

Sub-keV dark photon search with S2-only data in PandaX4T

Shuaijie Li

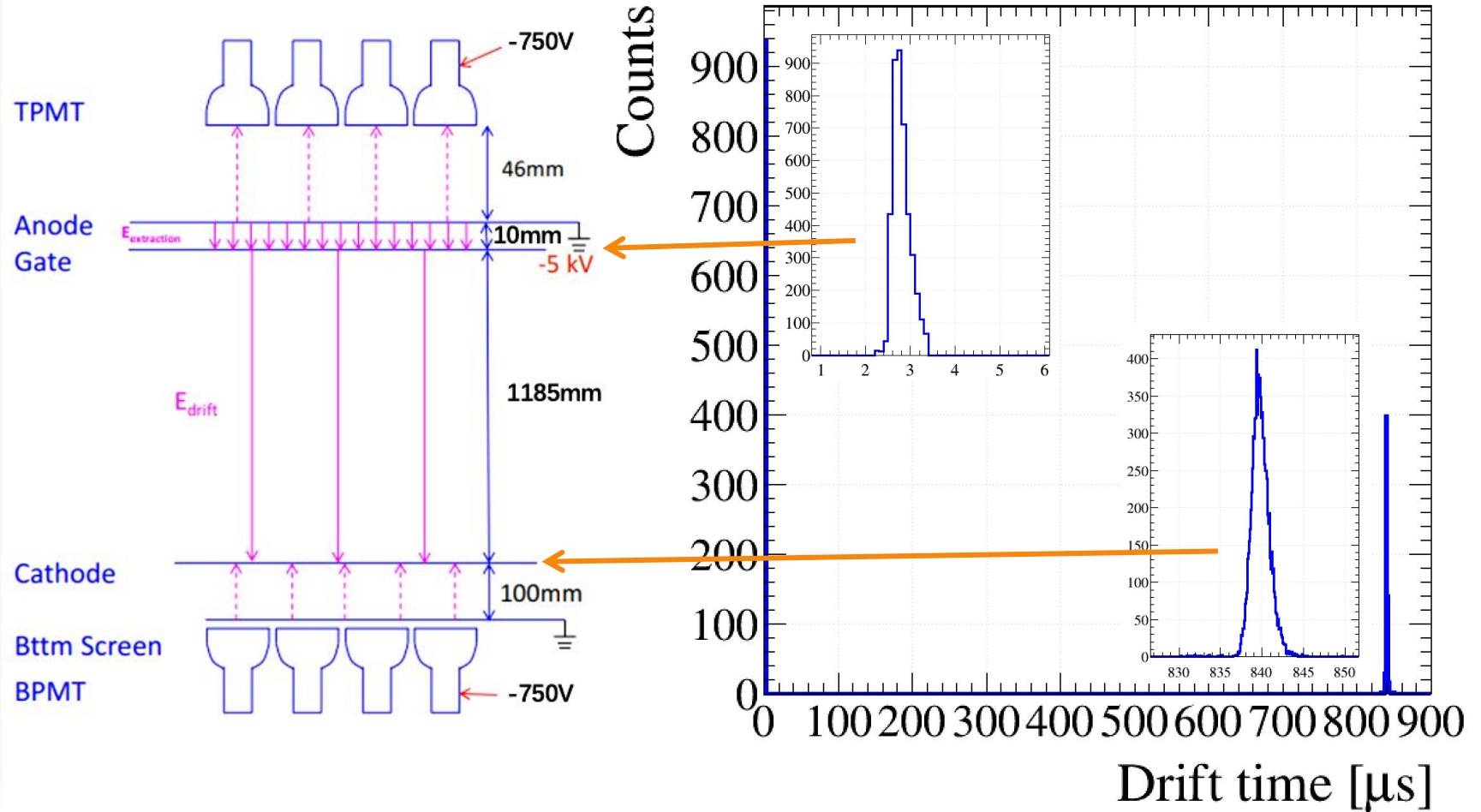
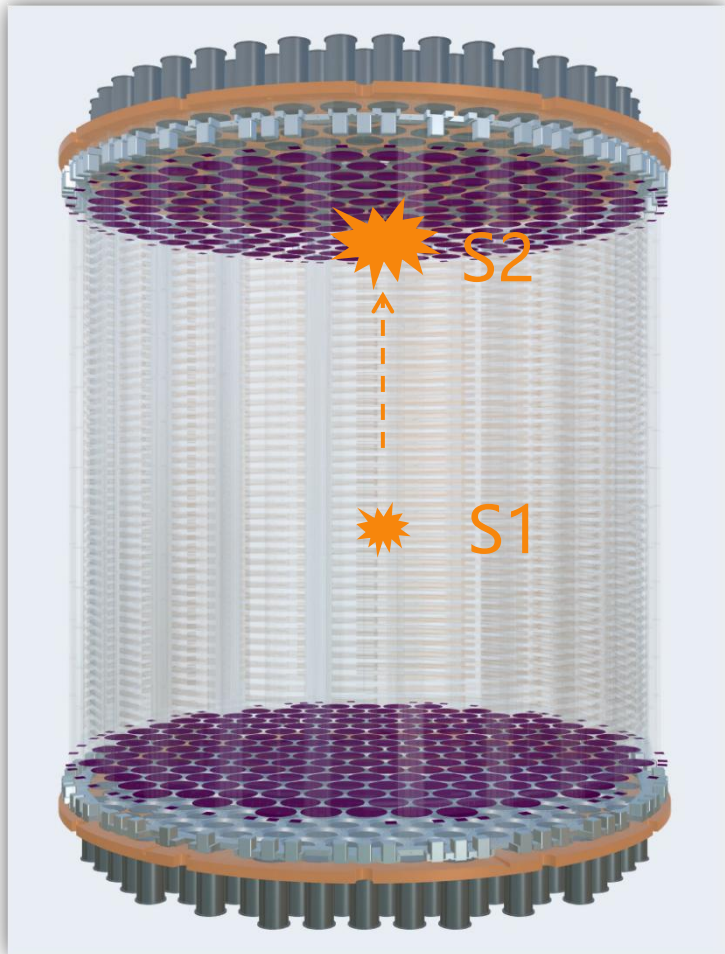
Yalong Hydro (CJPL) / SJTU

On behalf of the low energy group

August, 2025

PandaX-4T experiment

➤ Dual-phase Time Projection Chamber



Dark photons as DM candidate

- DPs are the gauge bosons of dark U(1), which kinetically mix (κ) with the SM photons. Since κ is a free parameter, DPs < 1 MeV are possible.
- Dark photons produce a monochromatic energy signal by kinetically mixing with the SM photons. (or ALP, B–L Gauge Boson...)

$$R_{\text{DPs}} = \frac{4.7 \times 10^{23}}{A} \kappa^2 \left(\frac{\text{keV}/c^2}{m_V} \right) \left(\frac{\sigma_{pe}}{\text{b}} \right) \text{d}^{-1} \text{kg}^{-1}$$

κ : kinetic mixing parameter

σ_{pe} : photo-electric cross section

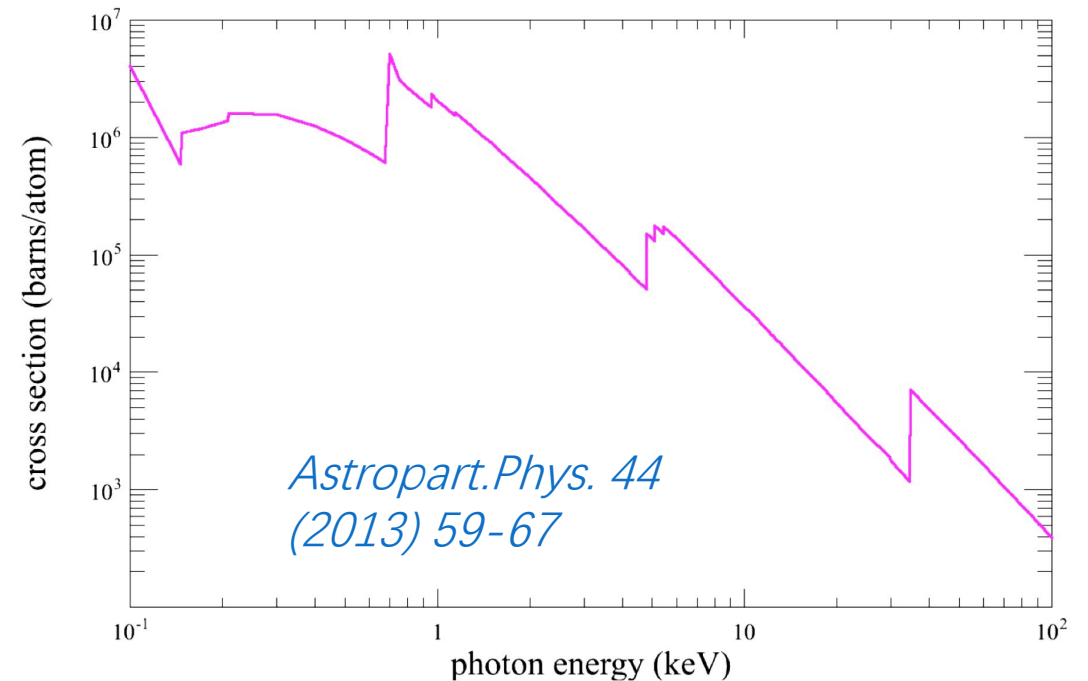
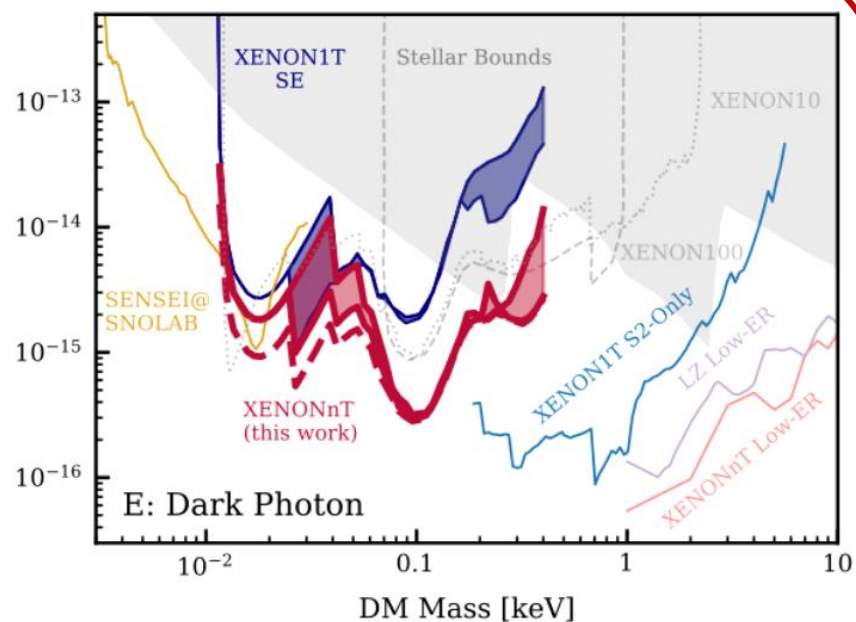


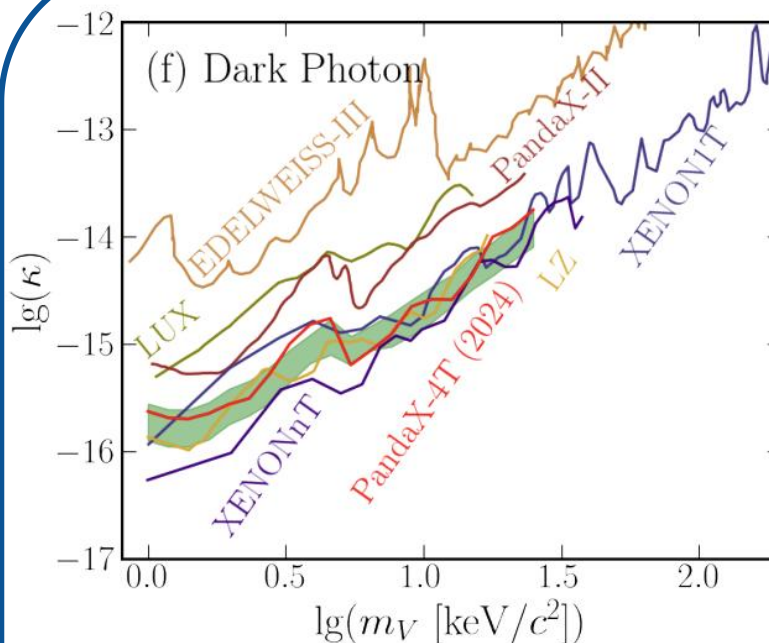
Figure 1: Photo-electric cross section on Xenon atom [27].

Dark photons results

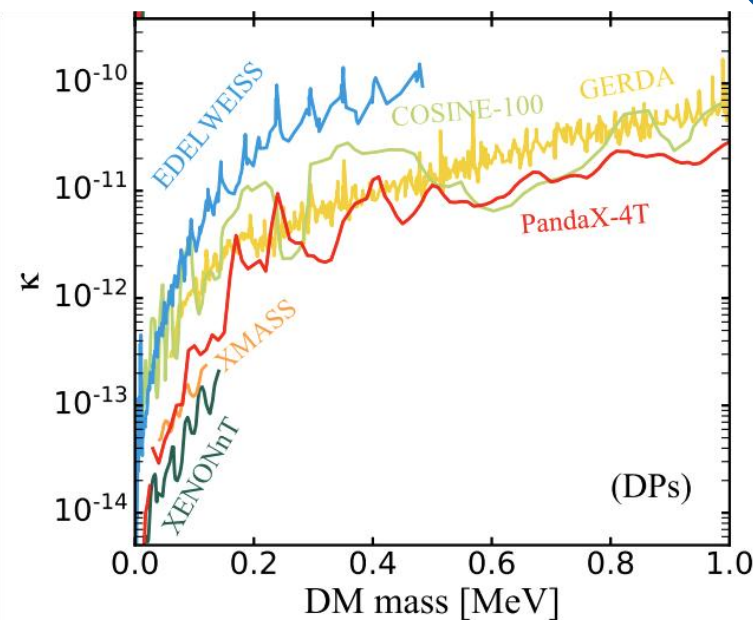
- S2-only channel analysis in LDM and boron-8 neutrino decreased the low energy threshold to < 0.1 keV



PHYSICAL REVIEW LETTERS 134, 161004
(2025)



PHYSICAL REVIEW LETTERS 134, 041001
(2025)



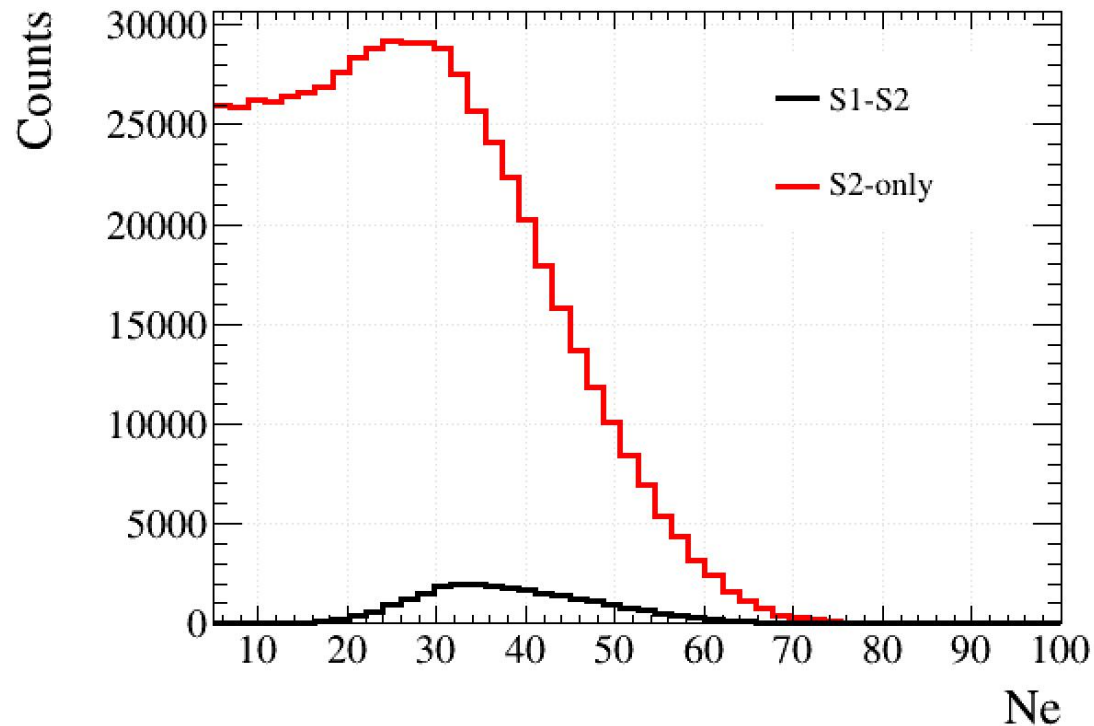
PHYSICAL REVIEW LETTERS 134, 071004
(2025)

Sub-keV DPs searching in PandaX-4T

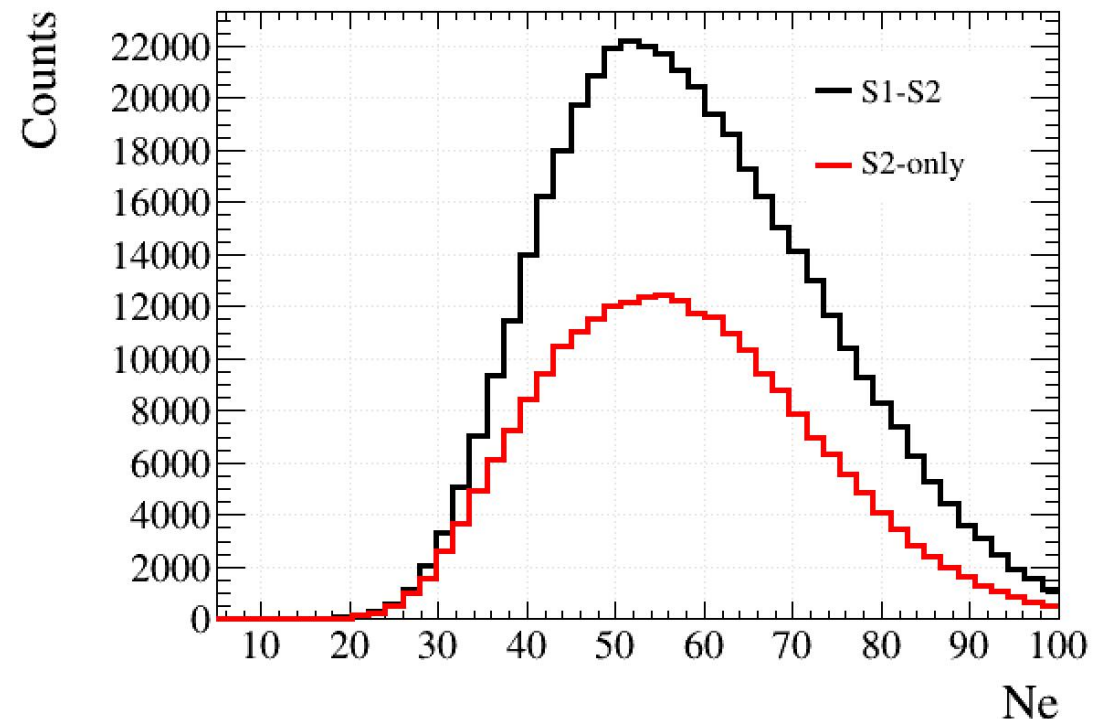
ROI for Sub-keV Dark Photons

- low efficiency of S1 makes the energy threshold to 1 keV for S1-S2 evts
- S2-only data are dominant below 1 keV

flat ER nest simulation (0-1 keV)



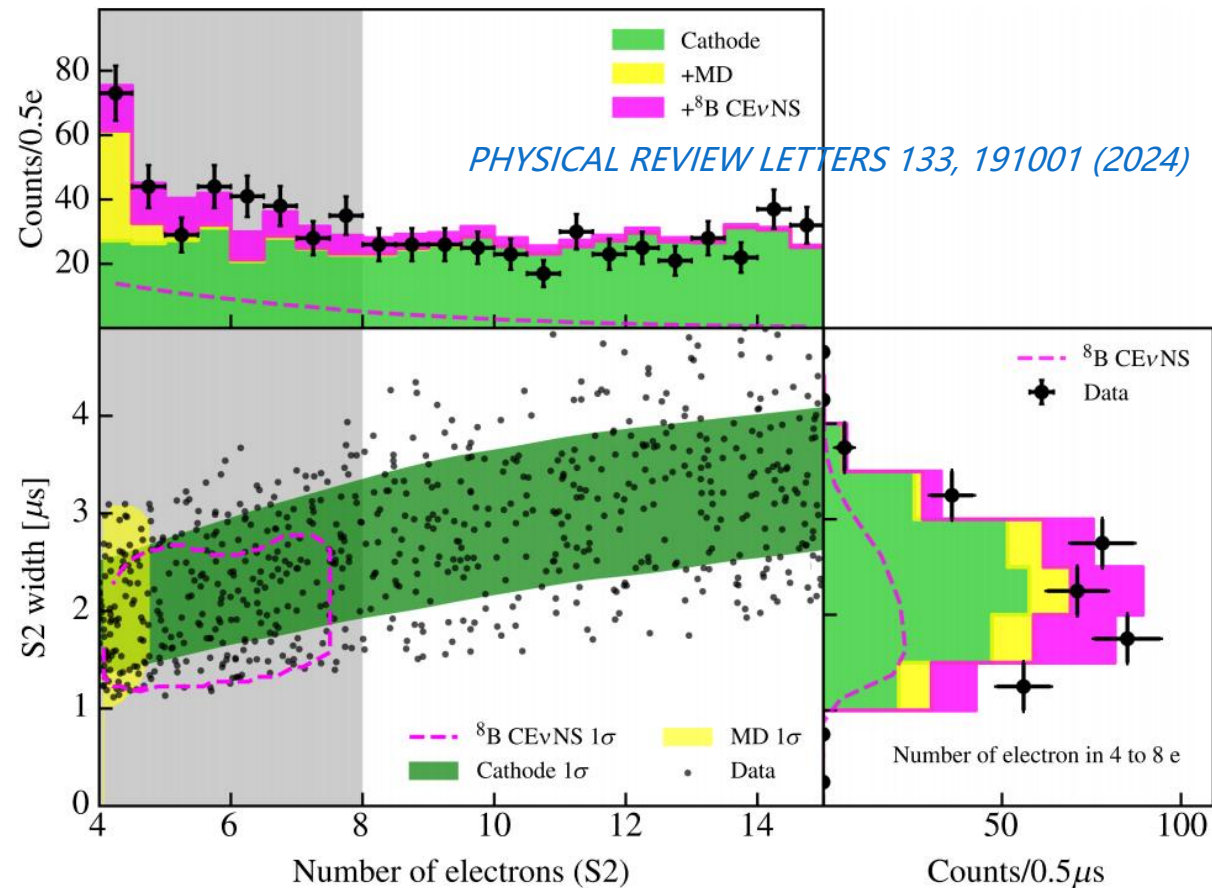
flat ER nest simulation (1-2 keV)



Use 65e as the upper edge of ROI to search Sub-keV DPs

ROI for Sub-keV Dark Photons

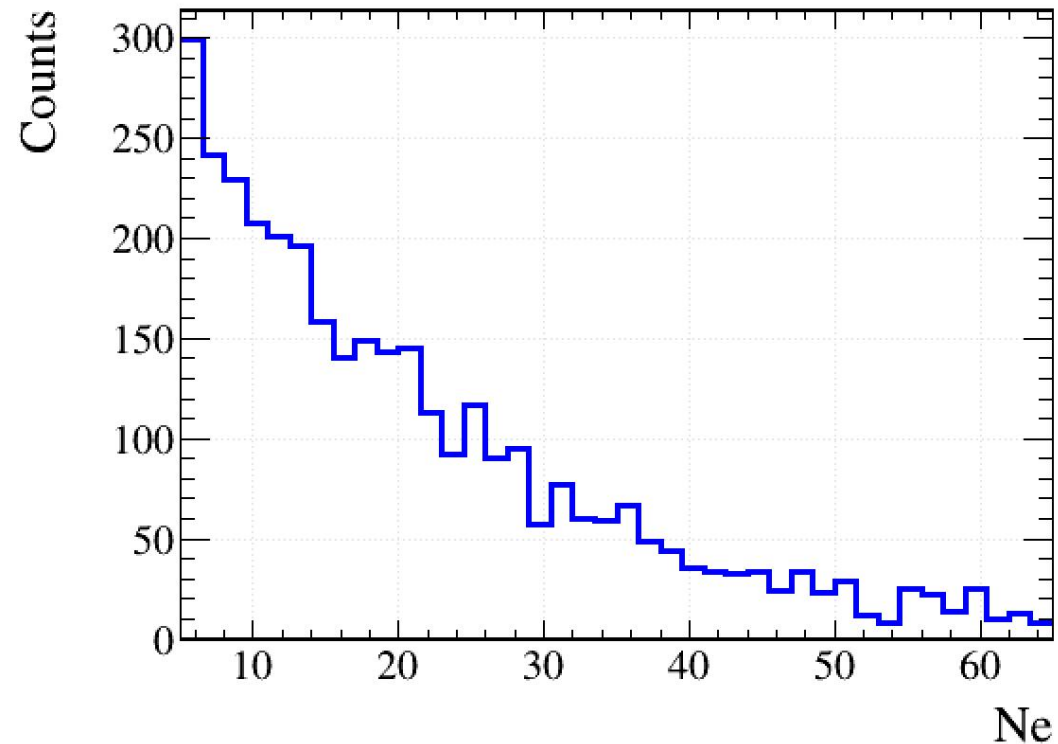
➤ <5 e S2-only data are dominant by MD background



Use 5e as the lower edge of ROI to search Sub-keV DPs

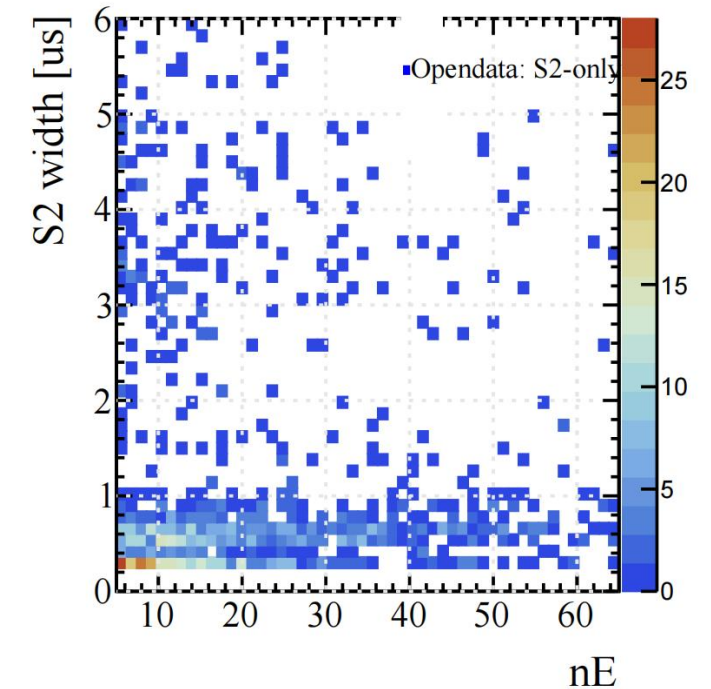
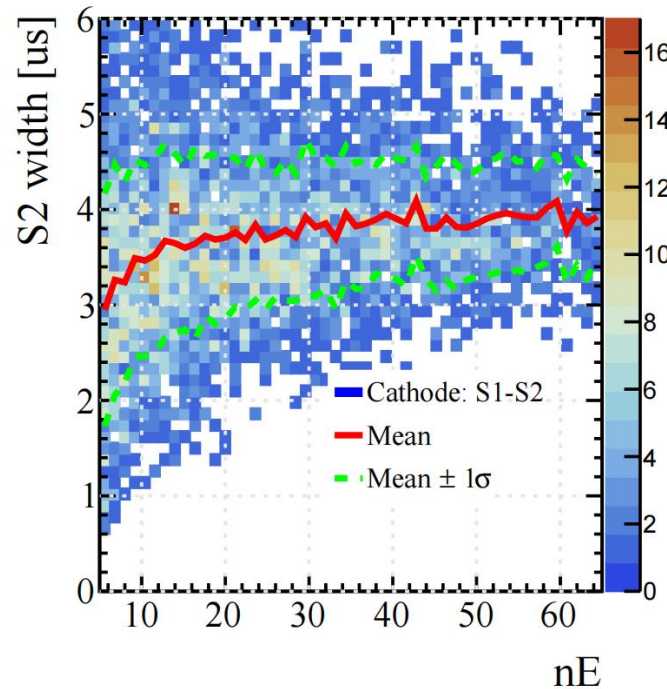
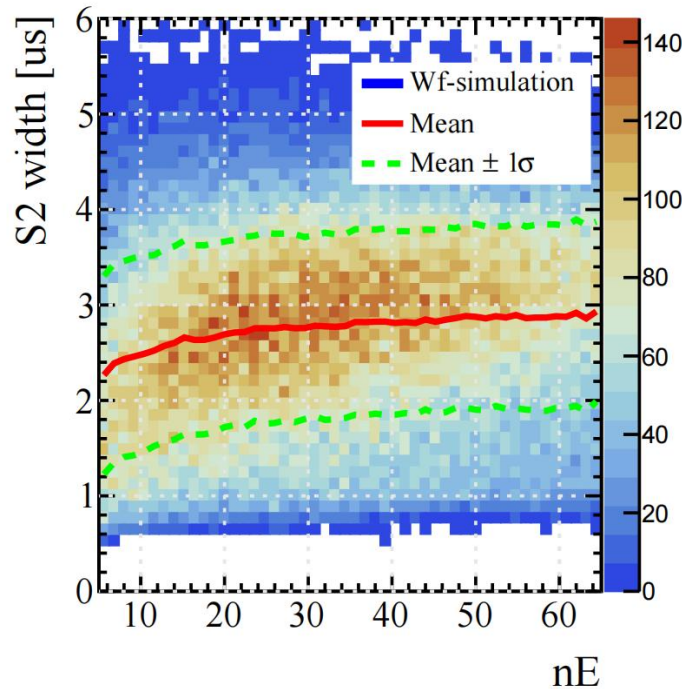
Data

- Randomly opened data in run0 (7.3 d) and run1 (6.4 d)
- all data: run0+run1 1.04 tonne*year: S2-only in ROI, blind
- Waveform simulation: 1-100 e flat S2 simulation
- Neutron calibration data: (including secondary scattered neutron evts)



Selections

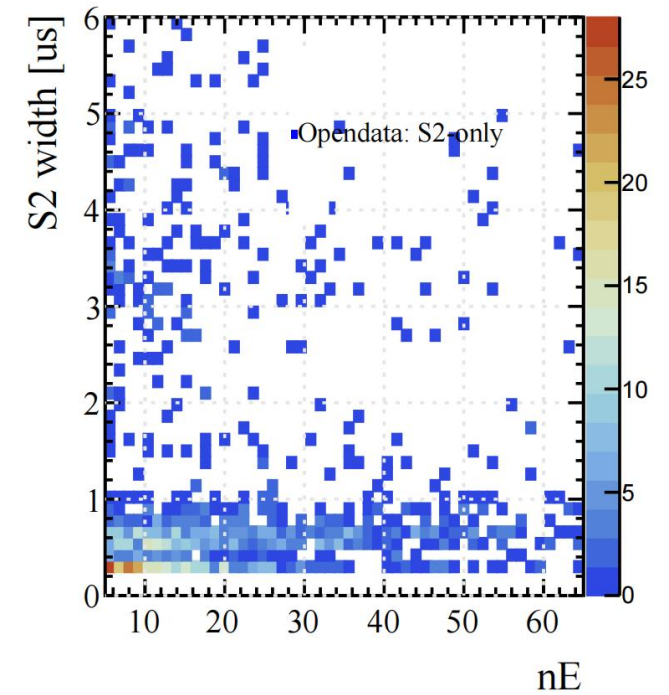
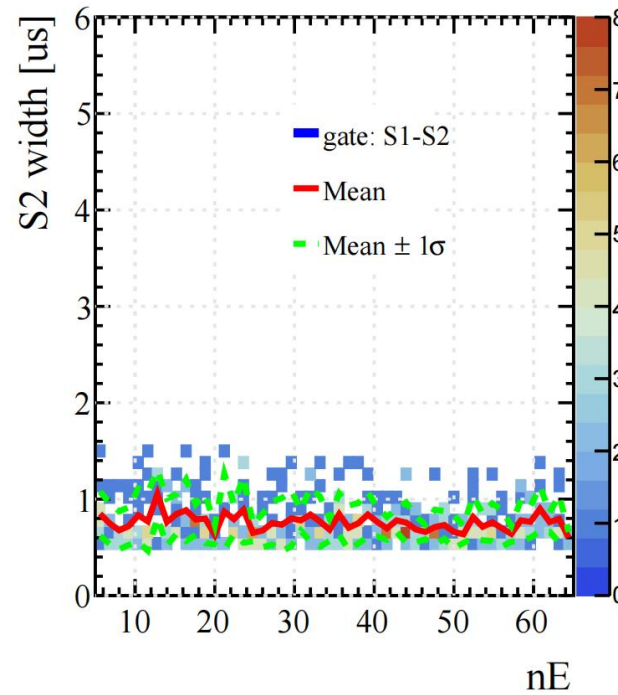
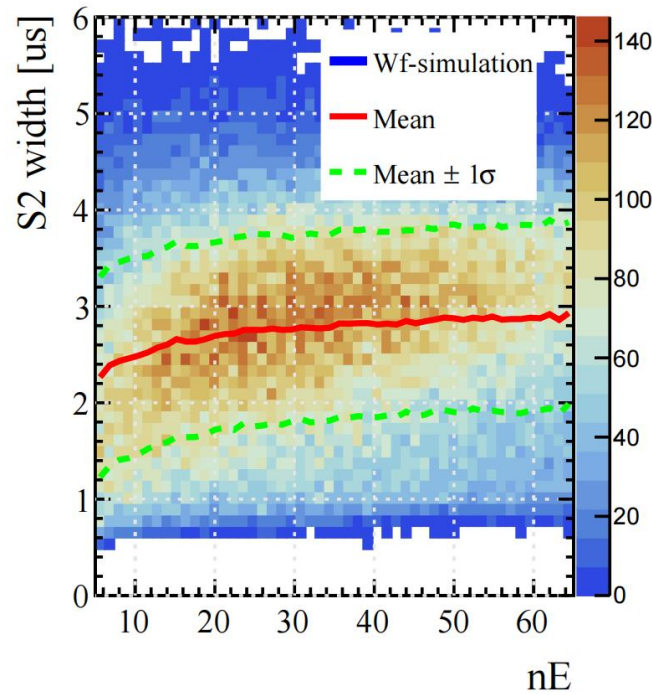
- Inherited the selections from boron-8 analysis
- By using basic cuts, the opendata shows dominance by cathode and gate events in ROI. (run0 as eg.).



Adjust suitable selections for S2-only data above 8 e

Selections

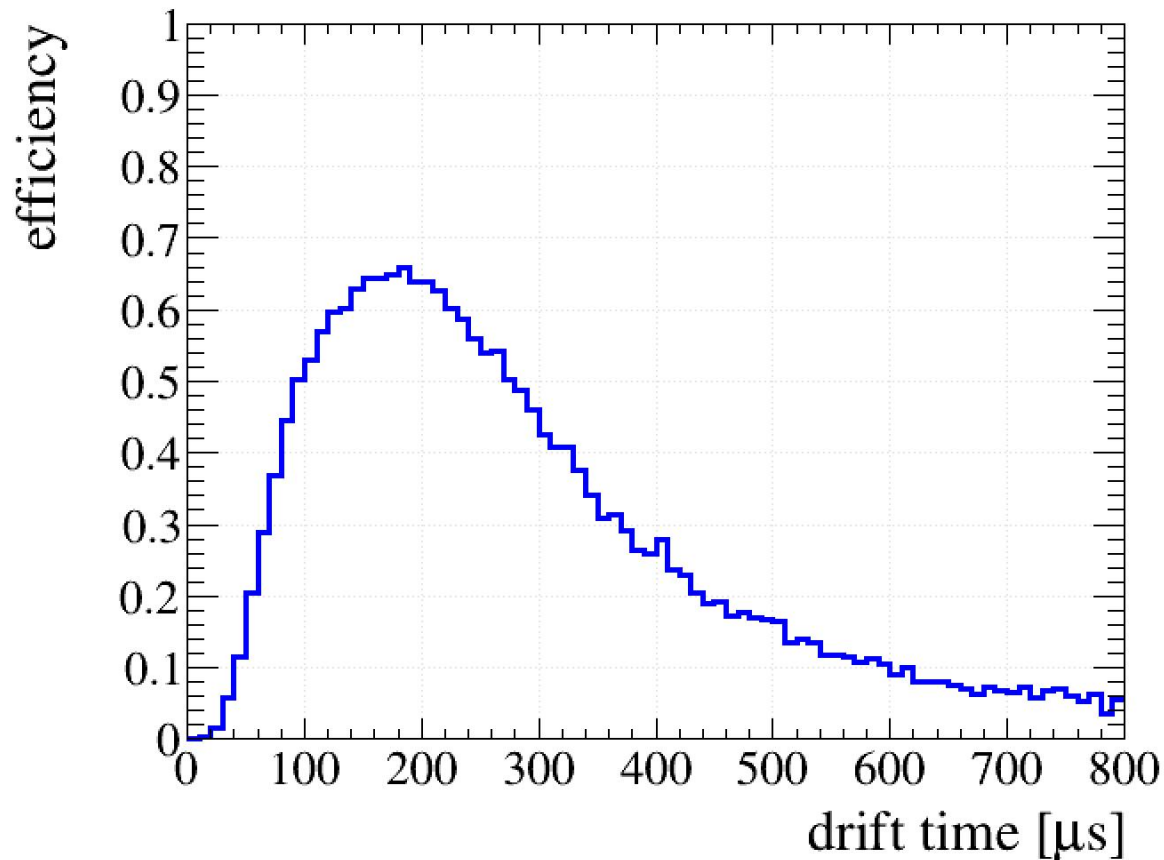
- Inherited the selections from boron-8 analysis
- By using basic cuts, the opendata shows dominance by cathode and gate events in ROI. (run0 as eg.).



Adjust suitable selections to suppress electrode evts

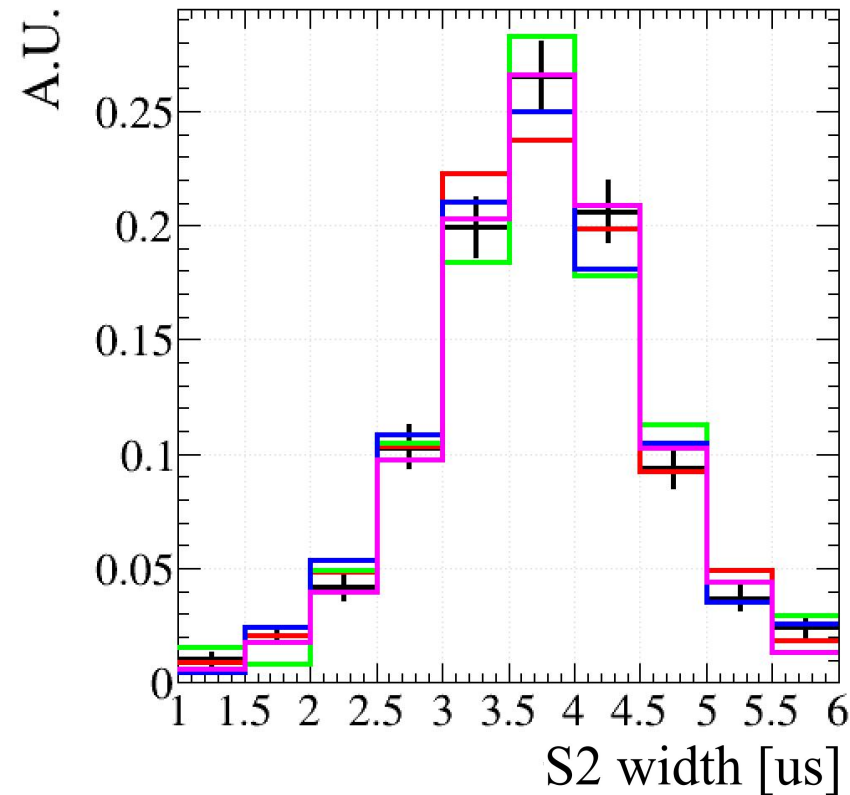
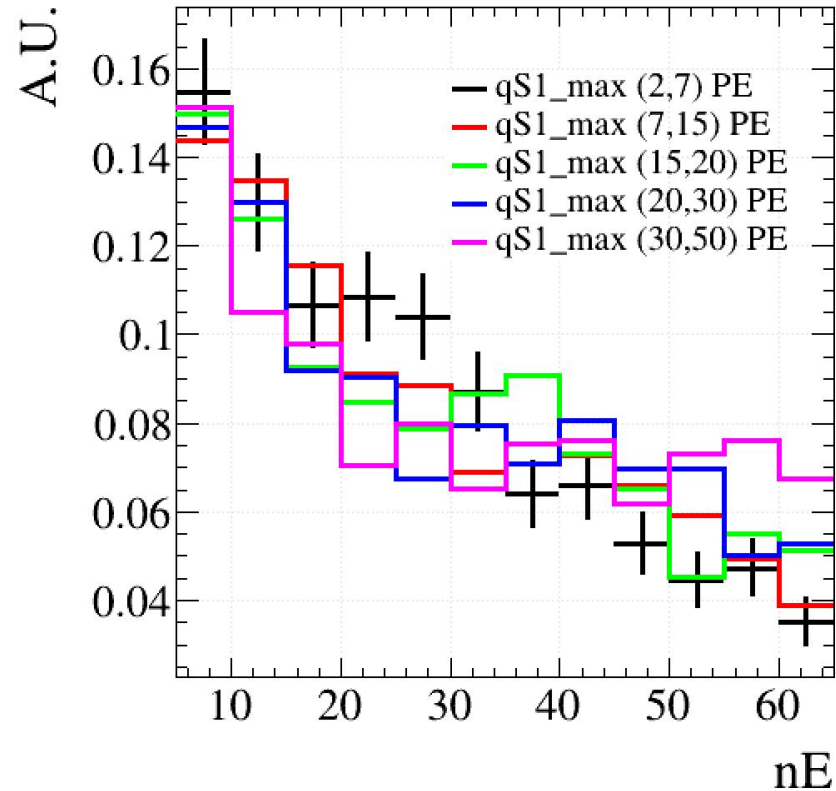
Electrode backgrounds efficiency

- use wf-simulation data to check the drift time efficiency
- The efficiency curve shows strong ability to reduce electrode backgrounds



Background - cathode

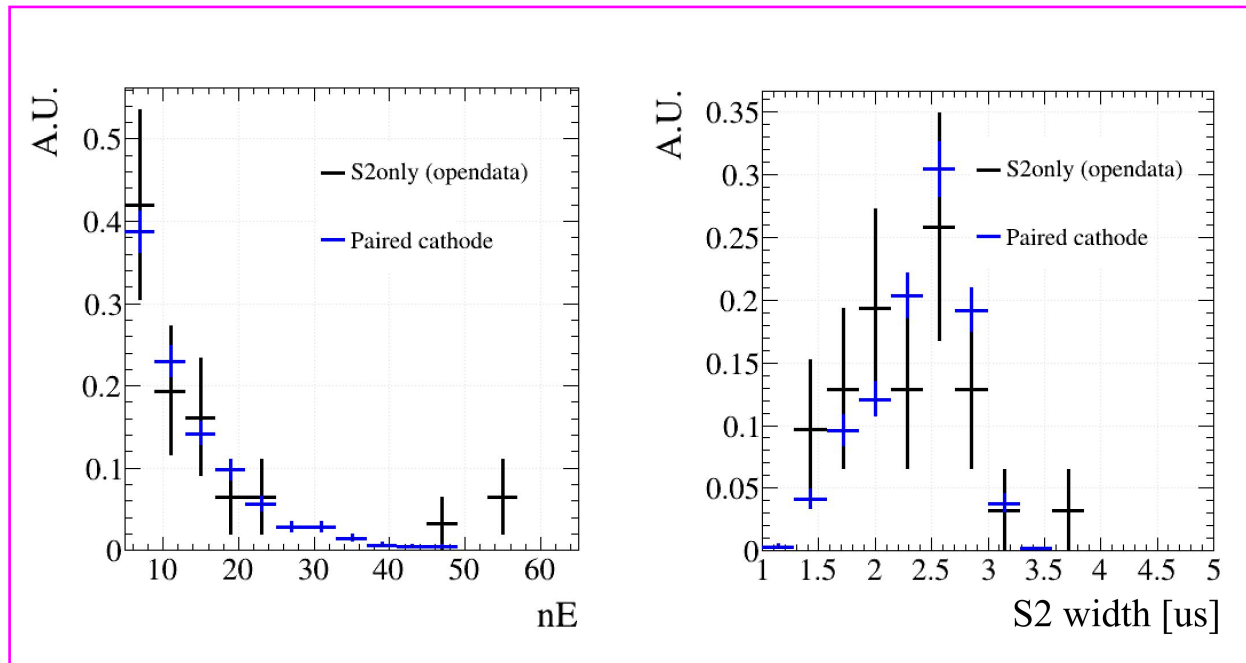
- Paired cathode S2 spectrum is independent with S1



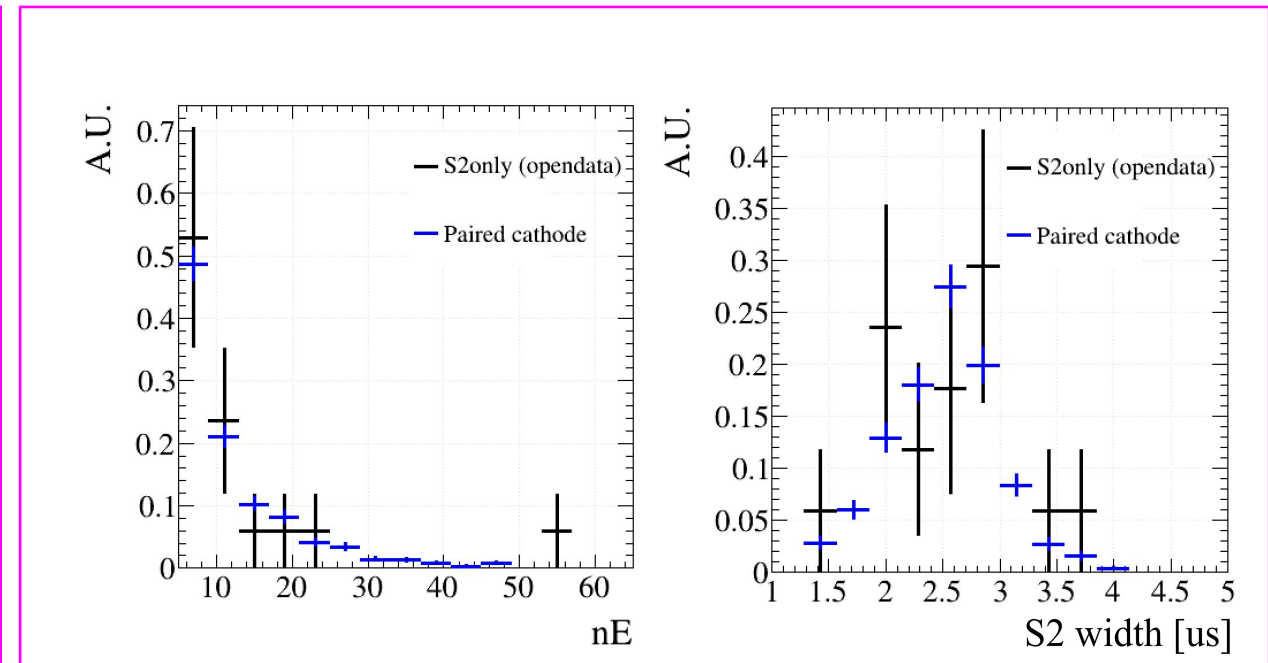
S2-only cathode has the similar features with paried cathode

compare paired cathode evts and open data (all cuts)

- Paired cathode events and S2-only events (opendata) have similar features



run0

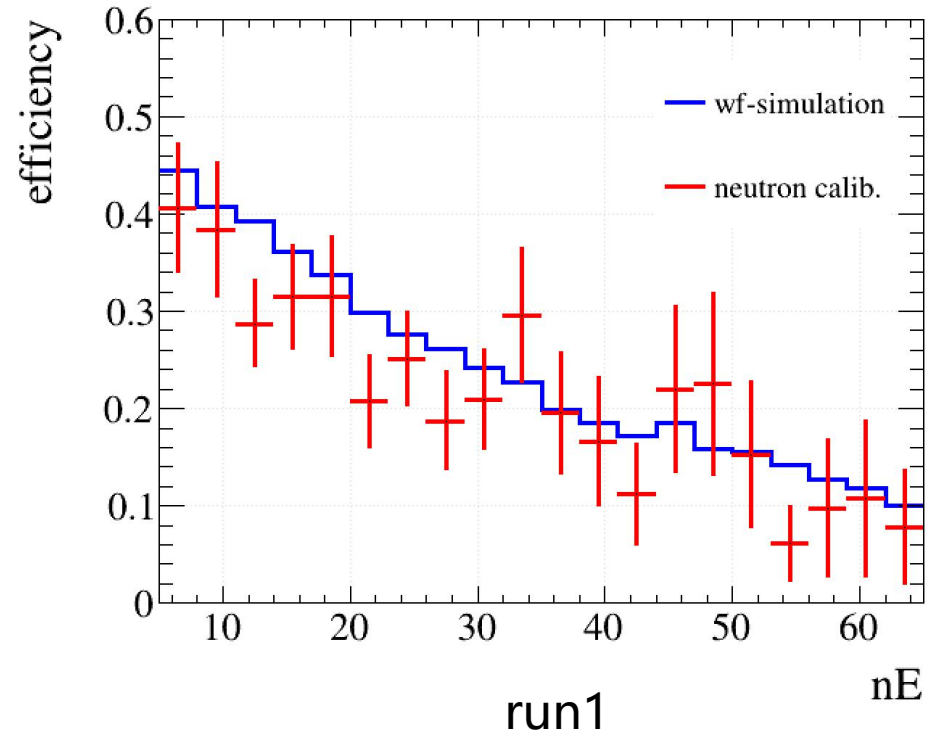
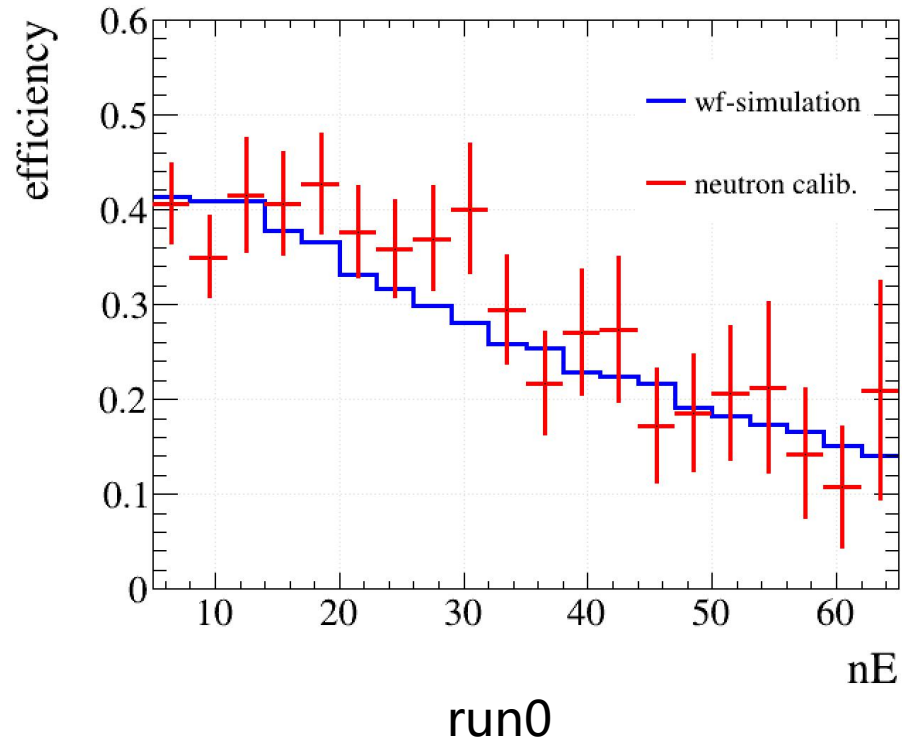


run1

Cathode contribution is non-negligible

Signal efficiency

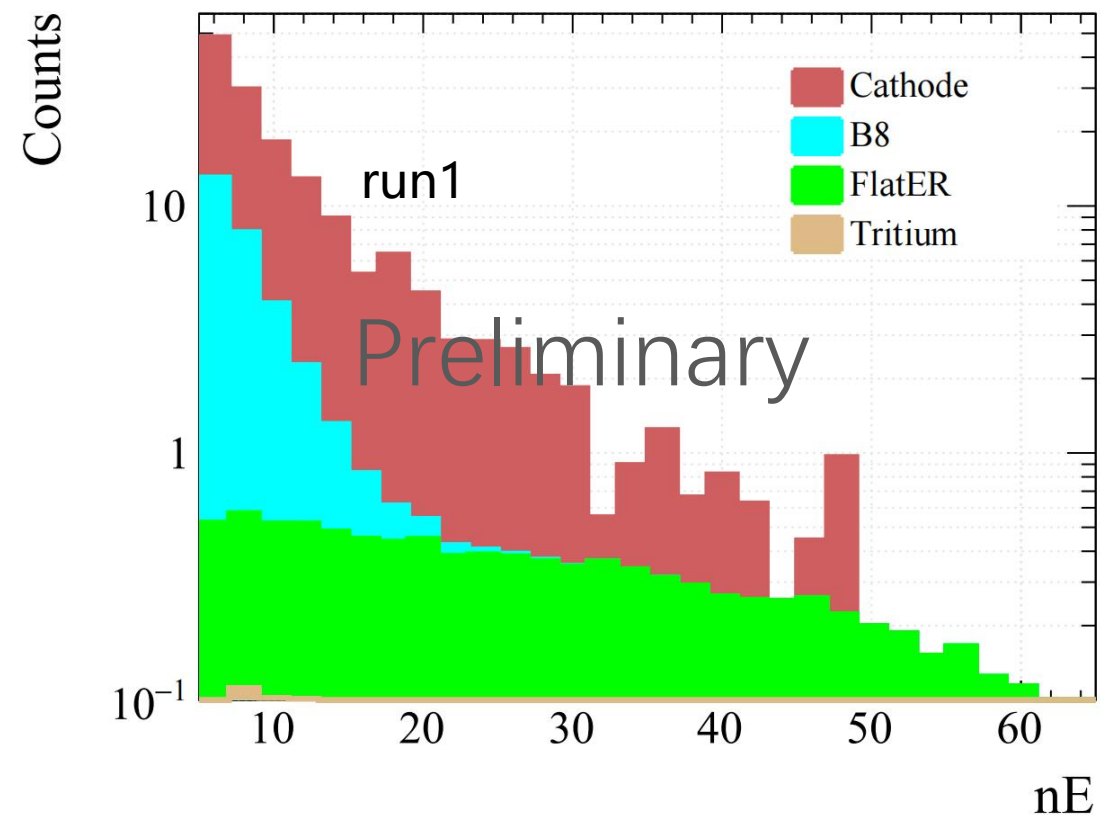
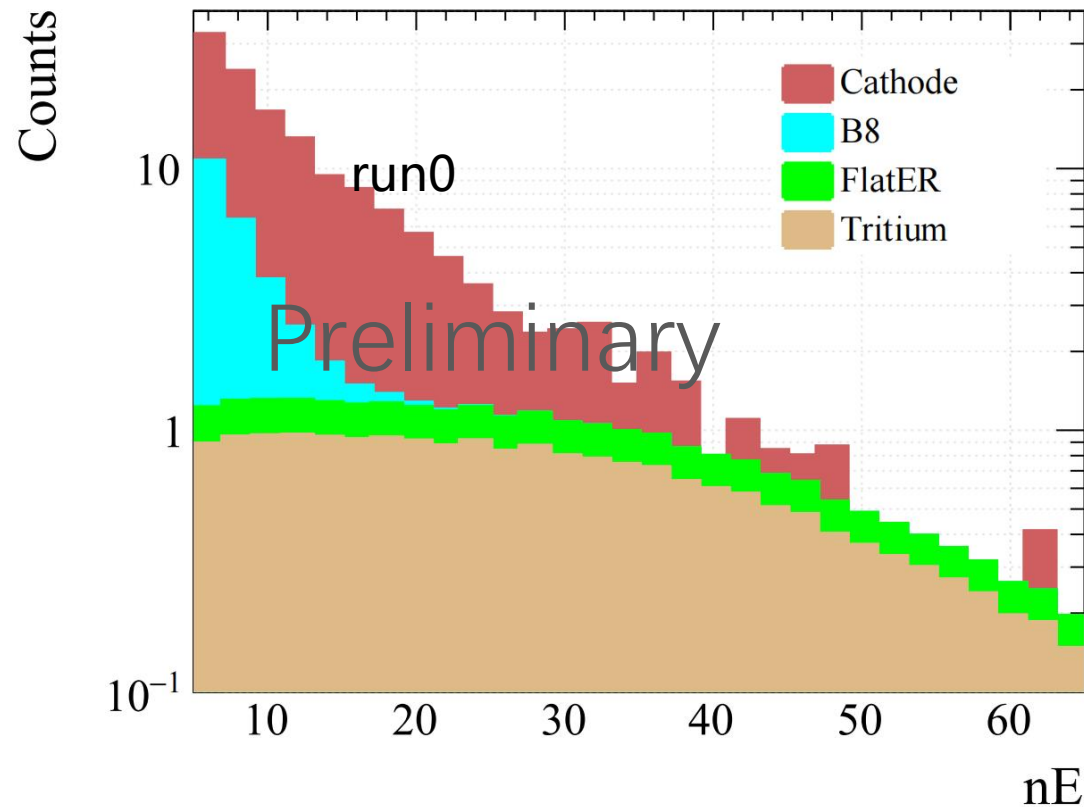
- compare the efficiency performance between wf-simulation and neutron calibration



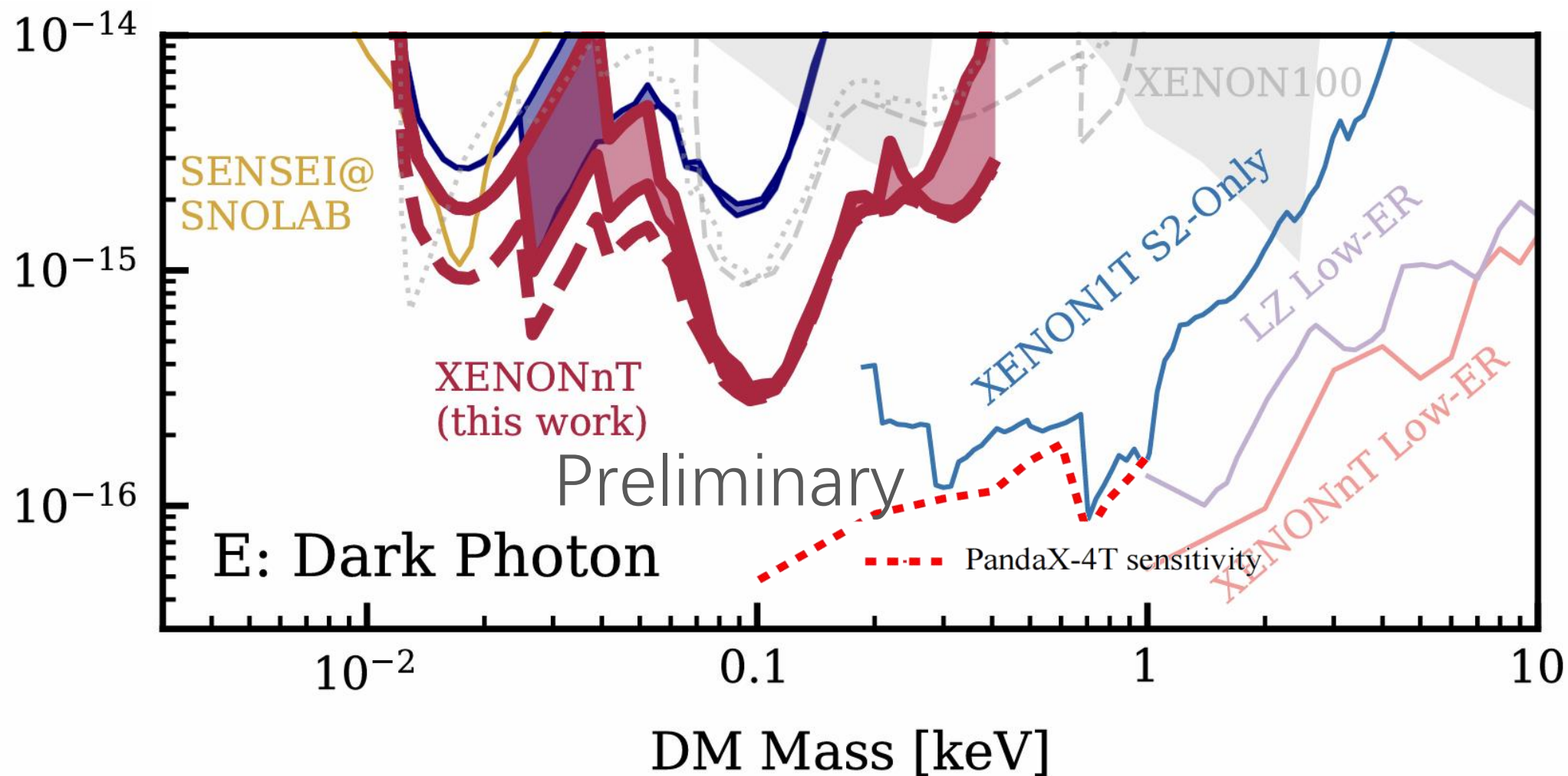
Consistent with wf-simulation and calibration

Expect backgrounds

- Consider the S2-only data consists of cathode, b8, ER events.
- Cathode contributions are scale from large width region in which the cathode S2-only events are dominant.
 - Open all S2-only data in 40-50 e, width in 3-5 us for sensitivity test



Sensitivity



Summary and Plans

- To search Sub-keV dark photons, we expanded the S2-only energy region from 4-8 e to 5-65e.
- The selections have a high ability to veto electrode backgrounds.
- The sensitivity is competitive with existing constraints for Sub-keV dark photons.
- Unblind data...

Thanks!

Back Up

Xenon1T exposure

- s2only use 258.2 days run1 data of Xenon1T, fiducial mass ~ 1 ton, efficiency $\sim 10\%$, exposure: 22 ton*day

Data selection.— We use the main science run (SR1) of XENON1T [5, 6] with a livetime of 258.2 days, after excluding time when the data acquisition was insensitive, the muon veto fired, or a PMT showed excessive pulse rates [7]. Ref. [5] derived a $\sim 4\%$ shorter livetime because it excluded time just after high-energy events. Backgrounds from these periods are mitigated by other methods here.

We report the first dark matter search results from XENON1T, a ~ 2000 -kg-target-mass dual-phase (liquid-gas) xenon time projection chamber in operation at the Laboratori Nazionali del Gran Sasso in Italy and the first ton-scale detector of this kind. The blinded search used 34.2 live days of data acquired between November 2016 and January 2017. Inside the (1042 ± 12) kg fiducial mass and in the $[5, 40]$ keV_{nr} energy range of interest for WIMP dark matter searches, the electronic recoil background was $(1.93 \pm 0.25) \times 10^{-4}$ events/(kg \times day \times keV_{ee}), the lowest ever achieved in such a dark matter detector. A profile likelihood analysis shows that the data is consistent with the background-only hypothesis. We derive the most stringent exclusion limits on the spin-independent WIMP-nucleon interaction cross section for WIMP masses above 10 GeV/c², with a minimum of 7.7×10^{-47} cm² for 35-GeV/c² WIMPs at 90% confidence level.

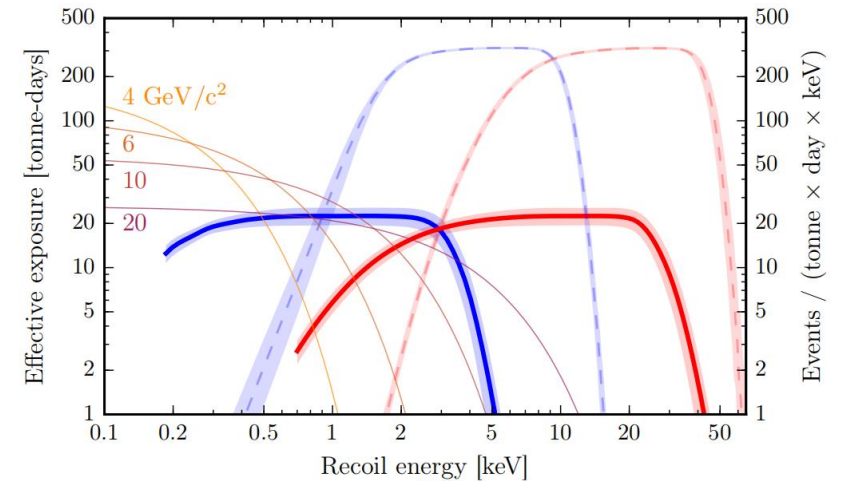


FIG. 1. Effective remaining exposure after event selections for NR (red) and ER (blue) signals of different energies, for $S2 \in [150, 3000]$ PE, on the left y-axis. Dashed lines show the same for XENON1T's main analysis [5], and shaded bands show $\pm 1\sigma$ systematic uncertainties. Thin lines show the expected differential event rate for 4, 6, 10, and 20 GeV/c² spin-independent (SI) DM-nucleus scattering with $\sigma = 10^{-43}$ cm², under the nominal signal model, on the right y-axis.

PandaX4t: 1.04 ton*year * 0.2 \sim 76 ton*day effective exposure

data

- all data are after s2-only recluster algorithm
- waveform simulation and neutron calibration
- all paired events, and s2-only events in open run

Dataset summary

Combine adaptive deadtime, bad files

| Run0 | duration | livetime |
|-------|----------|----------|
| Set1 | 1.955 | 0 |
| Set2 | 13.540 | 10.300 |
| Set3 | 5.534 | 0 |
| Set4 | 37.224 | 26.514 |
| Set5 | 36.607 | 27.787 |
| Total | 94.858 | 64.6 |

Combine fiducial volume, eburst, off pmt

| | run0 | run1 | Total |
|--------------------|------|------|-------|
| FV mass(ton) | 2.78 | 2.16 | / |
| Exposure(ton·year) | 0.49 | 0.55 | 1.04 |

Published s2only: 0.55ton·year



| run0 | livetime | totaltime |
|------|----------|-----------|
| 3026 | 9.26652 | 11.9689 |
| 3031 | 17.2275 | 22.287 |
| 3180 | 4.51674 | 5.75401 |
| 3181 | 3.13397 | 3.97455 |
| 3424 | 12.3921 | 20.3345 |
| 3439 | 10.5874 | 20.3059 |
| 3506 | 16.7903 | 22.3684 |
| 3509 | 12.96 | 17.4442 |
| 3550 | 3.07622 | 23.9055 |
| 3552 | 17.8586 | 23.4861 |
| 3600 | 3.91894 | 5.12163 |
| 3601 | 7.19873 | 15.3361 |
| 3730 | 14.3723 | 18.8436 |
| 3731 | 17.8457 | 23.5222 |
| 3777 | 18.2868 | 23.9892 |
| 3778 | 0.555916 | 0.71985 |
| 3805 | 13.2594 | 17.3507 |
| 3806 | 4.3591 | 5.71446 |
| 3839 | 18.2629 | 24.0268 |
| 3840 | 12.3248 | 16.3861 |
| | 7.3 day | |

| run1 | livetime | totaltime |
|------|----------|-----------|
| 4679 | 9.40012 | 15.0668 |
| 4680 | 15.3629 | 24.5382 |
| 4735 | 14.3074 | 25.9843 |
| 4737 | 3.44349 | 6.23841 |
| 4858 | 10.9969 | 16.3836 |
| 4883 | 2.90539 | 4.33747 |
| 4907 | 15.6828 | 23.3187 |
| 4909 | 1.27248 | 1.88857 |
| 4930 | 11.6747 | 27.1523 |
| 4931 | 13.6777 | 23.3445 |
| 4950 | 14.2889 | 23.7954 |
| 4951 | 13.7678 | 22.9999 |
| 5015 | 16.2411 | 23.9126 |
| 5016 | 23.3837 | 34.3499 |
| 5141 | 0.772023 | 1.12422 |
| 5142 | 15.813 | 23.0934 |
| 5216 | 16.0451 | 23.6257 |
| 5217 | 16.2356 | 23.9009 |
| 5245 | 16.0138 | 23.5223 |
| 5246 | 16.4769 | 24.1366 |
| | 6.4 day | |

How the interaction happens

16

CHAPTER 1. INTRODUCTION

The dark-photon dark matter is non-relativistic and interacts with ordinary matter mostly through the photo-electric process in which a photon (with energy $m_{A'}$) is captured by an atom, with atomic number Z , with a cross section given, for ordinary photons, by

$$\sigma_{p.e.} = 4\alpha^4 \sqrt{2} Z^5 \frac{8\pi r_e^2}{3} \left(\frac{m_e}{\omega} \right)^{7/2}, \quad (1.33)$$

where ω is the photon energy and r_e the classical radius of the electron $r_e = \alpha/m_e$. The cross section for the dark photons is that of ordinary photons rescaled by the mixing parameter ε :

$$\sigma_{A'} = \varepsilon^2 \sigma_{p.e.}. \quad (1.34)$$

This scenario is made accessible to the experiments by considering the rate of absorption of the dark photon by the detector [93, 94]:

$$\Gamma_{A'} = \frac{\rho_{A'}}{m_{A'}} \sigma_{A'} v_{A'}, \quad (1.35)$$

where the density $\rho_{A'}$ is estimated from the relic density (or the flux from the Sun).