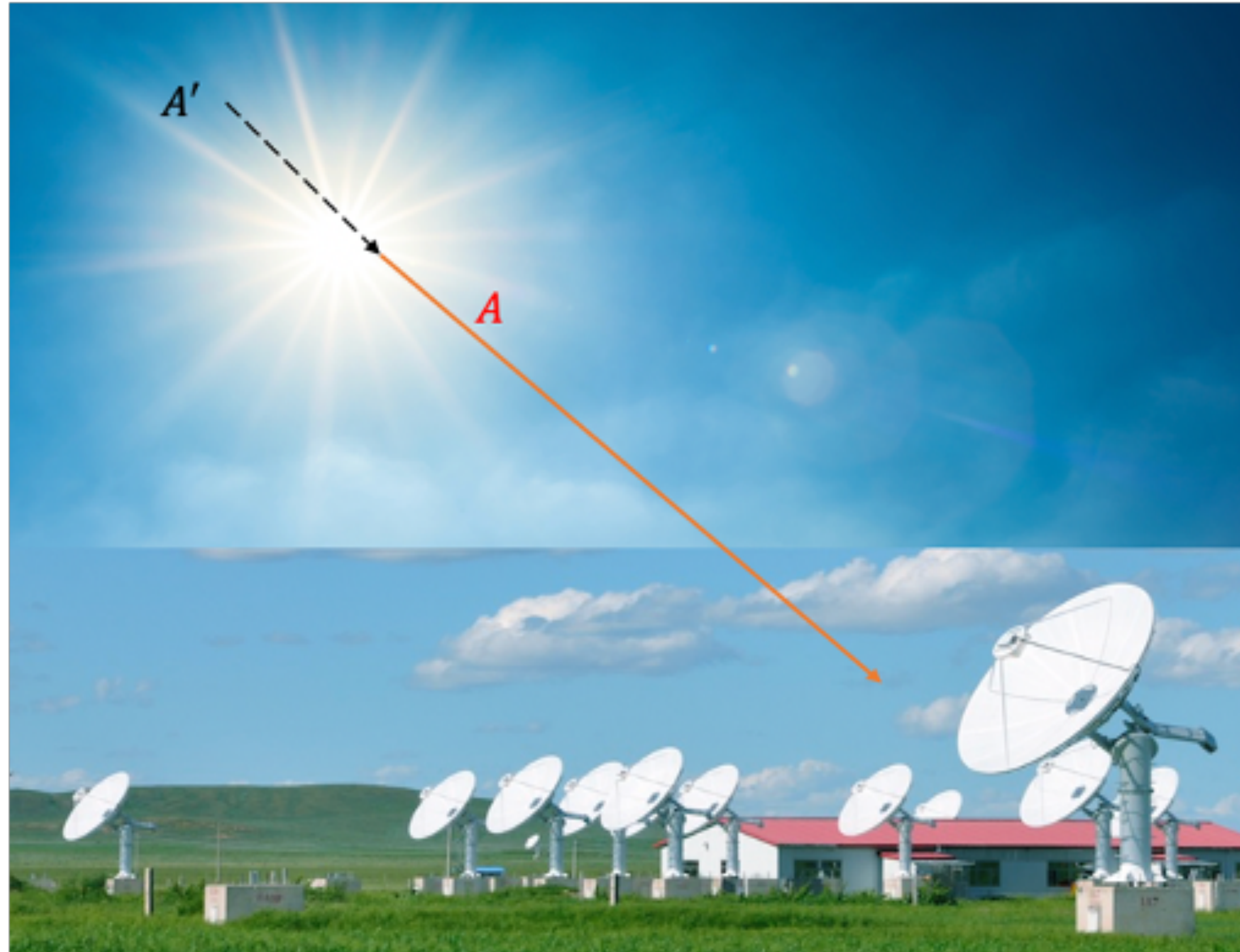
中子星观测与超轻轴子暗物质

- 轴子在中子星(Magnetar)强磁场中转化为单频光子



Radio astronomy and ultralight bosonic dark matter

Solar Physics



Radio telescope at Earth



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Direct detection of DPDM using **dish antenna** radio telescope

- Lagrangian

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{4}F'_{\mu\nu}F'^{\mu\nu} + \frac{1}{2}m_{A'}^2 A'_\mu A'^\mu - \epsilon e A'_\mu j_{\text{em}}^\mu + e A_\mu j_{\text{em}}^\mu.$$

- Extended Maxwell Eqs

$$\nabla \cdot \mathbf{E}' = -\epsilon\rho - m_{A'}^2 A'^0,$$

$$\nabla \cdot \mathbf{B}' = 0,$$

$$\nabla \times \mathbf{E}' + \frac{\partial \mathbf{B}'}{\partial t} = 0,$$

$$\nabla \times \mathbf{B}' - \frac{\partial \mathbf{E}'}{\partial t} = -\epsilon\mathbf{J} - m_{A'}^2 \mathbf{A}',$$

- Perfect conductor

$$\mathbf{J} = \sigma(\mathbf{E} - \epsilon\mathbf{E}_D)$$

$$\nabla \cdot \mathbf{J} + \frac{\partial \rho}{\partial t} = 0$$



FAST radio telescope

Direct detection of DPDM using dish antenna radio telescope

- The feature of current on a metal conductor plate induced by DPDM

$$i_{tot,x} = i_{up,x} + i_{down,x} \approx -2i\epsilon m_{A'} A'_x,$$

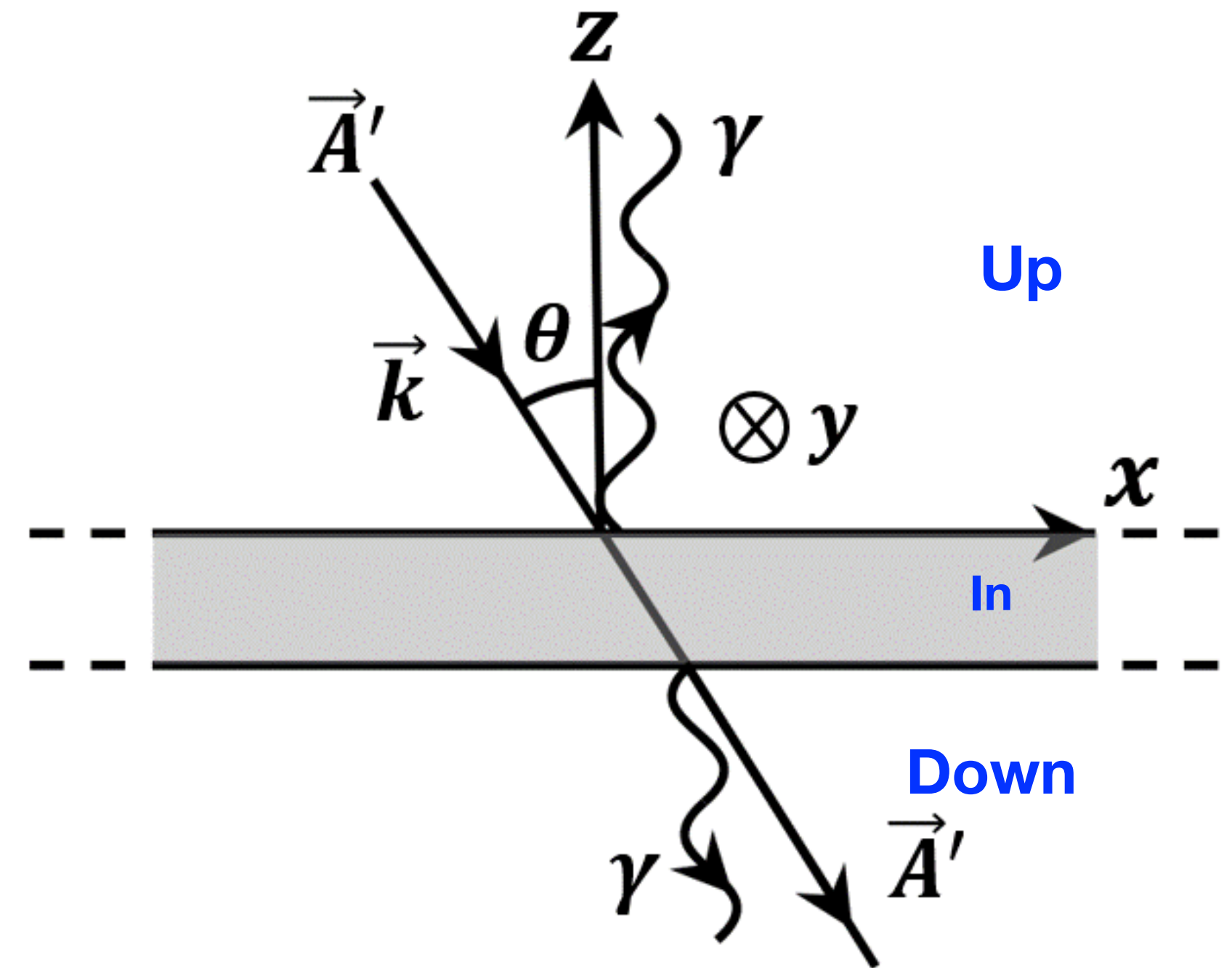
$$i_{tot,y} = i_{up,y} + i_{down,y} \approx -2i\epsilon m_{A'} A'_y,$$

$$\mathbf{J} = 0.$$

- Solving the reflected EM field
 - (always perpendicular to the surface)
 - Oscillating dipole unit

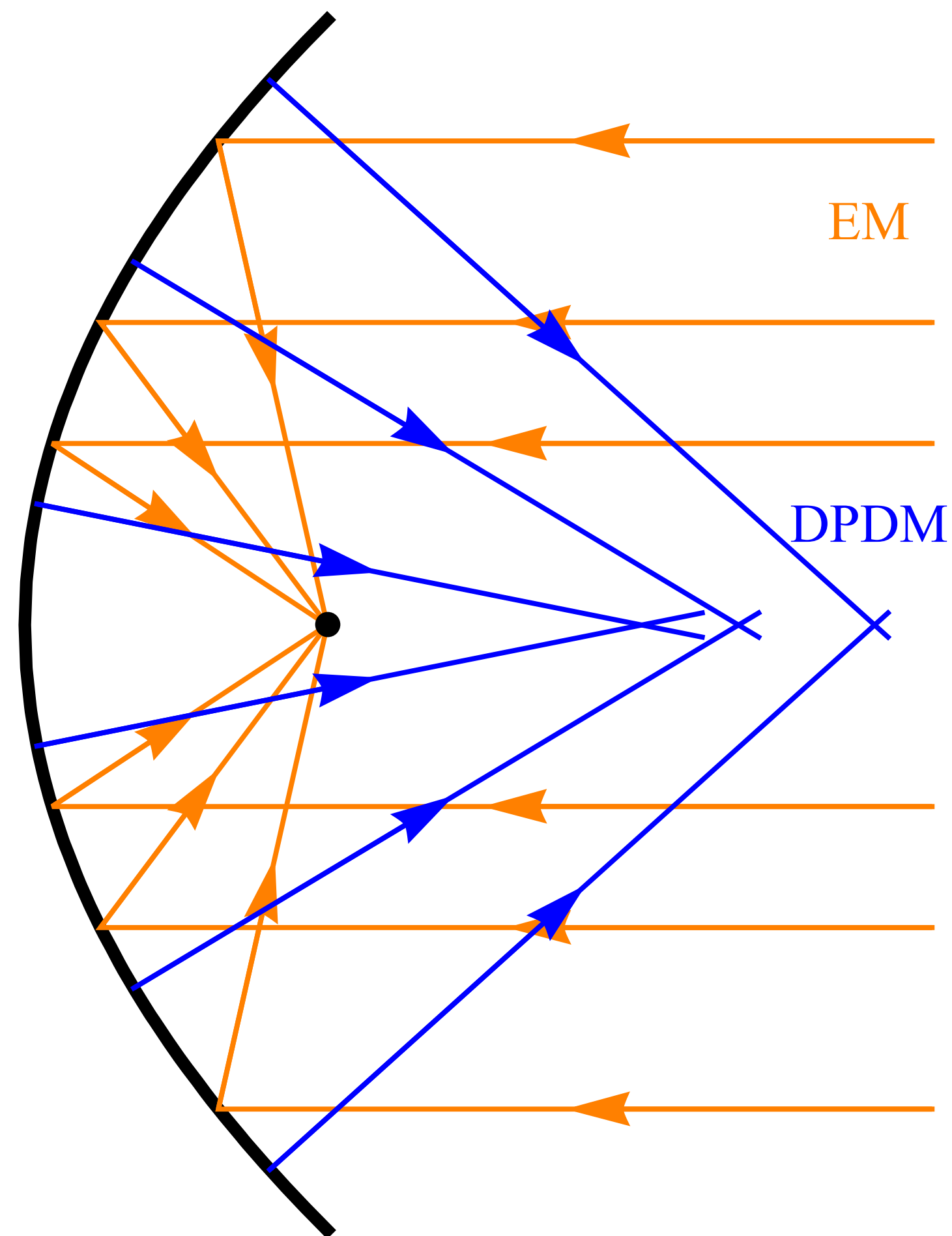
$$\boxed{d\mathbf{p} = 2\epsilon \mathbf{A}'_{\parallel} dS} \longrightarrow \boxed{\mathbf{B} = -\frac{\epsilon m_{A'}^2}{2\pi} \int dS_1 \mathbf{A}'_{\parallel} \times (\mathbf{r} - \mathbf{r}_1) \frac{e^{im_{A'}|\mathbf{r}-\mathbf{r}_1|}}{|\mathbf{r} - \mathbf{r}_1|^2}}$$

- Regular shapes of reflector can be solved
- General shapes need numerical integration

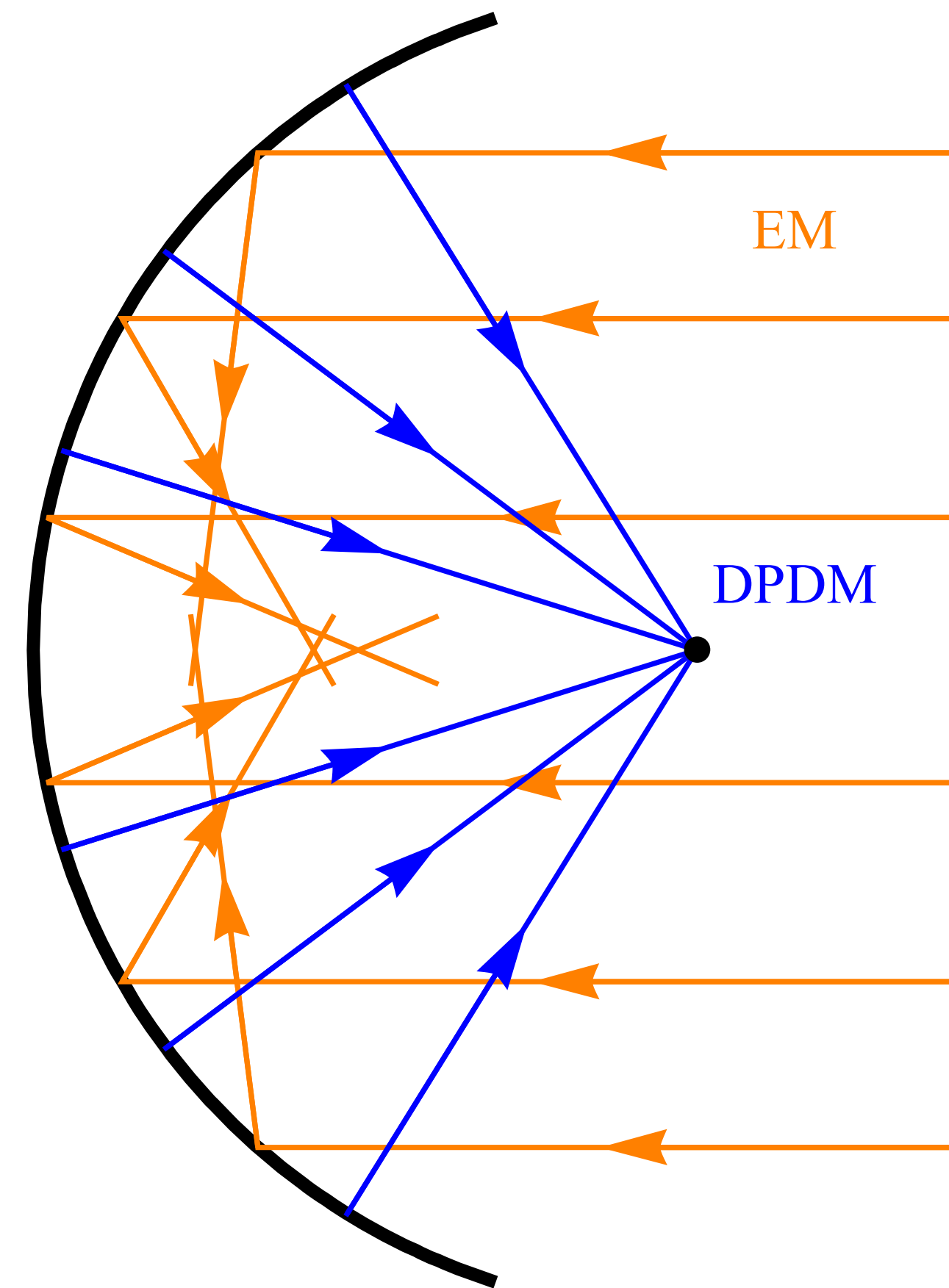


Direct detection of DPDM using **dish antenna** radio telescope

- The signal feature for DPDM with different mirrors



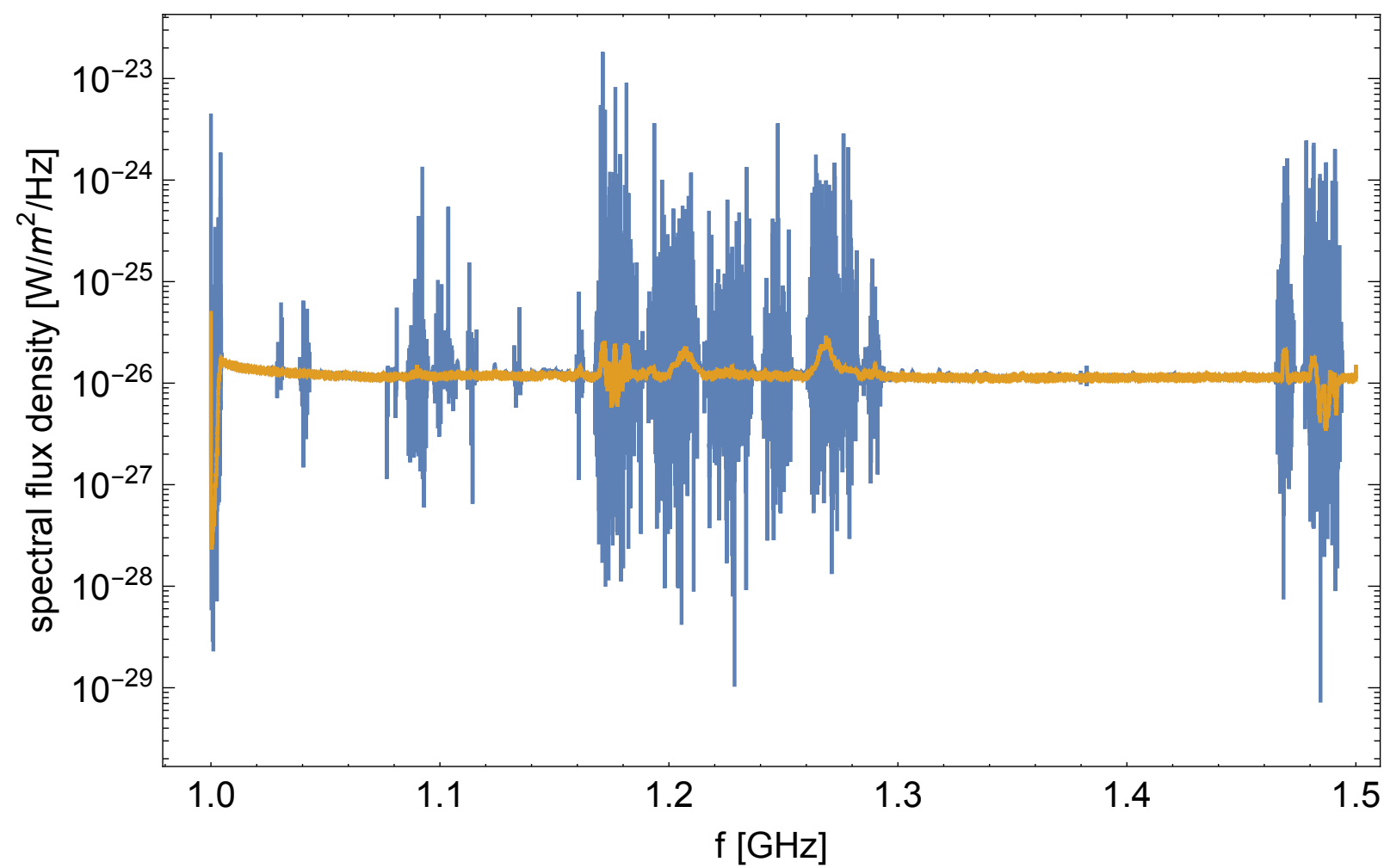
**Parabolic
Mirror**



**Spherical
Mirror**

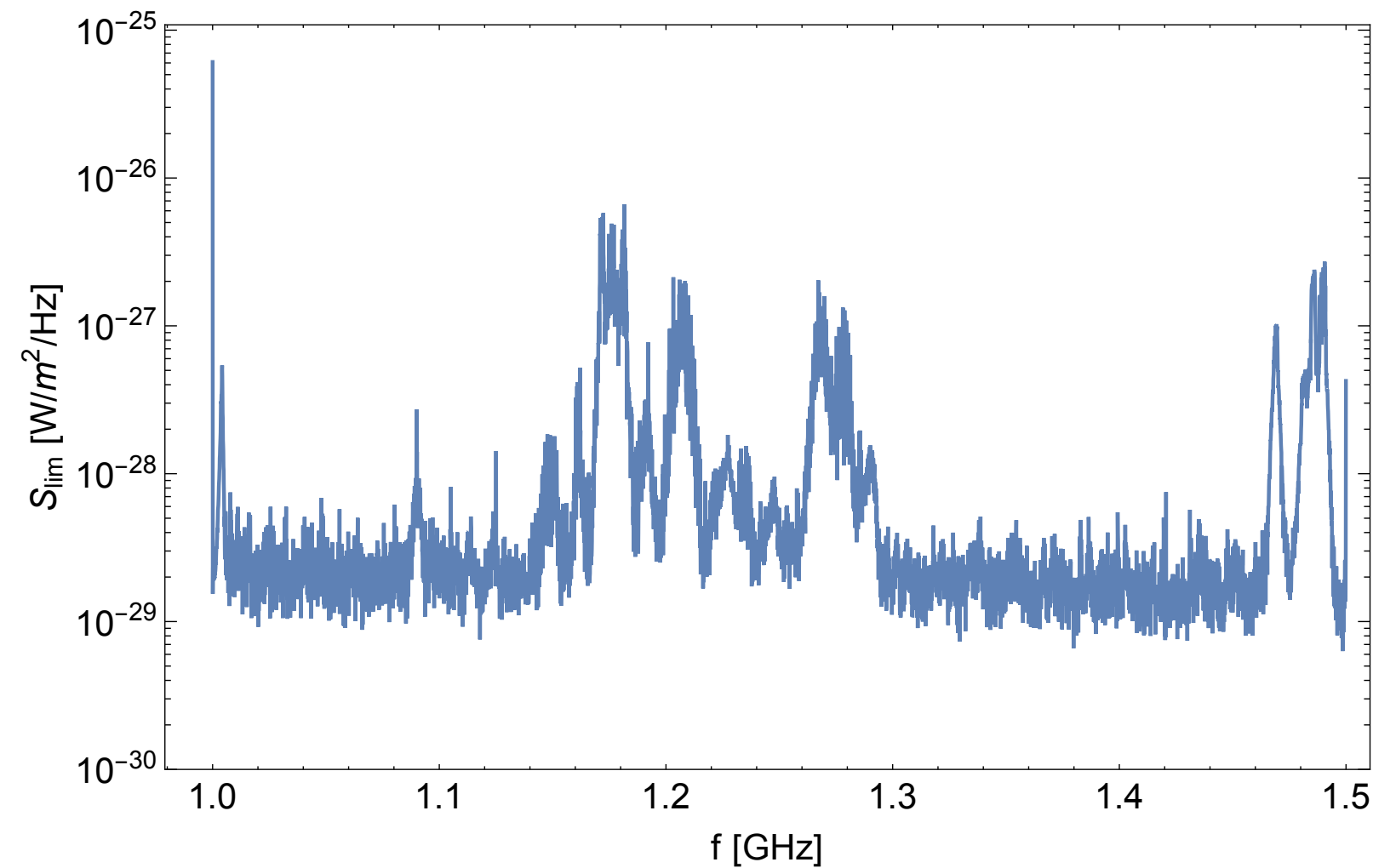
Constraints from FAST observation data

- ‘Bump hunting’ in the frequency data
- Using likelihood-based statistical test



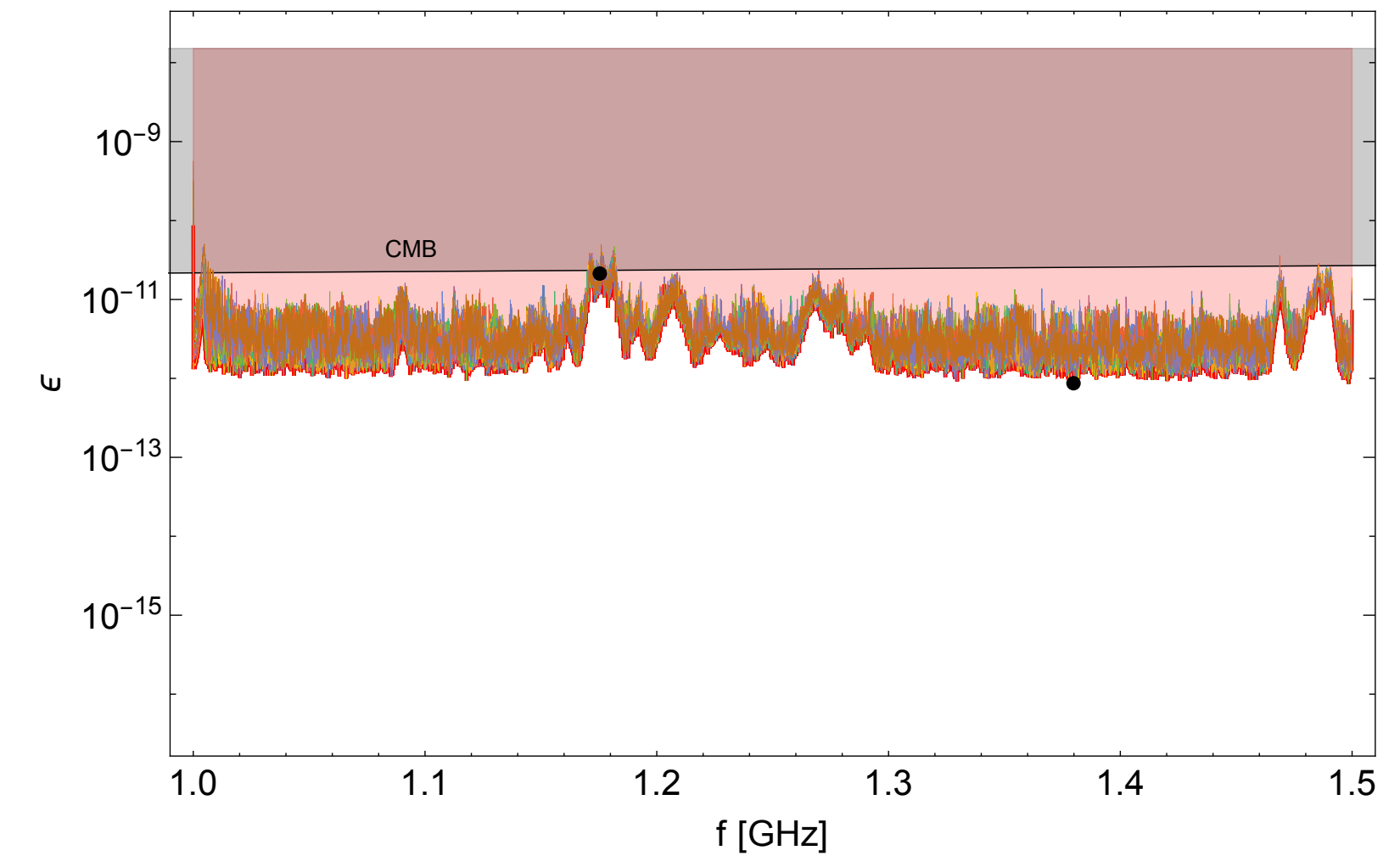
FAST data

$$S_{\text{obs}} \sim 10^{-26} \text{ W m}^{-2} \text{ Hz}^{-1}$$



FAST limit on line signal

$$S_{\text{lim}} \sim 10^{-29} \text{ W m}^{-2} \text{ Hz}^{-1}$$

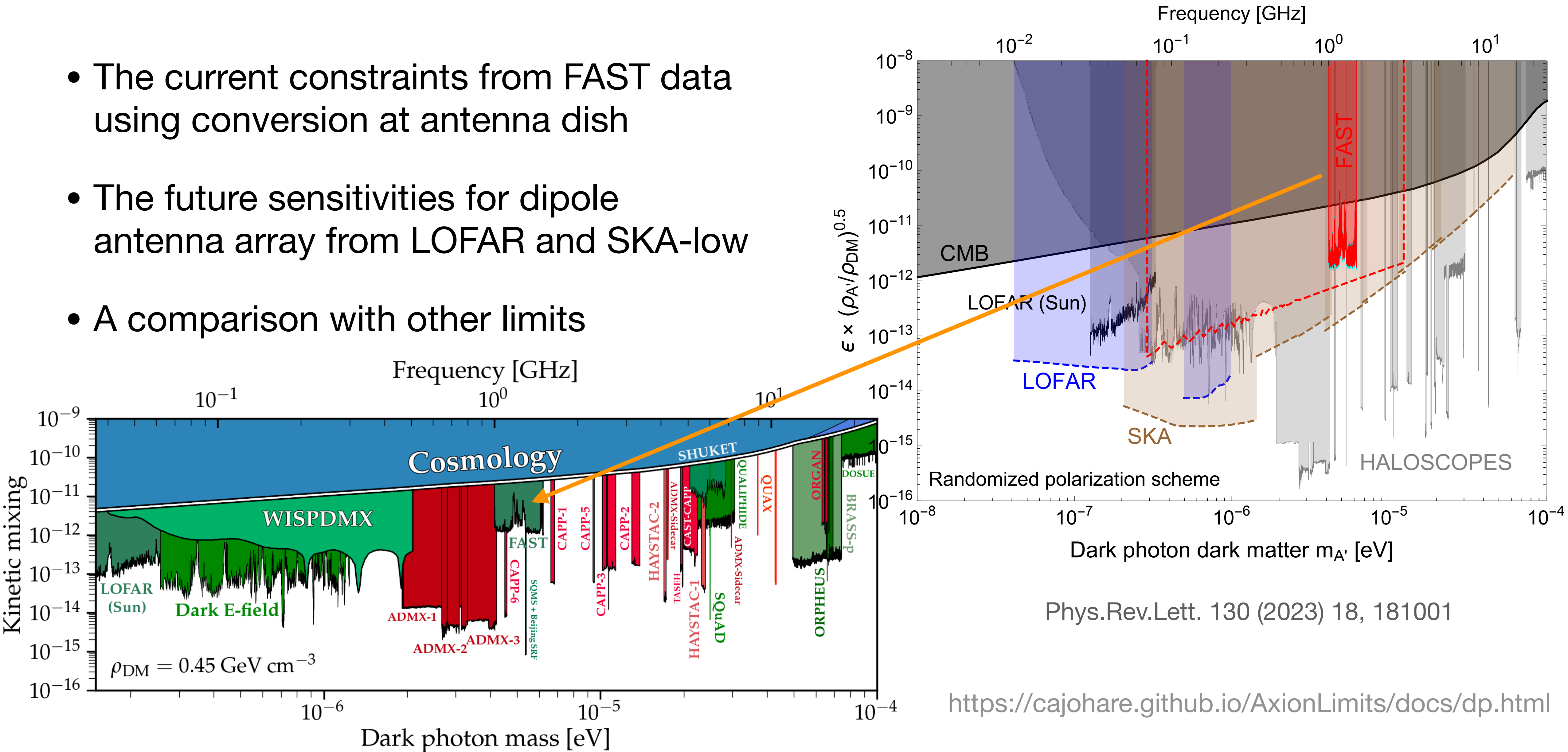


FAST limit on DPDM

$$\epsilon \sim 10^{-12}$$

The results for direct detection of DPDM using FAST radio telescope

- The current constraints from FAST data using conversion at antenna dish
- The future sensitivities for dipole antenna array from LOFAR and SKA-low
- A comparison with other limits



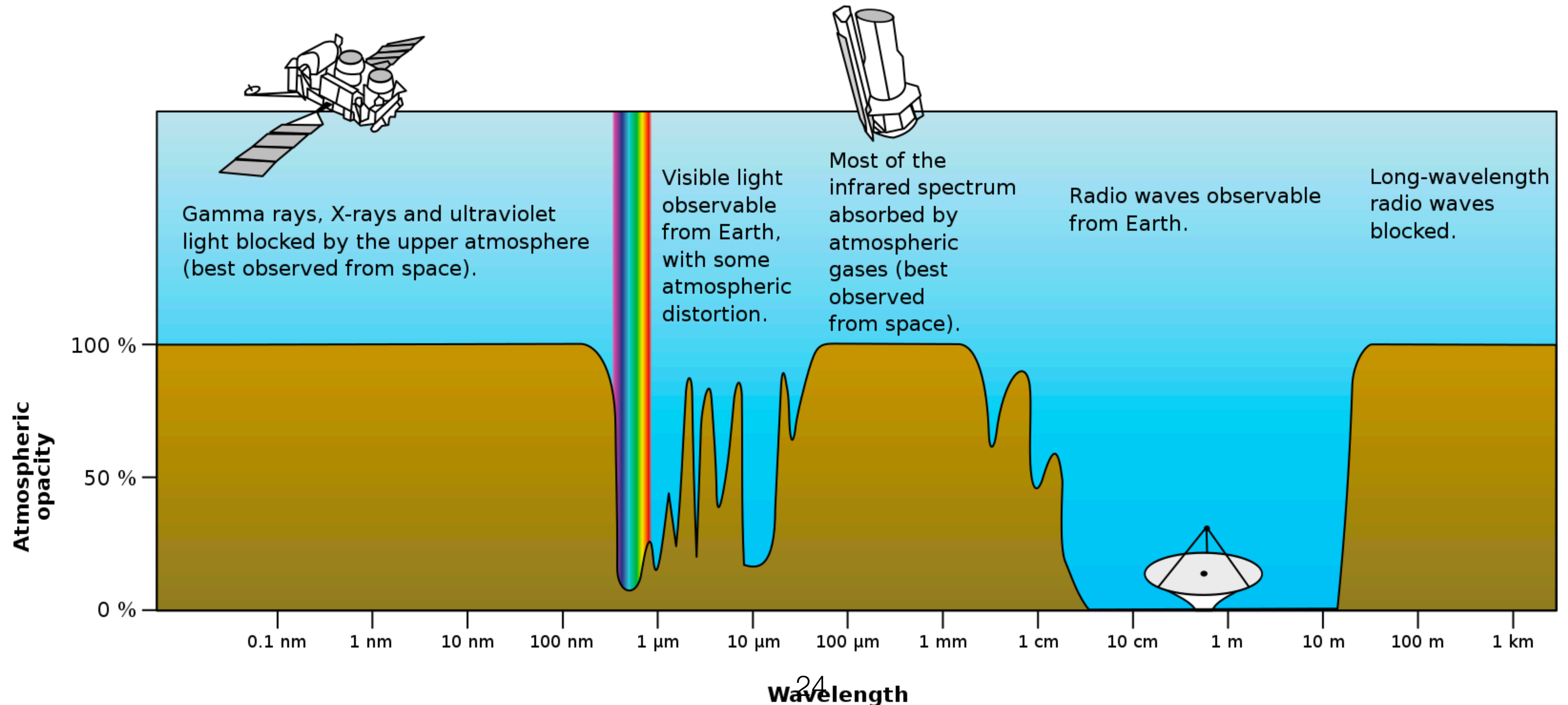
未来与展望： 太阳物理观测与超轻暗物质探测

- “千眼天珠”稻城太阳射电望远镜 (Daocheng Solar Radio Telescope)
- 313 **parabolic antennas** of 6-meter diameter each
- Operational frequency: 150 MHz — 450 MHz



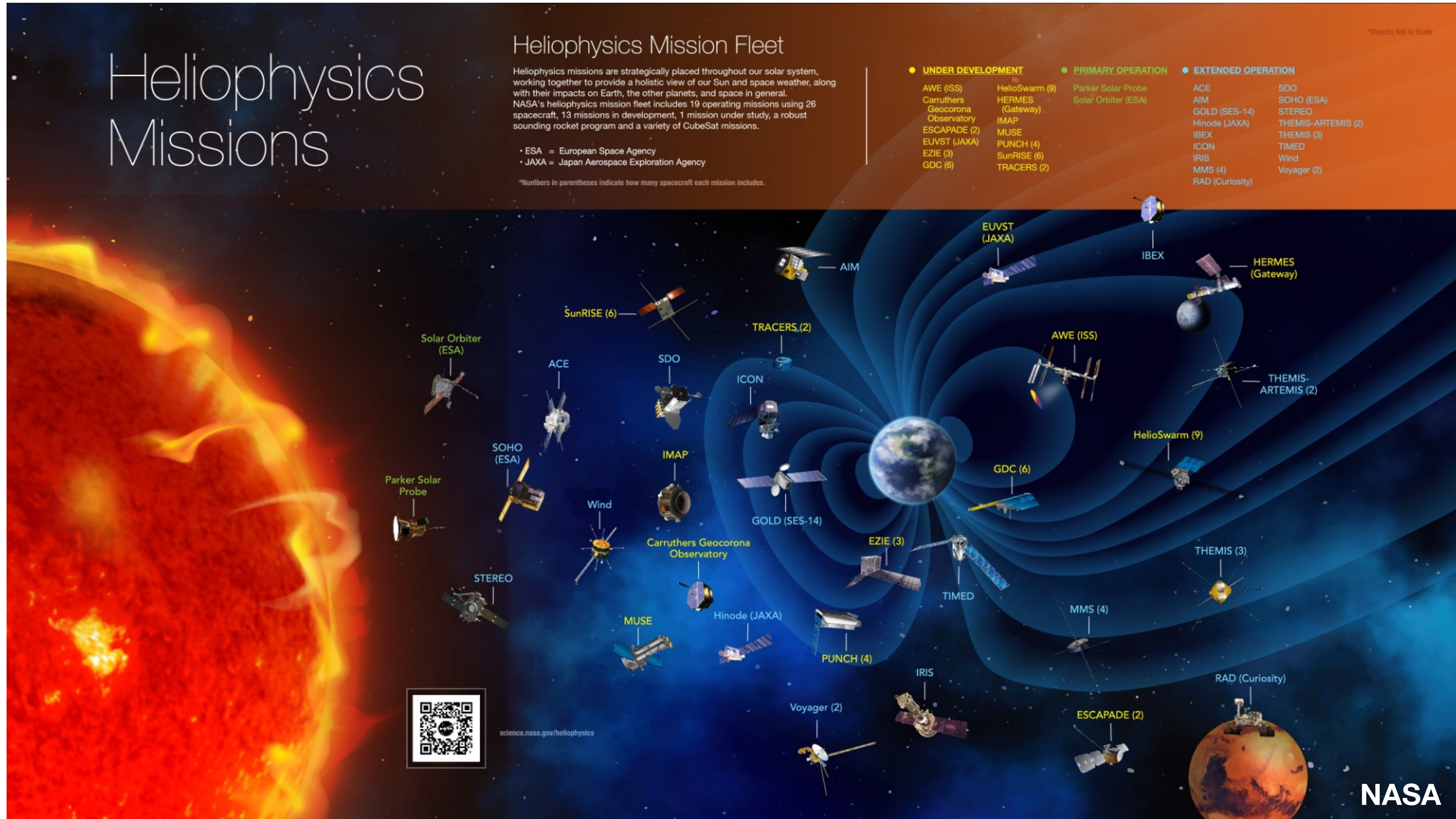
未来与展望：突破Radio Window

- How to detect the frequencies outside the Radio Window?

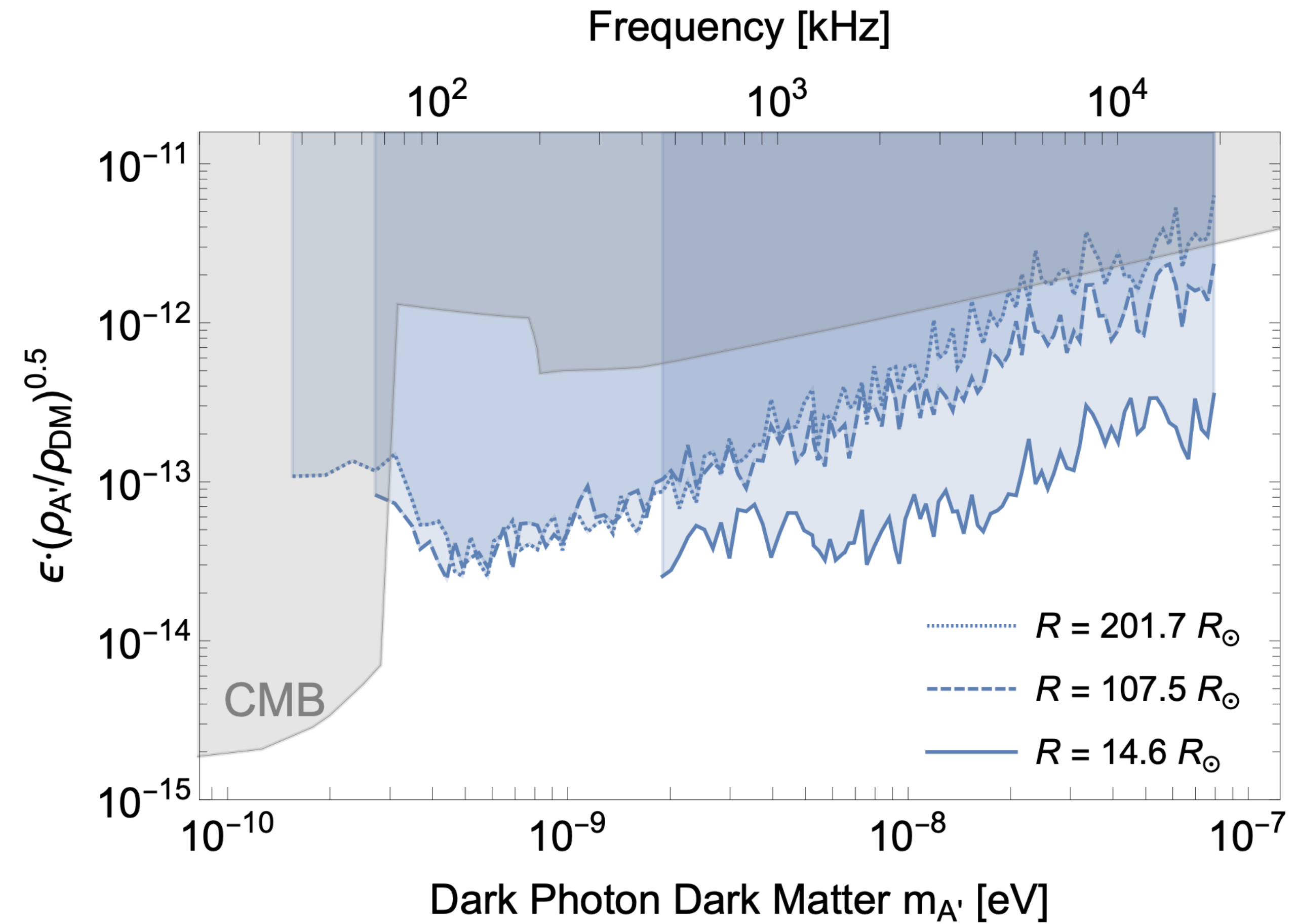
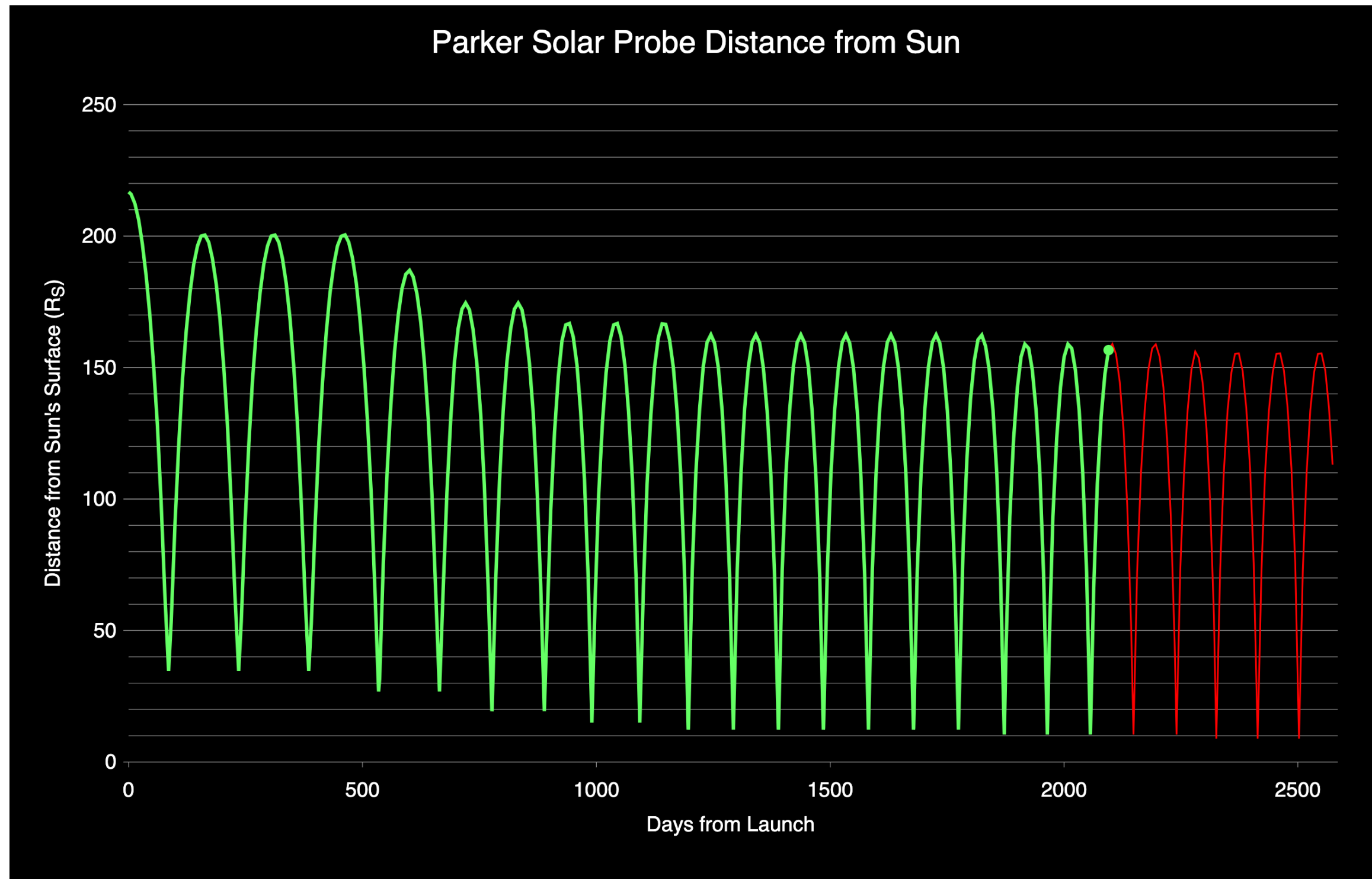


未来与展望：空间天文射电探测

- How to detect the frequencies outside the Radio Window?
- Solar signal: Go Space



Parker Solar Probe preliminary results



Summary

- 超轻玻色型暗物质探测与天文学望远镜观测可以交叉合作
 - 暗光子暗物质可以在太阳等离子体环境中转化为单频光子信号
 - 太阳物理观测数据可以探测超轻暗物质
 - LOFAR、SKA、Daocheng Solar Radio Telescope
 - 暗光子暗物质可以在望远镜反射面或天线阵列转化为单频信号
 - 观测空白天区可以探测超轻暗物质
 - FAST、LOFAR
- 未来可以通过太空射电望远镜探测射电窗口以外的质量区间

