SSMB Online Workshop, Tsinghua, 2020-12

Beam Quality Preservation in a Triplet Bend Achromat (TBA) with stable optics

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Congratulations!

- Success of the phase I experiment
- Impressive progress on key physics and technology R&D
- More and more close to be reality
- Wish great success of planned SSMB project

Before entering the topic

- The presented work has some similarity to the SSMB (strong focusing) lattice design
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What we want: when a high quality beam (low emittance, short bunch length) passes thorough a beam line, one can keep the beam distribution in 3D space (except in d dimension) as much as possible.

Previous TBA design study

• In Ref. [1], Venturini have designed a TBA satisfying

– min. α_p (or R56 in Linac notation) & min. $\Delta\epsilon$ (due to CSR)



[1] M. Venturini Phys. Rev. Accel. Beams 19, 064401 (2016).

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$$\begin{array}{c}
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R^{S_{l} \rightarrow S} & R^{S \rightarrow S_{f}} \\
\hline R_{B} & M & R_{B} & M & R_{B} \\
\hline L_{B}, \theta & L_{B}, \theta & L_{B}, \theta \\
\hline s_{i} & s_{1} & s_{2} & s_{3} & s_{4} & s_{f} \\
\hline A^{I} & A^{II} & A^{II} & R_{B} \\
\hline R^{S_{i} \rightarrow S_{f}} & R_{B} = \begin{pmatrix} 1 & L_{B} & 0 & L_{B}\theta/2 \\
0 & 1 & 0 & \theta \\
\theta & L_{B}\theta/2 & 1 & L_{B}\theta^{2}/6 \\
0 & 0 & 0 & 1 \\
\end{pmatrix} \quad M = \begin{pmatrix} m_{11} & m_{12} & 0 & 0 \\
m_{21} & m_{22} & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & 0 & 0 & 1 \\
\end{pmatrix}$$

To satisfy the above two conditions, the transfer matrix from the exit of the 1st dipole to the entrance of the 2nd dipole should have

$$m_{11} = 7 / 4$$

$$m_{12} = -9L_B / 8$$

$$m_{21} = 15 / (2L_B)$$

$$m_{22} = -17 / 4$$

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Transfer matrix of TBA:

/139	$117L_B$	0	0	
4	8	0		٦
165	139	0		F
$2L_B$	4	0	0	C
0	0	1	0	
/ 0	0	0	1/	

Not applicable to multi-TBA beamline

 If with only one of such cell with Tr(M) > 2, it is Ok.



Not applicable to multi-TBA beamline

- If with only one of such cell with Tr(M) > 2, it is Ok.
- But, if having multi such cells, it will be difficult to control the optical functions and emittance growth.







Two TBA cells. left: emittance, middle: β_{x}

CSR immune, isochronous & periodic stable?

- Three conditions instead of two conditions
 - min. α_p (or R56 in Linac notation) & min. Δε (due to CSR) & Tr(M_{total}) <= 2

CSR immune, isochronous & periodic stable?

- Three conditions instead of two conditions
 - min. α_p (or R56 in Linac notation) & min. Δε (due to CSR) & Tr(M_{total}) <= 2
- After derivations (ignored here), we found a **new** solution



The new solution: a singular point

• An example with $L_B = 0.4$ m



 Exactly at the point (m₁₁, m₂₁[m⁻¹]) = (-2, 0), the periodic beta functions are very large or even unstable.



• But near the point, one can find solutions satisfying periodic stability criterion, and having small R56 and small emittance growth induced by CSR.



- To verify this result, we did PSO optimization of practical lattice, with eight variables and two objectives: α_p (and high α_p) & $\Delta\epsilon$ due to CSR
- Of course, some practical constraints were considered, to ensure the periodic optical parameters are at a reasonable level.



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- Final solutions converged to a small area close to the point (m₁₁, m₂₁[m⁻¹]) = (-2, 0)
- Accord with the analytical prediction very well



High order $\alpha_{\rm p}$ also optimized



High order α_p also optimized



Further practical lattice optimization

• For 1 GeV beam, such a TBA can be realized with a cell length of about 8 m.



Another outcome: mitigate MBI effect

- Another outcome is that we find near the point, the microbunching instability (MBI) can be well controlled.
- According to the previous MBI studies [2], to mitigate the MBI effect, it is preferred to have
 - phase advance close to integer times of π between adjacent dipoles
 - moderate beta function

[2] Tsai, C.-Y., S. Di Mitri, D. Douglas, R. Li, and C. Tennant. "Conditions for Coherent-Synchrotron-Radiation-Induced Microbunching Suppression in Multibend Beam Transport or Recirculation Arcs." *Physical Review Accelerators and Beams* 20, no. 2 (February 22, 2017): 024401.

Another outcome: mitigate MBI effect

- Another outcome is that we find nearby the point, the microbunching instability (MBI) can be well controlled.
- According to the previous MBI studies [2], to mitigate the MBI effect, it is preferred to have
 - phase advance close to integer times of π between adjacent dipoles
 - moderate beta function

Near the point, the phase advance of the adjacent dipoles in the TBA cell is intrinsically close to π .

[2] Tsai, C.-Y., S. Di Mitri, D. Douglas, R. Li, and C. Tennant. "Conditions for Coherent-Synchrotron-Radiation-Induced Microbunching Suppression in Multibend Beam Transport or Recirculation Arcs." *Physical Review Accelerators and Beams* 20, no. 2 (February 22, 2017): 024401.



Another outcome: mitigate MBI effect

• Studies showed that even with 30 repetitive TBA cells (totally 360 degrees), the MBI gain factor can be well controlled to a sufficiently small level.



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work to be submitted
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Summary

- This study suggests a new direction of TBA cell design which could promise low α_p (and higher order terms) and small emittance growth induced by CSR, and satisfy the periodic stability criterion.
- It is also found that such a design is helpful in **mitigating the MBI effect**.

Summary

• We hope such a TBA design would be helpful in different accelerator beam line designs.







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Thanks!

• backups

纵向相空间动力学:考虑非线性项
$$z_f = z_i + R_{56}\delta_i + T_{566}\delta_i^2 + U_{5666}\delta_i^3 + O(\delta_i^4)$$

$$\alpha_1 = \frac{1}{L_0} \oint \frac{\eta_1(s)}{\rho} ds$$

$$\Delta \mathbf{L} = \mathbf{R}_{56} \boldsymbol{\delta}_i + \mathbf{T}_{566} \boldsymbol{\delta}_i^2 + \boldsymbol{U}_{5666} \boldsymbol{\delta}_i^3 + \boldsymbol{O}(\boldsymbol{\delta}_i^4)$$

$$\alpha_{2} = \frac{1}{L_{0}} \int \frac{\eta_{1}'(s)}{2} + \frac{\eta_{2}(s)}{\rho} ds$$

$$\frac{\Delta L}{L_0} = \alpha_1 \delta_i + \alpha_2 \delta_i^2 + \alpha_3 \delta_i^3 + O(\delta_i^4)$$

$$egin{array}{c} lpha_1
ightarrow 0 \ lpha_2
ightarrow 0 \end{array}$$

$$\begin{pmatrix} R_{11} & R_{12} & R_{13} & R_{14} & R_{15} & R_{16} \\ R_{21} & R_{22} & R_{23} & R_{24} & R_{25} & R_{26} \\ R_{31} & R_{32} & R_{33} & R_{34} & R_{35} & R_{36} \\ R_{41} & R_{42} & R_{43} & R_{44} & R_{45} & R_{46} \\ R_{51} & R_{52} & R_{53} & R_{54} & R_{55} & R_{56} \\ R_{61} & R_{62} & R_{63} & R_{64} & R_{65} & R_{66} \end{pmatrix}$$





图1. Coherent Synchrotron Radiation(CSR)效应的物理图像[1]



图2. CSR效应引起的束团相空间分布与束团密度分布图像[2]



CSR bursting 流强不稳定阈值[3]

 G. Wüstefeld, BESSY, Short Bunches & CSR, EPAC'08, June 23rd. 2008
 Marco Venturini, Robert Warnock, Ronald Ruth, and James A. Ellison Phys. Rev. ST Accel. Beams 8, 014202 – Published 28 January 2005

[3] J. M. Byrd, W. P. Leemans, A. Loftsdottir, B. Marcelis, Michael C. Martin, W. R. McKinney, F. Sannibale, T. Scarvie, and C. Steier Phys. Rev. Lett. 89, 224801 – Published 8 November 2002

CSR cancellation in TBA

Model
 2D CSR kick model
 Requirements

 achromatic condition
 CSR cancellation



Fig. 1. Schematic layout of a symmetric TBA and the corresponding physical model of the CSR effect in a TBA with three point-kicks.

[4] Huang, Xiyang & Jiao, Yi & Xu, Gang & Cui, Xiaohao. (2015). Suppression of the emittance growth induced by coherent synchrotron radiation in triple-bend achromats. Chinese Physics C. 39. 10.1088/1674-1137/39/5/057001



$$M_{12} = \begin{pmatrix} -\frac{r_2\rho + 2m_{12}(\theta_1 + \theta_2)S_1}{2r_1\rho} & m_{12} \\ \frac{1}{m_{12}} \left(\frac{r_2S_2}{4r_1S_1} + \frac{m_{12}(\theta_1 + \theta_2)S_2}{2r_1\rho} - 1 \right) - \frac{S_2}{2S_1} \end{pmatrix}, (16)$$



$$X_{k} = \begin{pmatrix} x_{k} \\ x_{k}' \end{pmatrix} = \begin{pmatrix} \rho^{1/3} \kappa [\theta \cos(\theta/2) - 2\sin(\theta/2)] \\ \sin(\theta/2) (2\delta + \rho^{1/3} \theta k) \end{pmatrix}$$

CSR效应引起的发射度增长:

$$\begin{split} \epsilon &= \sqrt{\left(\epsilon_0\beta_f + x_{f,rms}^2\right)\left(\epsilon_0\gamma_f + {x'}_{f,rms}^2\right) - \left(\epsilon_0\alpha_f - x_{f,rms}^2{x'}_{f,rms}^2\right)^2} = \sqrt{\epsilon_0^2 + \epsilon d\epsilon} \\ d\epsilon &= \gamma_f x_{f,rms}^2 + 2\alpha_f x_{f,rms} x'_{f,rms} + \beta_f {x'}_{f,rms}^2 \end{split}$$

work to be submi⁵ Y Jiao, X. Cui, X. Huang, and G. Xu, Phys. Rev. ST Accel. Beams 17, 060701 (2014).