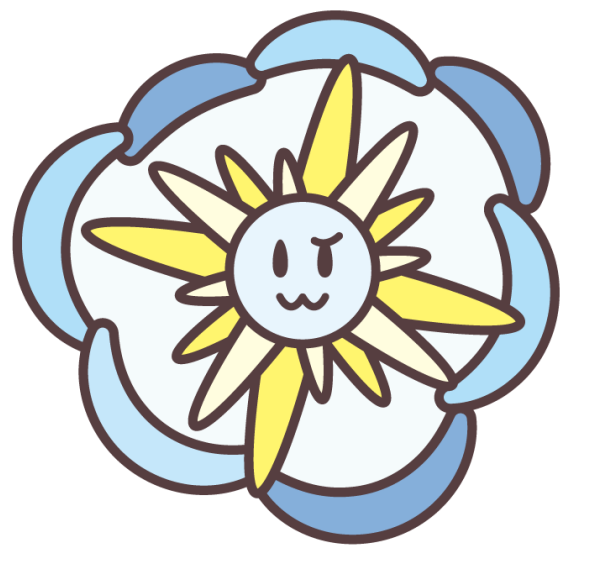


Diagnosing the origin of dense circumstellar material in a multi-energy neutrino astronomical approach

TOHOKU
UNIVERSITY

Yosuke ASHIDA (Tohoku U) & Ryo SAWADA (ICRR, U Tokyo)

Based on the work published in **The Astrophysical Journal 982, 93 (2025)**

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1. Introduction

□ Origin of dense circumstellar material (CSM)??

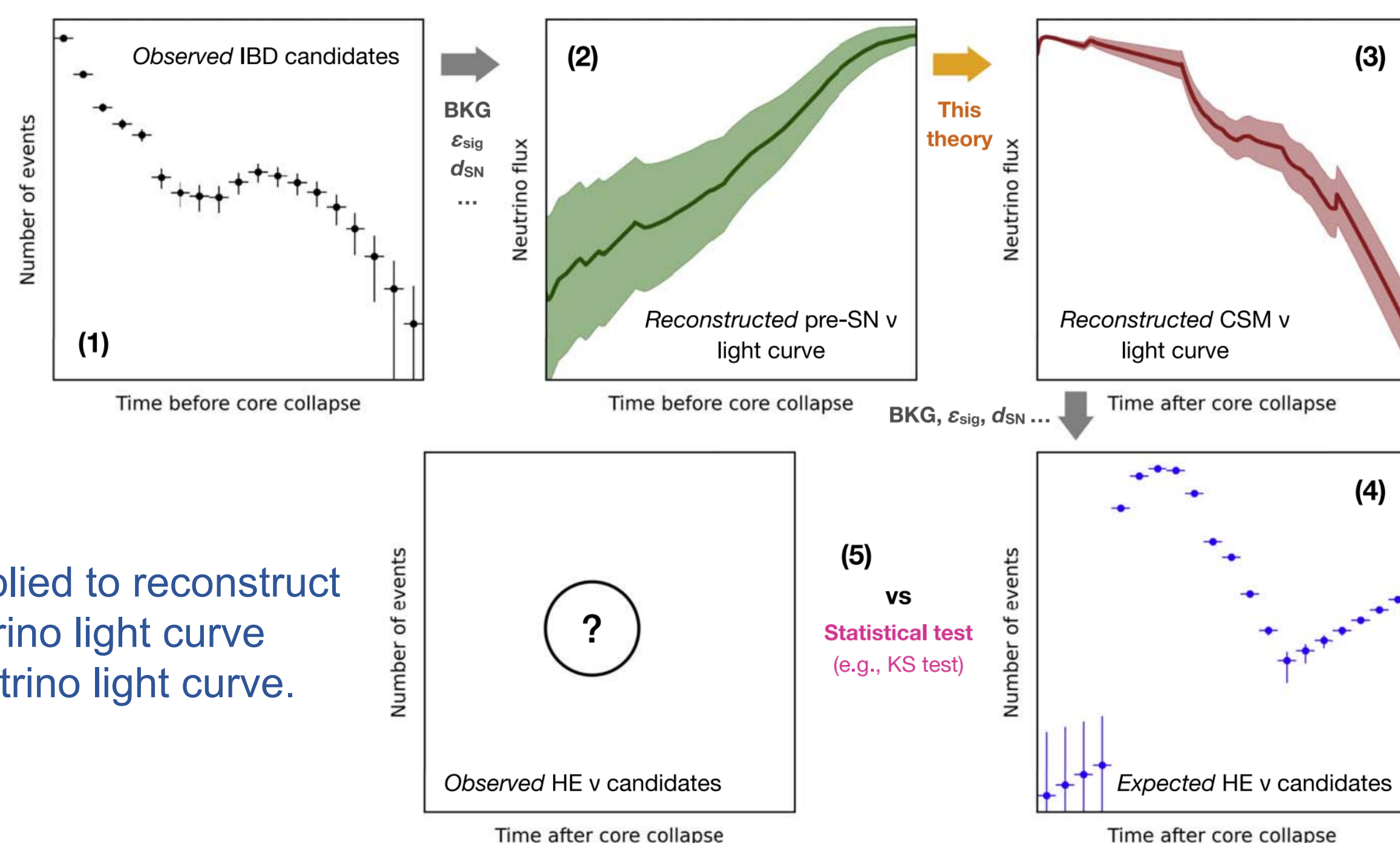
- Core convection
- Binary star process
- Pre-explosion activity

□ Mass loss may be caused by a weaker gravitational binding due to pre-SN neutrino emission?

- Thermal MeV neutrinos** (pre-SN activity);
e.g., C. Kato et al. *ApJ* 848, 48 (2017)
- Non-thermal high-energy neutrinos** (ejecta-CSM interactions);
e.g., K. Murase *PRD* 97, 081301 (2018)

□ Detecting multi-energy neutrinos could solve the mysterious origin of dense CSM.

Approach



3. Demonstration

□ JUNO: Based on *J Phys. G* 43, 030401 (2016); 1.8~4.0 MeV

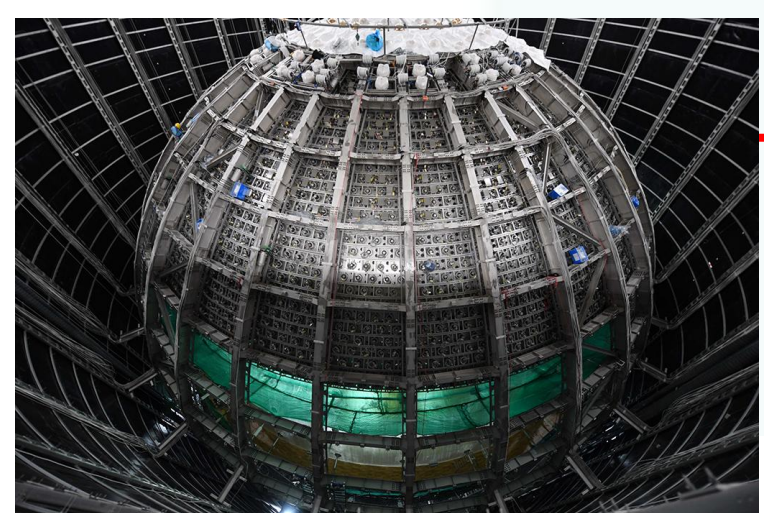
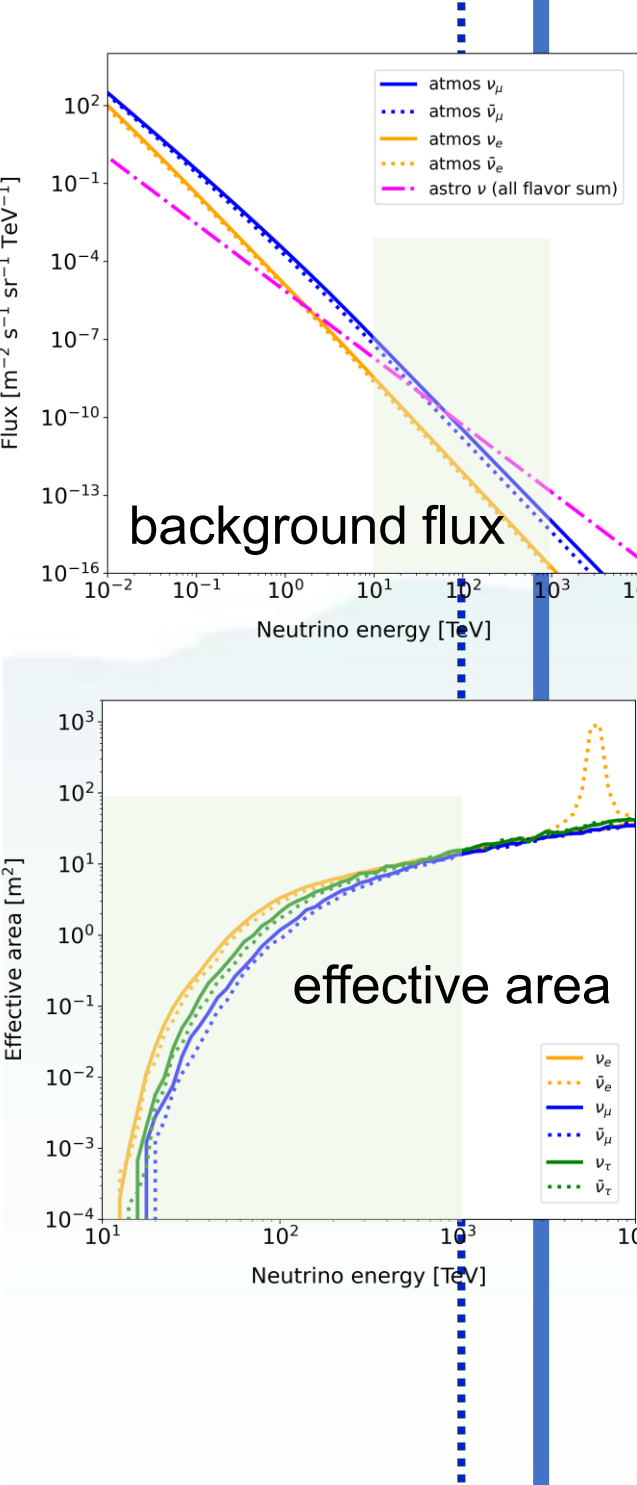
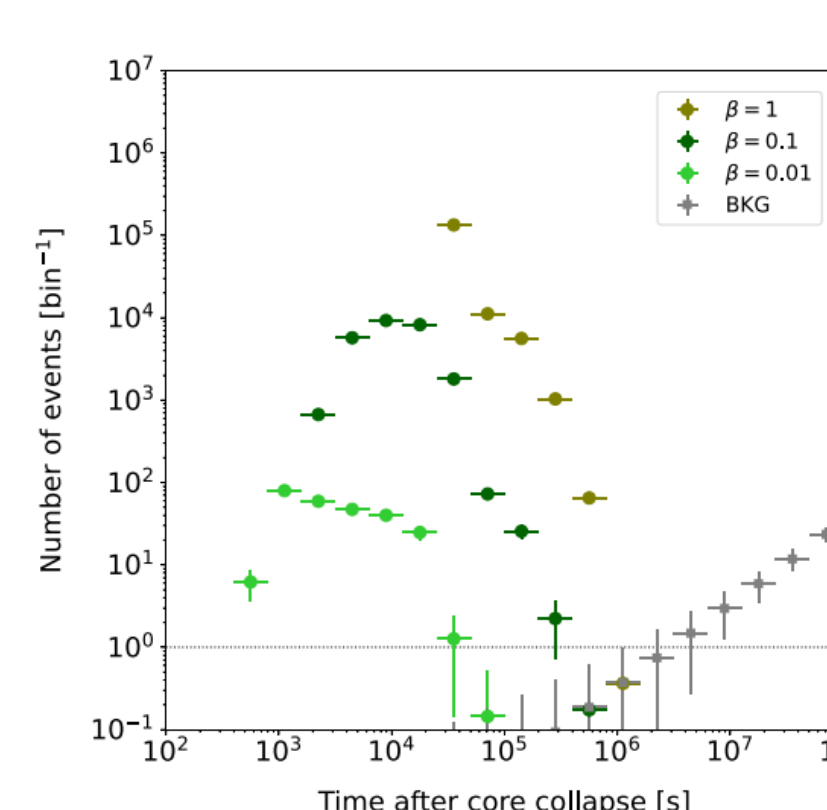
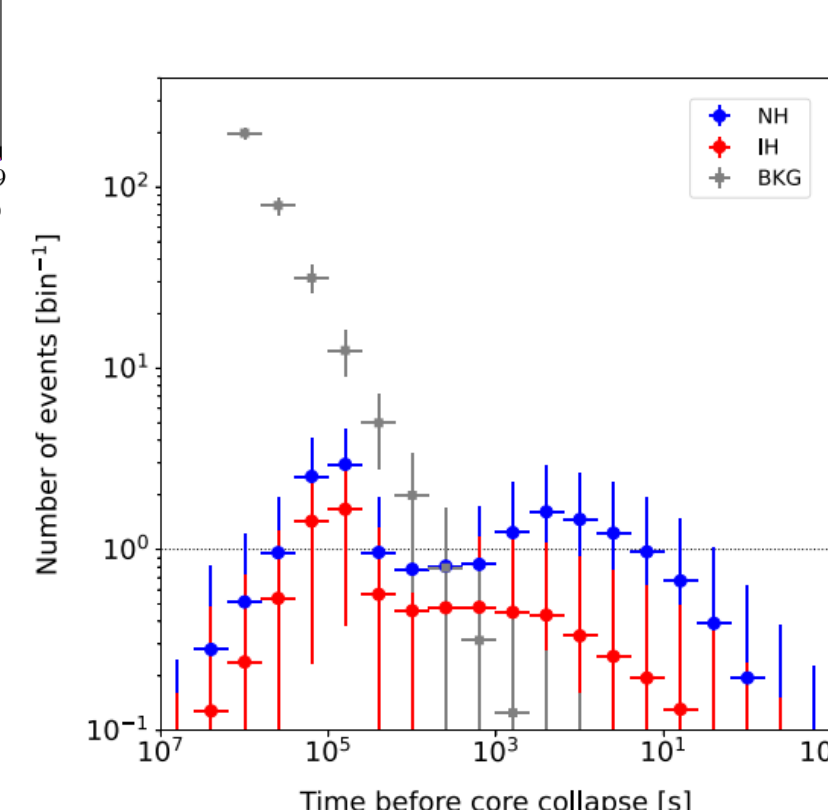
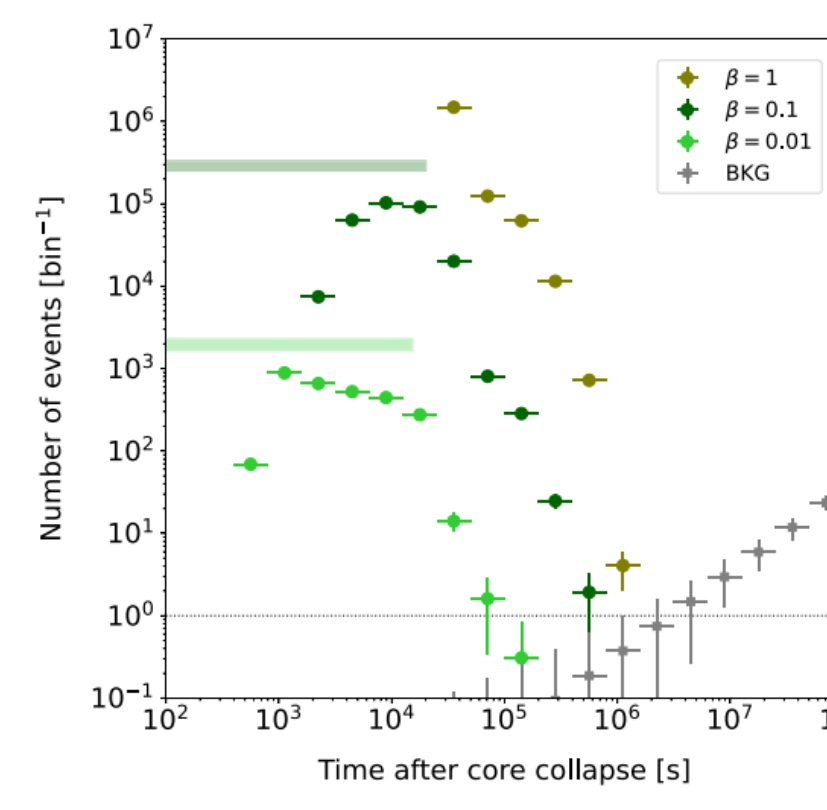
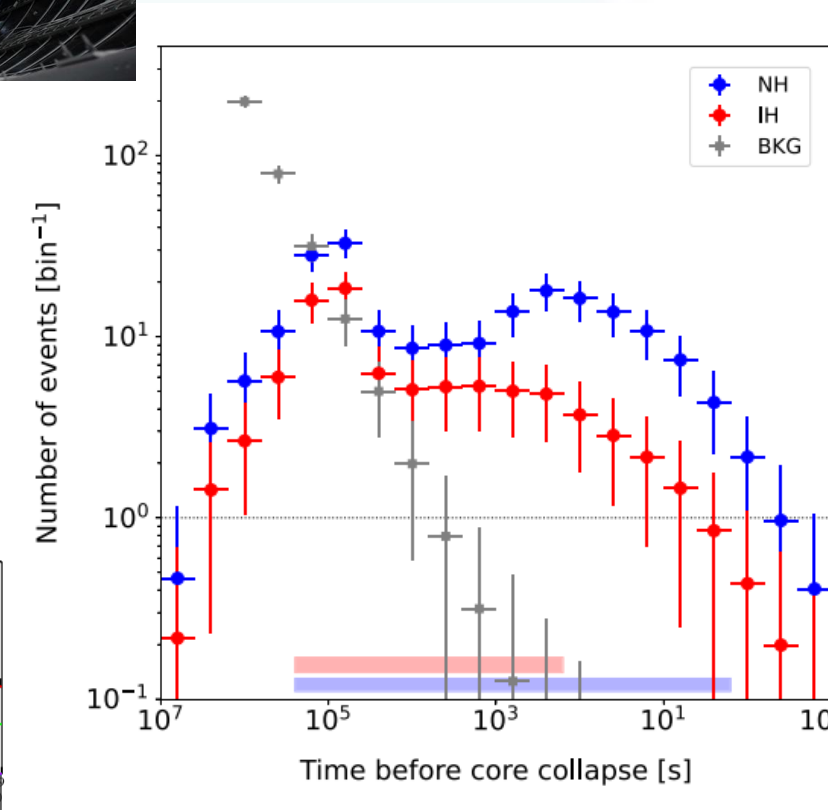
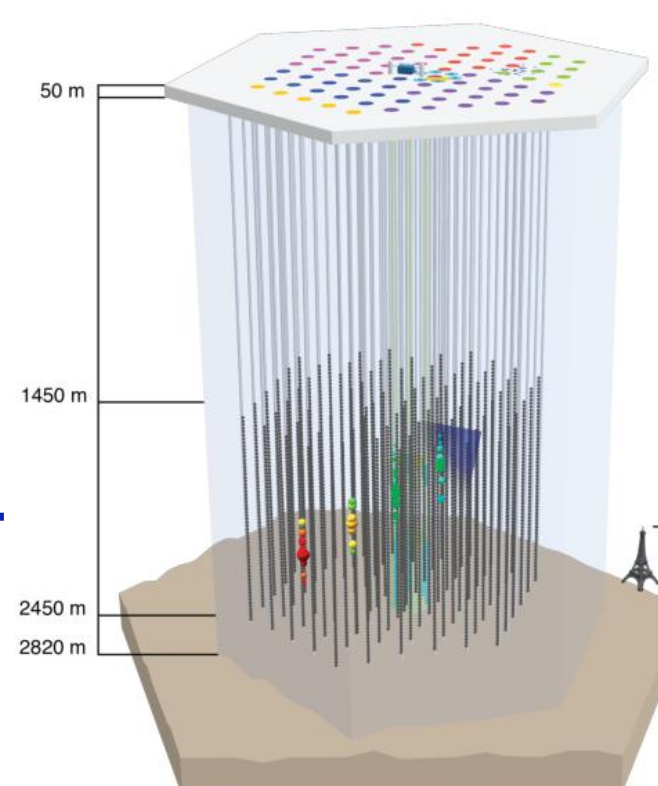
- Detection channel: Inverse Beta Decay (IBD)
- Background: reactor ν , geo ν , spallation, etc ($\sim 18 \text{ day}^{-1}$ in total)

□ IceCube: Based on *HESE* effective area; $10^1 \sim 10^3 \text{ TeV}$

- Detection channel: Deep Inelastic scattering (DIS)
- Background: atmospheric ν (*HAKKM2014*), astrophysical ν (*ESTES*)

□ Synchronized detection

- JUNO: $N_{\text{sig}} / \sqrt{N_{\text{sig}} + N_{\text{bkg}}} > 2$ and $N_{\text{sig}} > 1$
- IceCube: Converted via the time relation

JUNO
(China)IceCube
(South Pole)

4. Conclusion & Discussion

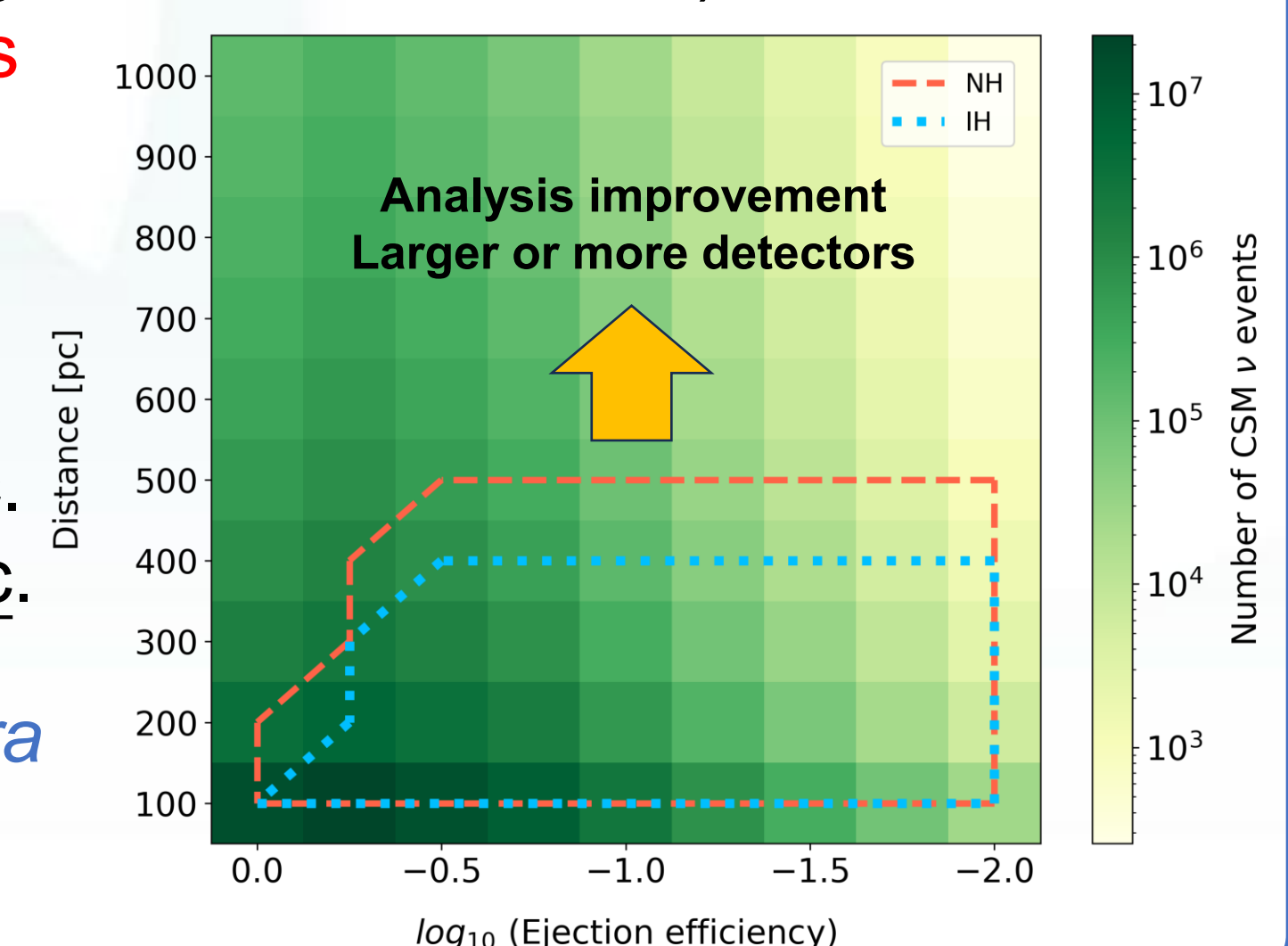
□ The proposed idea is applicable for the range up to **$\sim 500 \text{ pc}$** , which could be extended to **$\mathcal{O}(1) \text{ kpc}$** .

- Analysis energy window in JUNO
- Background reduction (especially reactor neutrinos...)
- Combination with other detectors** (e.g., SK-Gd, Hyper-K).

□ Candidate Wolf-Rayet stars

- A few within $\sim 1 \text{ kpc}$; the closest is γ^2 Velorum (WR 11) at $\sim 340 \text{ pc}$.
- Increased to > 10 within a few kpc.

□ Multi-energy neutrino astronomy era is now explored!



□ Mass loss is assumed to originate from the neutrino release, resulting in CSM production. Core mass-loss rate is expressed as,

$$\dot{M}_{\text{core}} = \frac{L_{\text{pre-}\nu}}{c^2} = 6.8 \times 10^{-3} \left(\frac{L_{\text{pre-}\nu}}{10^{11} L_{\odot}} \right) M_{\odot} \text{ yr}^{-1}$$

□ CSM is reconstructed from pre-SN neutrinos. Mass-loss rate from the stellar surface is,

- $\dot{M}_{\text{sur}}(t) \approx \dot{M}_{\text{wind}} + \beta \cdot \frac{L_{\text{pre-}\nu}(t)}{c^2} \propto \rho_{\text{CSM}}(r, t)$
- $\dot{M}_{\text{wind}} = 10^{-6} M_{\odot} \text{ yr}^{-1}$ (steady-state wind mass-loss rate) $\ll \beta \dot{M}_{\text{core}}$
- β : mass-loss efficiency (considered 0.01~1.00)

□ Ejecta-CSM interactions will produce high-energy neutrinos.

- Roughly, $\rho_{\text{CSM}}(r, t) \propto \frac{L_{\text{pre-}\nu}(t_{\text{CSM}})}{r^2}$

$$t_{\text{CSM}} = t - \frac{r - R_{\text{sur}}}{v_{\text{CSM}}}, \text{ where } v_{\text{CSM}} = v_{\text{esc}} = \sqrt{\frac{2GM_{\text{sur}}}{R_{\text{sur}}}}$$

□ The 'onset' time depends on CSM density.

- When CSM is *shallow*, the collisionless shock wave is formed right after shock wave leaves a star.
- When CSM is *dense*, timescale of the collisionless shock wave formation follows the photon escape time.

